## Dark Matter & the CMB

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# *Outline*

- Particle Dark Matter Effects on the CMB
- $DM \leftrightarrow CMB$  anisotropies (Ionization)
- $DM \leftrightarrow 21$ cm (Temperature)
- <sup>l</sup> Forecasts for DM and PBHs
- Inhomogeneity from DM heating

### *Dark matter out there…*



 $\mathbf{f}$  $\overline{v}$  $\rightarrow$ 

SCIENCE

Heart of darkness: Scientists probe dark matter near Milky Way's core



#### DM got'be `matter', right?



## *Theory orders are placed.*



CMB covers a wide DM mass range!

#### *CMB and Dark Matter – how?*



# *DM probes from the CMB*

• CMB spectral distortion:

` coupled' DM, early/steady energy injection, DM-photon conversion, etc

- CMB polarization: pol. rotation in CPV medium
- CMB derivatives:

21cm maps of matter power-spectrum: spatial & temperature distributions

# *DM probes from the CMB*

• CMB spectral distortion:

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- CMB polarization: pol. rotation in CPV medium
- CMB derivatives:

21cm maps of matter power-spectrum: spatial & temperature distributions

## *Impact from steady (high-energy) injection*

- Deposit energy into IGM during the dark age of Universe
- (1) Ionize (fraction of) the IGM; (2) Heats the IGM
- A small energy budget for a large impact

On decay lifetime:



#### *The ` standard'ionization history*

Standard ionization evolution (pre-EoR)

$$
\frac{dX_e}{dt} = \left\{ (1 - X_e)\beta - X_e^2 n_b \alpha^{(2)} \right\}
$$
\nIonization rate (by radiation field):

\n

$$
\beta \equiv \langle \sigma v \rangle \Big(\frac{m_e T}{2\pi}\Big)^{3/2} e^{-\epsilon_0/T}
$$

 $\sim$ 

Recombination:

$$
\alpha^{(2)}\equiv\langle\sigma v\rangle
$$

Approx. capture rate<sub>2</sub>)  
to a non-ground state = 
$$
9.78 \frac{\alpha^2}{m_e^2} \left(\frac{\epsilon_0}{T}\right)^{1/2} \ln\left(\frac{\epsilon_0}{T}\right)
$$



x<sub>e</sub> reduces to a 10<sup>-4</sup> floor during the cosmic dark age and returns to unity during EoR

#### *DM Effect 1: ionization*

●More free electrons

•More CMB scattering  $\rightarrow$  Damping on  $C_l$ 

 $\frac{dX_e}{dt} = \left\{ (1 - X_e)\beta - X_e^2 n_b \alpha^{(2)} \right\}$ 

SM: H atom ionization and recombination

"Deposit Channels"

ionization from ground state

ionization from excited states

(+ other channels)

*The ` perturbed'ionization history*



"Broadens the last scattering surface"

#### *Redshift dependence in injection rate*

- Annihilation and/or Decay of WIMPs
- Energy release during dark ages

**DM Annihilation**: fast during high z,

 $\sim (z+1)^6$ 

Late time density clustering boosts the annihilation rate after  $z \sim O(50)$ 

$$
\left(\frac{dE}{dVdt}\right)_{\text{INJ}}^{\text{ann},\text{boosted}} = [1 + B(z)] \left(\frac{dE}{dVdt}\right)_{\text{INJ}}^{\text{ann}}
$$

$$
B(z) = \frac{\Delta_c \rho_c}{\rho_{\text{DM}}^2} \int_{M_{\text{min}}}^{\infty} MB_{\text{h}}(M) \frac{dn}{dM} dM
$$

**DM Decay**: a steady rate, unaffected by structure formation

 $\sim (z+1)^3$ 

## *Lagged energy deposition*

Injected high-energy particles lose energy by scattering, ionization, excitations, etc...

> Not instantaneously deposited into the IGM if particles are energetic  $(E >> KeV)$ : \* accumulative over earlier injection \* efficiency reduces at later time

Energy "fraction" into ionization (of H)



#### Numerical calculation

Implemented into

new physics induced excitation, scattering ter Lyman-α photons, etc.

Also see: Belotsky, Kirillov 2015 13 <sup>l</sup> Compute a history-dependent deposit "efficiency" *f* (E,z)

$$
f_c \equiv \left(\frac{dE_{tot}}{dVdt}\right)_{DEP} / \left(\frac{dE_{tot}}{dVdt}\right)_{INJ}
$$
 (from theory)

(from si

$$
f_{\rm c}(z_{\rm i}) \approx \frac{\sum_{s}\sum_{\rm j}\sum_{\rm k}E_{\rm j}^{s}I^{s}(z_{k},E_{\rm j}^{s}){\rm d}V(z_{\rm k}){\rm d}t(z_{\rm k})T_{\rm c,ijk}^{s}{\rm d}E_{\rm j}^{s}}{\sum_{s}\sum_{\rm j}E_{\rm j}^{s}I^{s}(z_{i},E_{\rm j}^{s}){\rm d}E_{\rm j}^{s}{\rm d}V(z_{\rm i}){\rm d}t(z_{\rm i})}
$$

- <sup>l</sup> Averaged over injection spectra *(j)* and species *(s)* and integrated over all previous redshift  $(z_k > z_i)$
- Electrons are more effective than gamma rays at large energy
- Photons extends to (much) lower mass range

Simulated eff.  $T_{ijk}$ `DarkHistory' Liu, Ridgway, Slatyer, 19'

#### *DM: impact on xe*

Annihilation: raises the  $x_e$  floor,

Decay: steady rise in  $x_e$ 



## *Xe on CMB Cl: damping & shift*



### *Current limits: WIMP annihilation*

Planck Collaboration: Cosmological parameters



Fig. 46. Planck 2018 constraints on DM mass and annihilation cross-section. Solid straight lines show joint CMB constraints on several annihilation channels (plotted using different colours), based on  $p_{\text{ann}} < 3.2 \times 10^{-28} \text{ cm}^3 \text{ s}^{-1} \text{ GeV}^{-1}$ . We also show the  $2\sigma$ preferred region suggested by the AMS proton excess (dashed ellipse) and the Fermi Galactic centre excess according to four possible models with references given in the text (solid ellipses), all of them computed under the assumption of annihilation into  $b\bar{b}$ (for other channels the ellipses would move almost tangentially to the CMB bounds). We additionally show the  $2\sigma$  preferred region suggested by the AMS/PAMELA positron fraction and *Fermi*/H.E.S.S. electron and positron fluxes for the leptophilic  $\mu^+\mu^-$  channel (dotted contours). Assuming a standard WIMP-decoupling scenario, the correct value of the relic DM abundance is obtained for a "thermal cross-section" given as a function of the mass by the black dashed line.

**PLANCK 18: `Cosmological parameters' Thermal WIMP mass limit: 10~30 GeV** 17

#### *PLANCK 18: Pol. data lifts EoR degeneracy*



## *Pol. EE peak shifts make the call*



Data: PLANCK18

TABLE I. Linear correlation coefficients between  $\langle \sigma v \rangle / m_{\chi}$ and cosmological parameters corresponding to Fig.3.

Polarization ani.  $C_l$ TE, EE peak location shift sensitive to higher z (~recombination) effects

# *Near future: How about more pol. data*



![](_page_19_Picture_2.jpeg)

![](_page_19_Picture_3.jpeg)

![](_page_19_Picture_52.jpeg)

+ BICEP3 data available

## *AliCPT: China ' s upcoming CMB pol. observatory*

![](_page_20_Figure_1.jpeg)

![](_page_21_Figure_0.jpeg)

#### *Forecast on WIMP lifetime (decay to photons)*

![](_page_22_Figure_1.jpeg)

![](_page_23_Figure_0.jpeg)

| Experiment          |      |     |
|---------------------|------|-----|
| Planck              | 24   | 85  |
| AdvACTPol           | 0.68 | 4.7 |
| <b>AliCPT</b>       | 21   | 78  |
| $A$ liCPT+ $Planck$ | 16   | 53  |
| <b>CLASS</b>        | 5.5  | 30  |
| Simons Array        | 0.35 | 1.5 |
| Simons Observatory  | 0.92 | 4.2 |
| $SPT-3G$            | 2.2  | 9.9 |

TABLE II. 95% C.L. upper limit on  $\Gamma_{\chi}$  (in  $10^{-26}$  s<sup>-1</sup>) at  $m_{\chi} = 10$  GeV.

| Experiment                  | $^{\circ}e$         | $\rightarrow \gamma \gamma$<br>$\chi\chi$ |
|-----------------------------|---------------------|-------------------------------------------|
| Planck                      | 39                  | 32                                        |
| <i>Planck</i> - Unclustered | 39                  | 33                                        |
| AdvACTPol                   | 330                 | 330                                       |
| AliCPT                      | 32                  | 22                                        |
| $A$ liCPT+ $Planck$         | 51                  | 42                                        |
| <b>CLASS</b>                | 49                  | 37                                        |
| Simons Array                | $1.1 \times 10^{3}$ | $1.0 \times 10^{3}$                       |
| Simons Observatory          | 310                 | 290                                       |
| $SPT-3G$                    | 140                 | 130                                       |

TABLE III. Expected 95% C.L. lower limit on  $m<sub>x</sub>$  (in GeV) assuming a thermal relic's annihilation cross-section  $\langle \sigma v \rangle = 3 \times 10^{-26}$  cm<sup>3</sup>/s.

## *Low mass PBHs, radiation*

has a dE/dt $\sim$ (1+z)<sup>3</sup> history  $\qquad$   $_{10^{-1}}$ 

Significant sensitivity<br>
n relevant mass range:  $\frac{10^{-2}}{\frac{10^{-3}}{\epsilon^2}}$ in relevant mass range:

PLANCK15 constraint: S.Clark., B.Dutta., Y.Gao, Y-Z.Ma,  $10^{-7}$ <br>L.E. Striggri 1612.07738 L.E. Strigari, 1612.07738  $10^{10}$   $10^{15}$ 

PLANCK18 & forecasts:<br>Fxtended RH mass distributions see: 10<sup>0</sup> Extended BH mass distributions, see: J.Cang., Y.Gao., Y-Z. Ma., 2011.12244

![](_page_25_Picture_153.jpeg)

![](_page_25_Figure_7.jpeg)

## *BH accretion radiation (solar-supermassive)*

![](_page_26_Figure_1.jpeg)

## *About (E-mode) Pol. Sensitivity...*

- Mostly via extra ionization, breaks degeneracy in  $\tau$  looks good!
- (Annihilation) Not very sensitive to clustering boost

### *Remaining Issue: EoR uncertainty (*τ*) washes out late-time DM injection*

![](_page_28_Figure_1.jpeg)

FIG. 4: The state-of-the-art measurement on  $x_{\text{H}_{I}}(z)$ , taken from Table II. The black and red dashed lines are two examples of the "tanh" model which cannot fit the data very well.

**Current Pol. data sensitivity MOSTLY** derives from injection right-after recombination time

EoR uncertainty needs future exp. input

![](_page_28_Figure_5.jpeg)

![](_page_29_Figure_0.jpeg)

Remember this bump?

poor low-z sensitivity due to EoR We need a late-time handle.

#### *DM ef ect #2: IGM temperature*

![](_page_30_Figure_1.jpeg)

## *IGM heating with DM*

• Injected particles raise IGM temperature

Scattering with bkg radiation  $\propto$  T<sub>CMB</sub> - T<sub>IGM</sub>

## *We may hear a lot from 21cm ...*

#### ●**Precision reionization history**:

Ionization fraction  $x_e$ , mean temperature  $T_G$ 

#### ●**Distribution of neutral Hydrogen gas**

temperature map & power spectrum

![](_page_32_Figure_5.jpeg)

![](_page_32_Figure_6.jpeg)

![](_page_32_Figure_7.jpeg)

#### Projected power spectrum sensitivities (from SKA white paper)  $33$

## *Neutral Hydrogen 21cm line*

Hyper-fine split between the singlet and triplet states of neutral Hydrogen atom

Spontaneous 21cm transition very slow: 10<sup>7</sup> yr

![](_page_33_Figure_3.jpeg)

$$
N_1/N_0 = 3 \, \mathrm{e}^{-0.068 \mathrm{K}/T_\mathrm{S}}
$$

Spin temperature determined by relative effectiveness between 21cm transitions, conversions by kinetic collisions, coupling to the CMB and Lyman- $\alpha$  photons.

$$
T_{\rm S} = \frac{T_{\rm CMB} + y_{\rm c}T_{\rm G} + y_{\rm Ly\alpha}T_{\rm Ly\alpha}}{1 + y_{\rm c} + y_{\rm Ly\alpha}},
$$

$$
y_{\rm c} = \frac{C_{10}}{A_{10}} \frac{T_{\star}}{T_{\rm G}},
$$

$$
y_{\rm Ly\alpha} = \frac{P_{10}}{A_{10}} \frac{T_{\star}}{T_{\rm Ly\alpha}},
$$

CMB develops [dark] absorption lines by running through neutral Hydrogen gas clouds with  $T_S < T_{CMB}$ .

 $\rightarrow$  Slices of high-z universe

#### **CMB's 21cm absorption windows**

(1) neutral Hydrogen presence  $(2)$  T<sub>s</sub> cooler than the CMB

Dark age window

![](_page_34_Picture_3.jpeg)

Early reionization window (first discovery claim from EDGES) Bowman, et.al. Nature 555, <sup>67</sup> (2018). Gas temperature decouples from  $CMB$   $\sim$  200  $C$   $Gyr$   $Gyr$   $C1$   $Gyn$   $FDCFS$   $35$ 

# *T<sup>21</sup> dependencies...*

- 21cm brightness relies on IGM temperature evolution
- Direct  $T_{\text{GAS}}$  measurements.

ionization

Gas density distribution

Optical depth: Cosmology modeldependent

Wouthuysen-Field:  $T_{spin} \sim T_{lya} \sim T_{GAS}$ <br>at cosmic dawn

![](_page_35_Figure_7.jpeg)

Gas spin temperature against CMB

### *Temperature evolution*

![](_page_36_Figure_1.jpeg)

21cm absorption lines whenever  $T_S$  is lower than  $T_{CMB}$ . Temperature sim. with HyRec

DM induced heating can suppress / erase the 21cm signal

![](_page_37_Figure_1.jpeg)

The average `brightness temperature' Z

#### *EDGES: claim of 21cm*

 $0.0$ 

 $-0.2$ 

 $-0.4$ 

 $-0.6$ 

 $-0.8$ 

60

EDGES 2018 J. D. Bowman,et.al. Nature 555, 67 (2018).

2020 (summer)

![](_page_38_Figure_3.jpeg)

![](_page_38_Figure_4.jpeg)

![](_page_38_Picture_5.jpeg)

EDGES: A Discovery

**~ Twice the LCDM signal !**

near 78 MHz? LOFAR & MWA (by 2020) Upper limits only.

Mid-Band

100

80

90

 $\nu$  [MHz]

70

Bowman et al. (2018)

110

120

 $(f)$ 

#### *WIMP involvement?*

DM cooling **(DM is cooler)**

Lower gas temperature via collisions: more 21cm signal

Explains the EDGES data \*needs large scattering xsec

DM heating **(DM releases energy)**

Raises gas temperature by energy injection: reduces 21cm signal

Most stringent bounds on DM annihilation, decays & other energy injections

**Large uncertainty at low frequency**; radio-frequency \*new physics

CMB uncertainties Non-standard cosmology Modified Friedmann Eq.<br>Dynamic DE, etc

$$
T_{21} \propto \frac{1}{H(z)} \left( 1 - \frac{T_{\gamma}}{T_{S}} \right)
$$

## *WIMP cooling as an explanation to the EDGES data*

![](_page_40_Figure_1.jpeg)

Milli-charged DM constrained to MeV range and tiny (<1%) fractions of relic density E.D.Kovertz, et.al. 18'

\* subleading abundance is OK if millicharged DM also has long-range force with the rest of DM (H.Liu, Outmezguine, Redigolo, Volansky 1908.06986)

*Discovery of 21cm means high WIMP sensitivity*

![](_page_41_Figure_1.jpeg)

Unlike CMB pol., 21cm is VERY sensitive to DM clustering boost

#### *WIMP lifetime bound @ 21cm discovery*

**Limit on TGAS rise: ΔT<sup>21</sup> < +100 or +150 mK at z=17**

S.Clark, B.Dutta, Y.Gao, Y.-Z.Ma, L.E.Strigari, 18'

![](_page_42_Figure_3.jpeg)

### *21cm has great expectations…*

![](_page_43_Picture_1.jpeg)

![](_page_43_Picture_2.jpeg)

#### 中国签约参与SKA项目

实验建设:SKA 一期项目的总的预计为6.5 亿 欧元,相当于50亿人民币左右。将由中国、南 非、英国、澳大利亚、荷兰、意大利、葡萄 牙,加拿大、德国、印度等共同承担。 \*中国出资规模占重要比例。 \*中电集团54研究所负责133盏天线建造。

#### *21cm with DM: Inhomogeneous heating…*

- w/o DM: inhomogeneity from matter fluctuations
- w DM:  $x_e$ , *T* also become inhomogeneous esp. for ρ<sup>2</sup> enhanced annihilation & quick *E* deposit
- Potential correction to 21cm spectrum (v.s. global signal)

SM   
\n
$$
\frac{dx_e}{dt} = C \left( \beta (1 - x_e) e^{-E_\alpha/k_B T} - x_e^2 n_H \alpha \right) + \frac{1}{n_H E_i} J_{d,i} + \frac{1 - C}{n_H E_\alpha} J_{d,\alpha}
$$
\n
$$
+ \text{transport terms, etc.}
$$
\n
$$
\frac{dT}{dt} = \frac{2T}{1 + z} \frac{dz}{dt} + \frac{8\sigma_T a_R T_\gamma^4}{3m_e c^2} \frac{x_e}{f_b} (T_\gamma - T) + \frac{2}{3k_B n_H f_b} J_{d,h}
$$

#### $x_e \& T$  inhomogeneity potentially affect the 21cm power spectrum

![](_page_45_Figure_1.jpeg)

(Mpc/pixel, deposit terms only, instantaneous deposition)

### *Dark Matter and CMB*

![](_page_46_Figure_1.jpeg)

AdvACTPol, AliCPT, Simons Obs., SPG3, S4, etc.& 21cm coming clo<sup>g</sup>e?

# BACKUP: light bosonic DM

#### *CMB on very light bosonic DM*

![](_page_48_Figure_1.jpeg)

## BACKUP: partially coupled DM

## *Ef ect 3: DM ` couples 'to matter*

![](_page_50_Figure_1.jpeg)

 $DM$  – matter scattering at low velocity:  $\langle v\sigma \rangle \sim v^{-n}$ Corrections in TT, EE, and lensing spectra

## FORECAST Method

### *Forecast likelihood*

$$
-2\ln \mathcal{L}(\{C_{\ell}\}|\{\hat{C}_{\ell}\}) =
$$
  

$$
f_{\rm sky} \times \sum_{\ell} (2\ell + 1) \{ \text{Tr}[\hat{C}_{\ell} C_{\ell}^{-1}] - \ln |\hat{C}_{\ell} C_{\ell}^{-1}| - 2]
$$

$$
C_{\ell} \equiv \begin{bmatrix} C_{\ell}^{TT} & C_{\ell}^{TE} \\ C_{\ell}^{TE} & C_{\ell}^{EE} \end{bmatrix}
$$

$$
\hat{C}_\ell \equiv \begin{bmatrix} \bar{C}^{TT}_\ell + N^{TT}_\ell & \bar{C}^{TE}_\ell \\ \bar{C}^{TE}_\ell & \bar{C}^{EE}_\ell + N^{EE}_\ell \end{bmatrix}
$$

$$
N_{\ell}^{\text{EE}} = \left[ \sum_{\nu} \omega_{E,\nu} \exp\left( -\ell(\ell+1) \frac{\theta_{\text{FWHM},\nu}^2}{8 \ln 2} \right) \right]^{-1}
$$

$$
N_{\ell}^{\text{TT}} = \frac{1}{2} N_{\ell}^{\text{EE}}
$$

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## *Exp. specifications (DM)*

![](_page_53_Picture_17.jpeg)

<sup>a</sup> AdvACTPol constraints would improve by a factor of 2 if choosing  $\ell_{\rm min} = 60$ .

## *Exp. specifications (PBH)*

![](_page_54_Picture_12.jpeg)