

非微扰方法
及其在高能物理中的应用
专题研讨会

利用格点QCD研究 ρ 介子

Study ρ meson on Lattice QCD

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Paper is preparing ...

中国科学技术大学

2024. 10. 27

合肥

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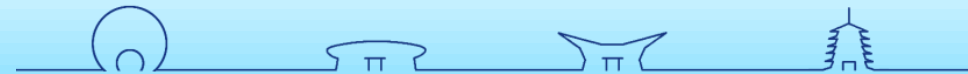
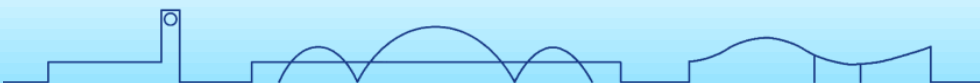


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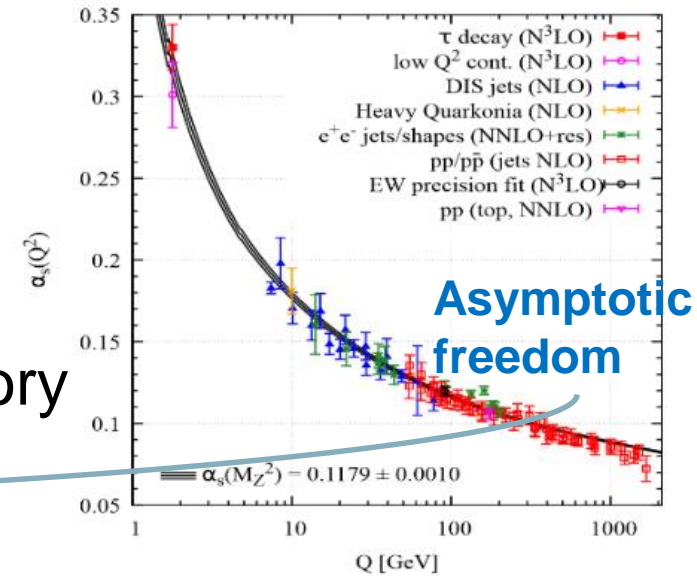
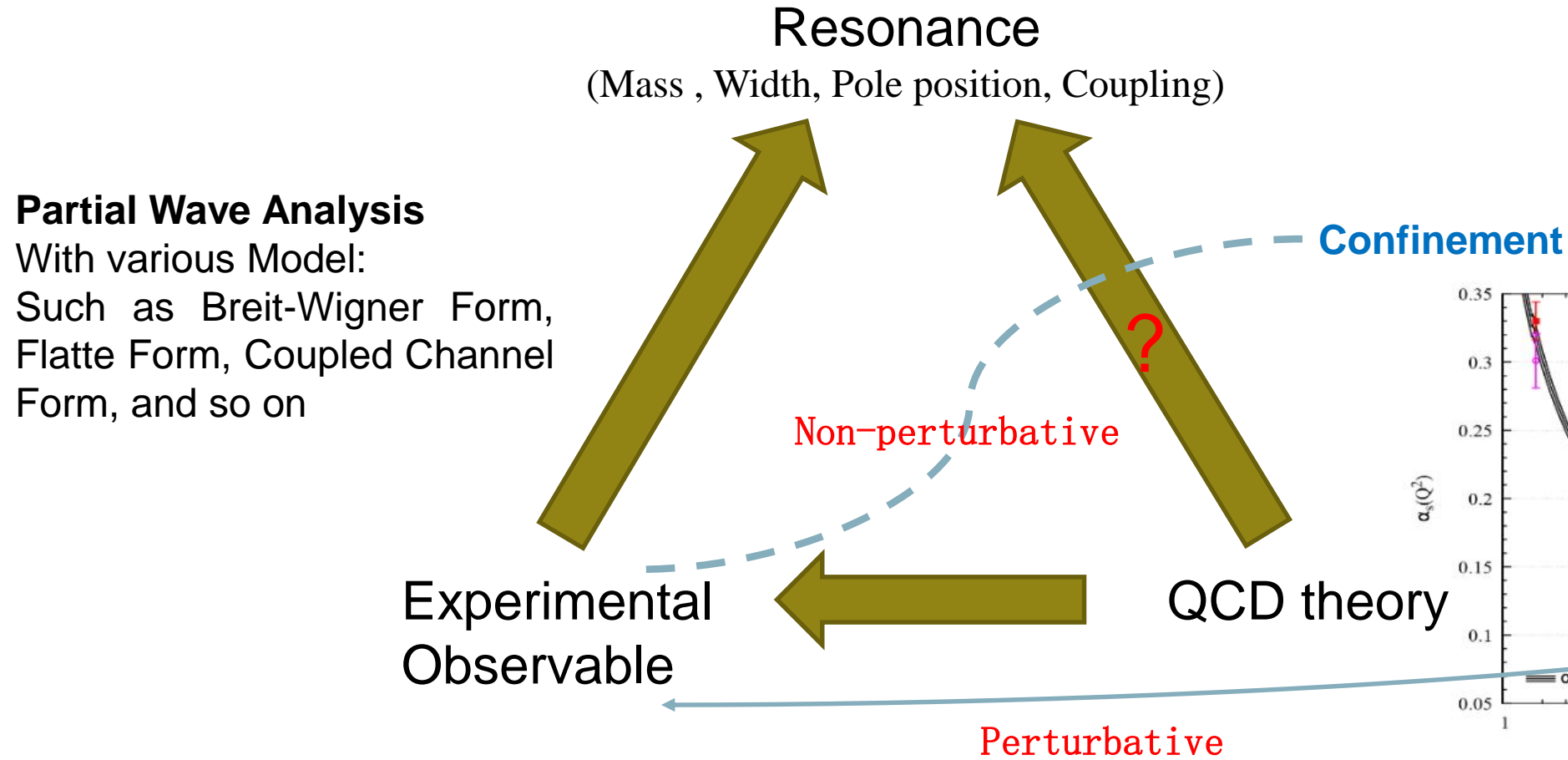


Outline

- Motivation
- Lattice QCD introduction
- ρ meson spectra on LQCD
- Summary



Motivation



Motivation: Spectrum & Scattering

- Spectrum: Table of the energy of the physical states, can reflect the inner structure information.
- Scattering: The line shapes of Cross section or T-matrix show the properties of the states and interaction.
- These observables can constrain the models.



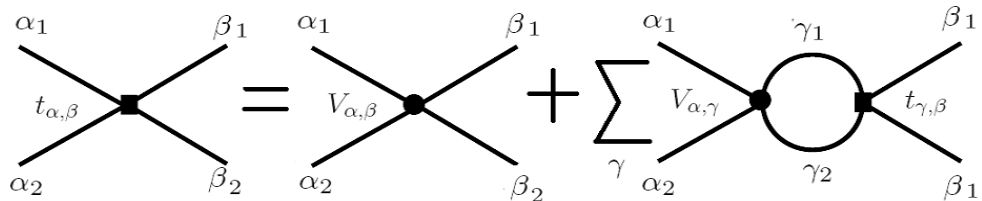
Powerful Parameters, Powerful Model



Motivation: Scattering

Lattice Spectrum !!!
The dimension of m_π

$$t_{\alpha,\beta}(k_\alpha, k_\beta, E) = V_{\alpha,\beta}(k_\alpha, k_\beta) + \sum_\gamma \int k_\gamma^2 dk_\gamma \frac{V_{\alpha,\gamma}(k_\alpha, k_\gamma) t_{\gamma,\beta}(k_\gamma, k_\beta, E)}{E - \sqrt{m_{\gamma 1}^2 + k_\gamma^2} - \sqrt{m_{\gamma 2}^2 + k_\gamma^2} + i\epsilon}$$



A diagram showing a propagator B_i between two vertices. The incoming lines are α_1, α_2 and the outgoing lines are β_1, β_2 . The propagator is represented by a horizontal line connecting two vertices.

$$g_{i,\alpha}^* \frac{1}{E - m_i} g_{i,\beta}$$

A diagram showing a vertex $V_{\alpha,\beta}$ where incoming lines α_1, α_2 meet outgoing lines β_1, β_2 at a single point.

$$V_{\alpha,\beta}$$

$$\Sigma_{ij}(E) = \sum_\alpha \int k^2 dk \frac{g_{i,\alpha}^* G_{i,\alpha}}{E - \sqrt{m_{\alpha 1}^2 + k^2} - \sqrt{m_{\alpha 2}^2 + k^2} + i\epsilon}$$

Λ

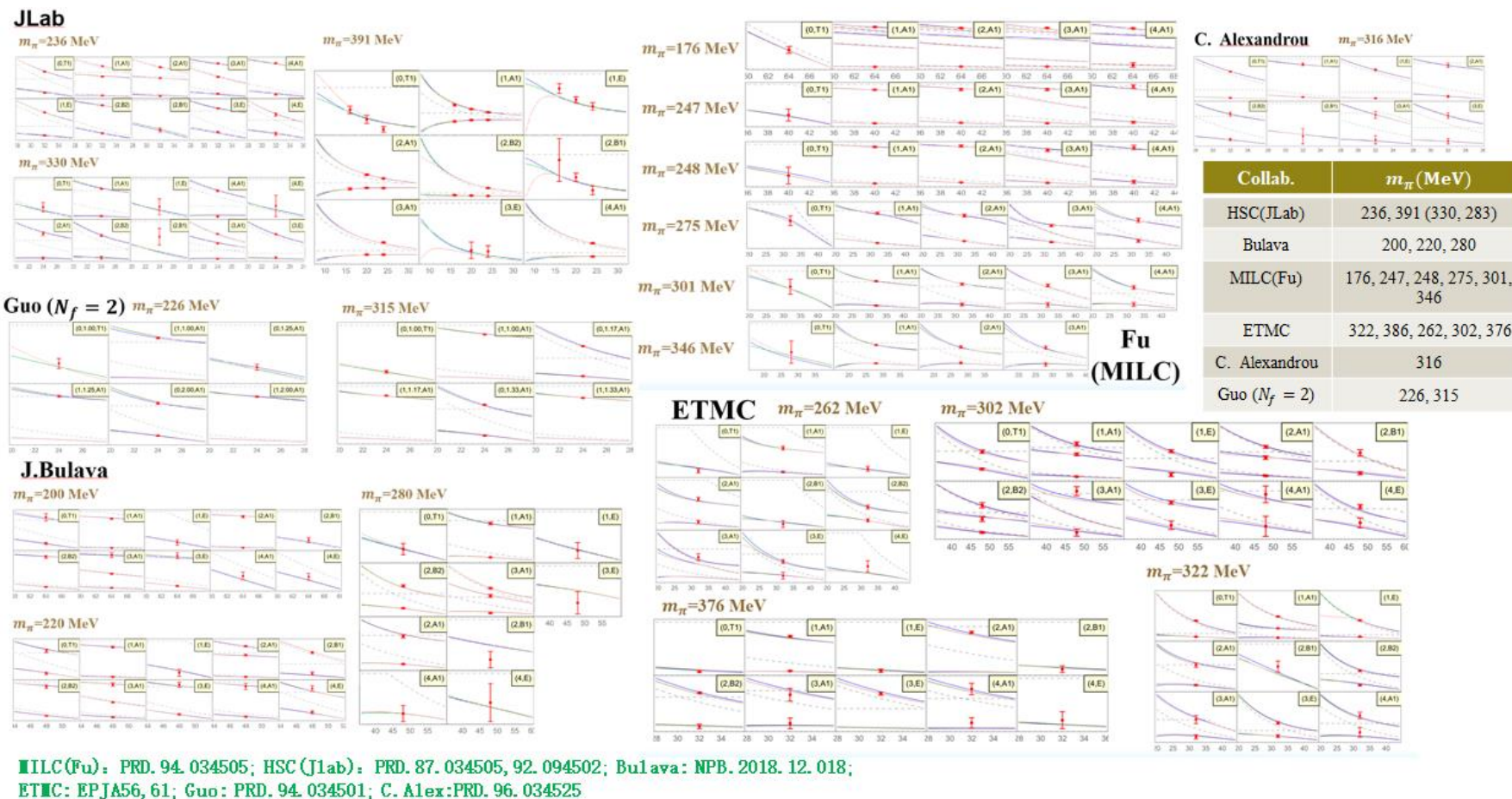
$$G_{i,\alpha} = g_{i,\alpha} + \sum_\beta \int \frac{k^2 dk \tilde{t}_{\alpha\beta} g_{j,\beta}}{E - \sqrt{m_{\beta 1}^2 + k^2} - \sqrt{m_{\beta 2}^2 + k^2} + i\epsilon} \quad \tilde{t}_{\alpha\beta} = v_{\alpha\beta} + \sum_\gamma \int \frac{k_\gamma^2 dk_\gamma v_{\alpha\gamma} \tilde{t}_{\gamma\beta}}{E - \sqrt{m_{\gamma 1}^2 + k_\gamma^2} - \sqrt{m_{\gamma 2}^2 + k_\gamma^2} + i\epsilon}$$

- $T^{-1}(E) \sim E - m_i - \Sigma(E)$
 $= E - m_i - \text{Re}(\Sigma(E)) - i\text{Im}(\Sigma(E))$
- $T^{-1}(E) \sim E - m_i - \text{Re}(\Sigma(E)) + i\Gamma(E)/2$
- $m_{phys} = m_i + \text{Re}(\Sigma(E))$
 $= m_0 + am_\pi^2 + \dots + \text{Re}(\Sigma(E))$



Motivation: π mass dependence

- The ρ is the best one!
- The largest lattice data of ρ .
- The lowest state can decay through strong interaction.
- We almost clearly understand ρ .

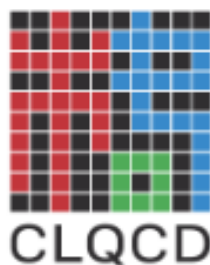


Motivation: π mass dependence

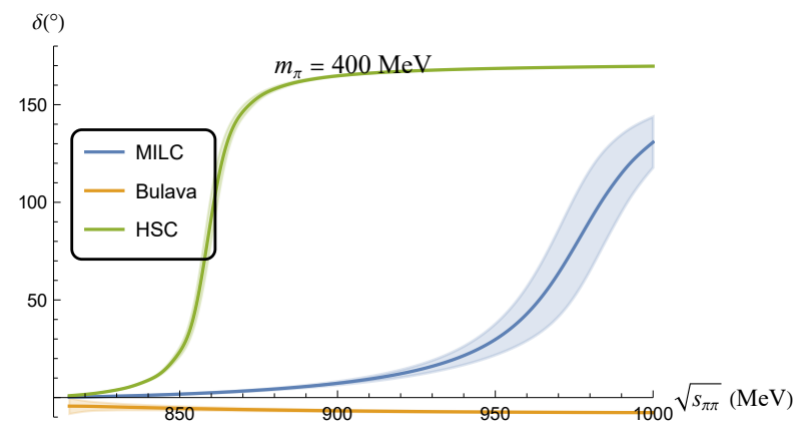
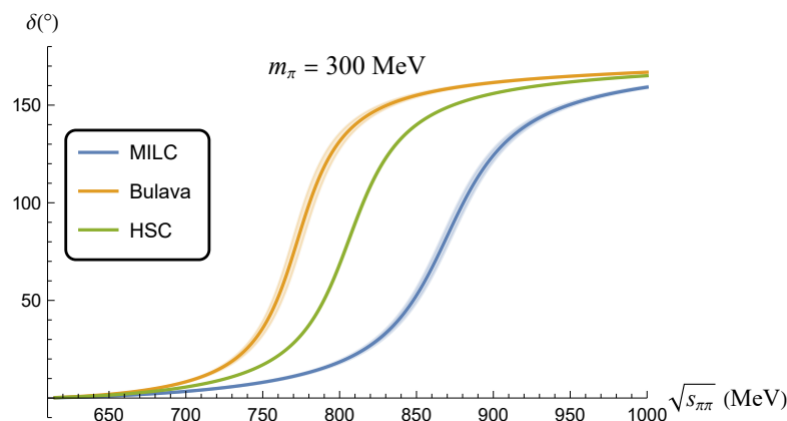
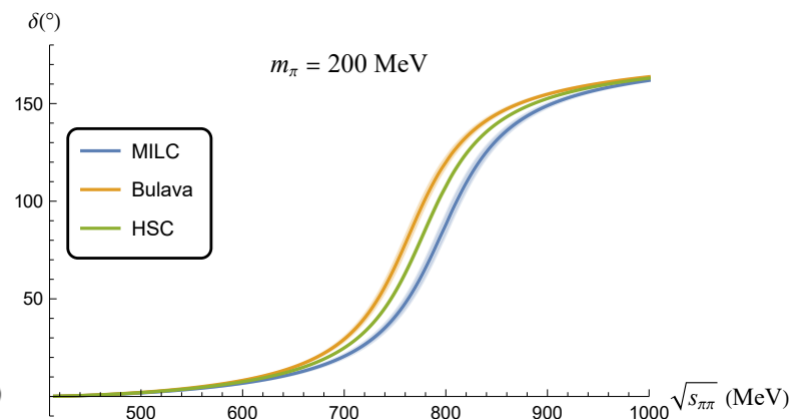
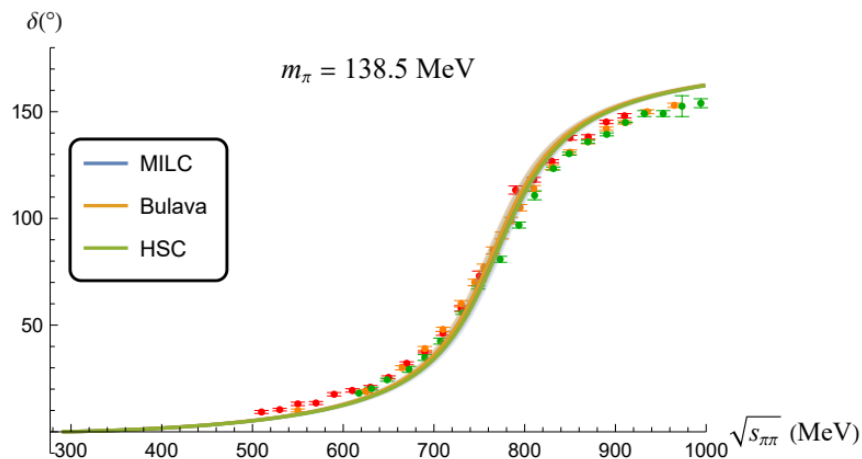
- The ρ is the best one!
- The largest lattice data of ρ .
- The lowest state can decay through strong interaction.
- We almost clearly understand ρ .

However

Lattice spacing effect,
Fermion action effect...

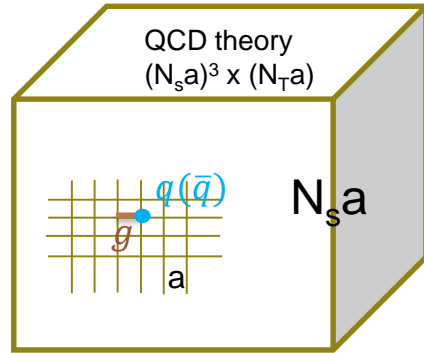


We need a systematic Lattice
data group to study ρ !



Lattice QCD introduction

1. QCD theory: on a box in the Euclid four space



2. $a \rightarrow$ UV cutoff, $N_s a \rightarrow$ Infrared truncation

3. Lattice QCD \rightarrow a model of statistical physics.

$$\langle O \rangle = \int D\phi O[\phi] P[\phi] \quad P[\phi] = \frac{1}{Z} e^{-S[\phi]} \quad Z = \int D\phi e^{-S[\phi]}$$

Φ : field quantity, $S[\phi]$: Action, $O[\phi]$: physical quantity

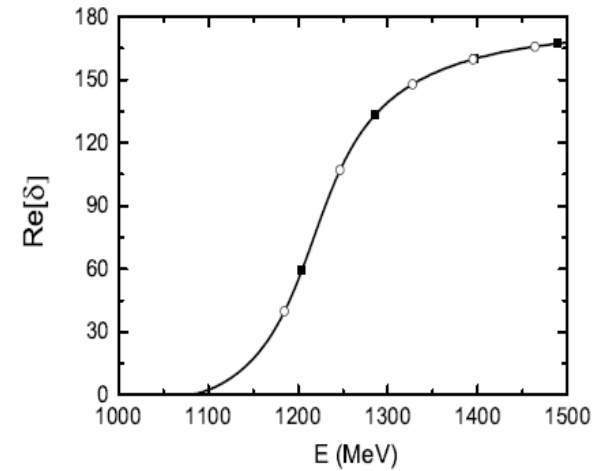
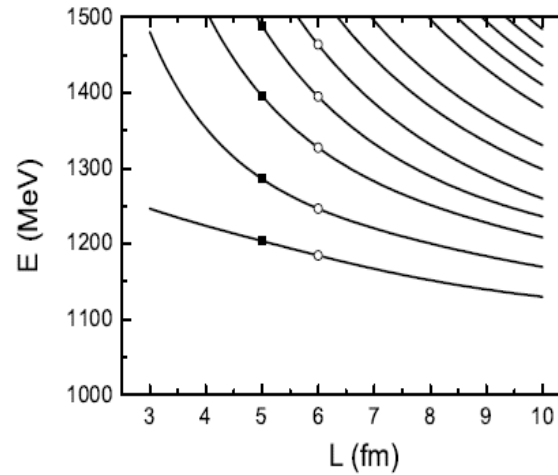
4. Monte Carlo method

5. Three steps for Lattice QCD to real world

a, Configuration

b, Measurement $\sum_{(\vec{y}-\vec{x}) \in \mathbb{Z}^3} e^{i\vec{p} \cdot (\vec{y}-\vec{x})} \langle T(\psi(t; \vec{y}), \psi^\dagger(t; \vec{y})) \rangle \sim \sum_{\Gamma, i} Z_i^\Gamma e^{-E_i^\Gamma t}$

c, Transformation



$$L \text{ — } E \text{ — } \delta_{\pi\pi}(E)$$

$$\delta(k) = \Delta(L) \bmod \pi$$

M. Luscher, NPB 354, 531 (1991).

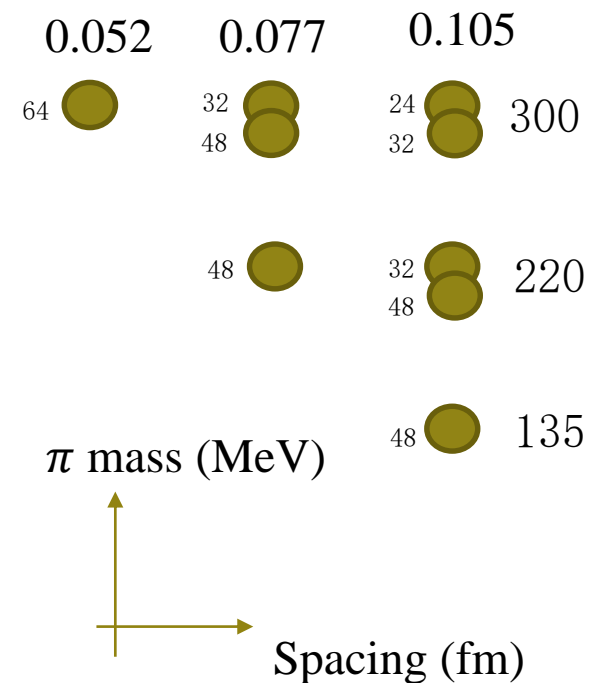


**Amplitude to
Resonance
information**



ρ meson spectra on LQCD -- Configuration

Name	Volume	Spacing	β	m_π/MeV	$m_\pi L$
C24P34	$24^3 \times 64$	0.10530 fm	6.20	340	4.38
C24P29	$24^3 \times 72$			292	3.75
C32P29	$32^3 \times 64$			292	5.01
C32P23	$32^3 \times 64$			228	3.91
C48P23	$48^3 \times 96$			225	5.79
C48P14	$48^3 \times 96$			135	3.56
F32P30	$32^3 \times 96$	0.07746 fm	6.41	303	3.81
F48P30	$48^3 \times 96$			303	5.72
F32P21	$32^3 \times 64$			210	2.67
F48P21	$48^3 \times 96$			207	3.91
H48P32	$48^3 \times 144$	0.05187 fm	6.72	321	4.06
H64P32	$64^3 \times 128$			321	5.41



CLQCD, PRD109
(2024) 5, 054507



ρ meson spectra on LQCD -- Operator

$[000]T_1^-$	$[001]A_1$	$[001]E_2$	$[011]A_1$
$\rho_{[000]}$	$\rho_{[001]}$	$\rho_{[001]}$	$\rho_{[011]}$
$\pi_{[001]}\pi_{[00-1]}$	$\pi_{[000]}\pi_{[001]}$	$\pi_{[0-10]}\pi_{[011]}$	$\pi_{[000]}\pi_{[011]}$
$\pi_{[011]}\pi_{[0-1-1]}$	$\pi_{[0-10]}\pi_{[011]}$	$\pi_{[0-1-1]}\pi_{[111]}$	$\pi_{[-100]}\pi_{[111]}$
$\pi_{[111]}\pi_{[-1-1-1]}$	$\pi_{[-1-10]}\pi_{[111]}$		$\pi_{[01-1]}\pi_{[002]}$
	$\pi_{[00-1]}\pi_{[002]}$		
$[011]B_1$	$[011]B_2$	$[111]A_1$	$[002]A_1$
$\rho_{[011]}$	$\rho_{[011]}$	$\rho_{[111]}$	$\rho_{[002]}$
$\pi_{[010]}\pi_{[001]}$	$\pi_{[-100]}\pi_{[111]}$	$\pi_{[000]}\pi_{[111]}$	$\pi_{[000]}\pi_{[002]}$
$\pi_{[110]}\pi_{[-101]}$	$\pi_{[110]}\pi_{[-101]}$	$\pi_{[100]}\pi_{[011]}$	
$\pi_{[0-11]}\pi_{[002]}$		$\pi_{[200]}\pi_{[-111]}$	

$$\rho^0 = \frac{1}{\sqrt{2}}(\bar{u}\Gamma^P u - \bar{d}\Gamma^P d)$$

$$\rho^{0\dagger} = \frac{1}{\sqrt{2}}(\bar{u}\bar{\Gamma}^P u - \bar{d}\bar{\Gamma}^P d)$$

$$\pi^+ = \bar{d}\Gamma^k u, \quad \pi^- = \bar{u}\Gamma^k d$$

$$\begin{aligned} \pi\pi &= \frac{1}{\sqrt{2}}(\pi^+(k_1)\pi^-(k_2) - \pi^-(k_1)\pi^+(k_2)) \\ &= \frac{1}{\sqrt{2}}(\bar{d}\Gamma^{k_1} u \bar{u}\Gamma^{k_2} d - \bar{u}\Gamma^{k_1} d \bar{d}\Gamma^{k_2} u) \end{aligned}$$

$$(\pi\pi)^\dagger = \frac{1}{\sqrt{2}}(\bar{d}\bar{\Gamma}^{k_2} u \bar{u}\bar{\Gamma}^{k_1} d - \bar{u}\bar{\Gamma}^{k_2} d \bar{d}\bar{\Gamma}^{k_1} u)$$



ρ meson spectra on LQCD -- correlation function

$$C_{ij} = \langle 0 | \mathcal{O}_i \mathcal{O}_j^\dagger | 0 \rangle = \sum_n Z_i^n Z_j^{n*} e^{-E_n t}$$

Computed based on wick contraction and quark propagator calculated based on the configuration.

To parameterize the correlation function !

Energy level in the finite volume

- Solve the generalized eigenvalue problem (**GEVP**)

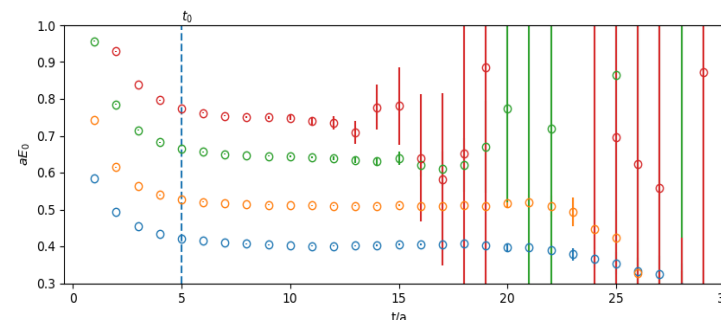
$$\mathcal{G}_{ij}(t_0 + dt) u_j^\alpha = e^{-m_a dt} \mathcal{G}_{ij}(t_0) u_j^\alpha,$$

$$\mathcal{G}^\alpha(t) = v_i^\alpha \mathcal{G}_{ij}(t) u_j^\alpha,$$

$$[\mathcal{G}^{-1}(t_0) \mathcal{G}(t_0 + dt)]_{ij} u_j^\alpha = c^\alpha u_i^\alpha$$

$$E^\alpha(t) = \frac{1}{n} \log \frac{\mathcal{G}^\alpha(t)}{\mathcal{G}^\alpha(t+n)},$$

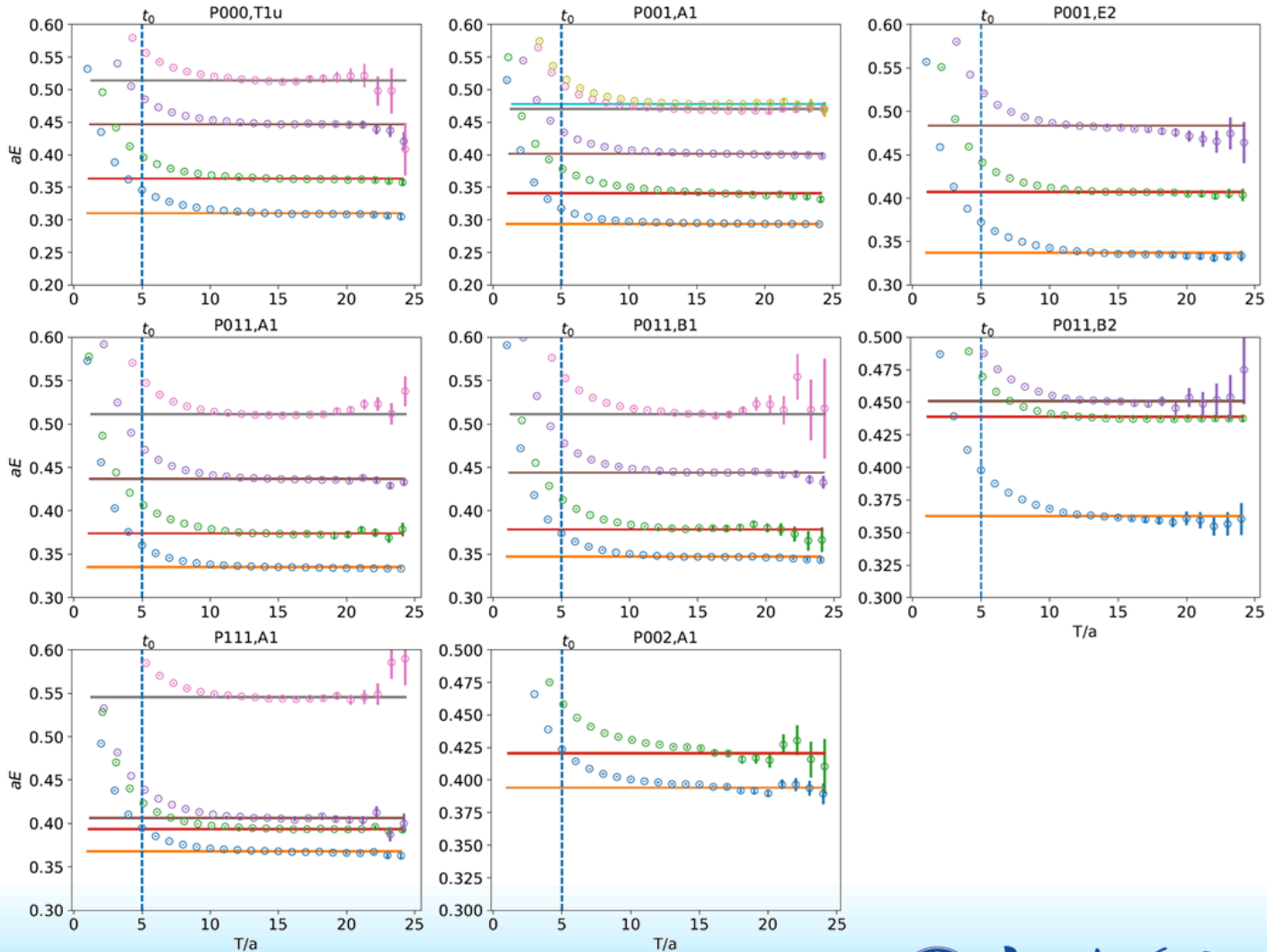
$$v_i^\alpha [\mathcal{G}(t_0 + dt) \mathcal{G}^{-1}(t_0)]_{ij} = c^\alpha v_j^\alpha,$$



Kiratidis, et al., PRD 91, 094509 (2015)



ρ meson spectra on LQCD -- Spectra



Configuration: F48P30

There are 8 different spectra with different total momentum and irreps, including following irreducible representations.:

$$p^2 = 0 : T_1^-,$$

$$p^2 = 1 : A_1 \text{ and } E_2,$$

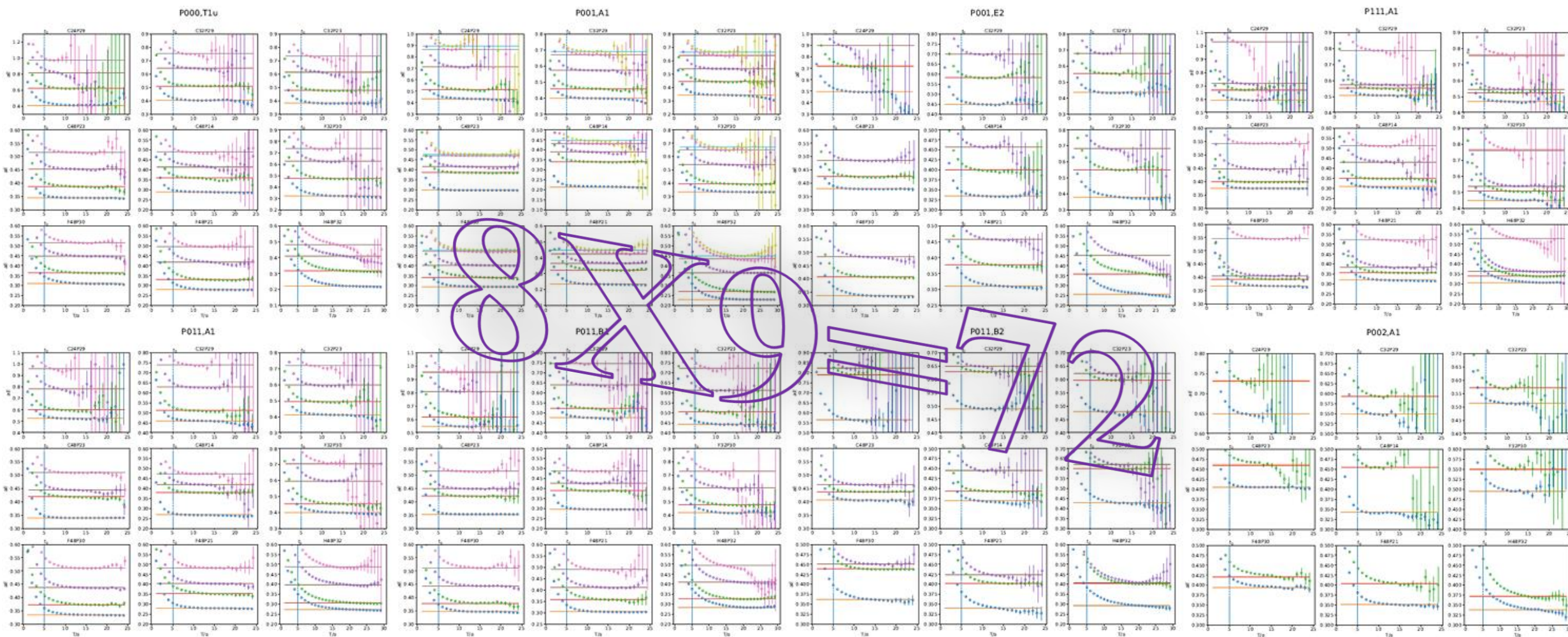
$$p^2 = 2 : A_1, B_1 \text{ and } B_2,$$

$$p^2 = 3 : A_1,$$

$$p^2 = 4 : A_1.$$

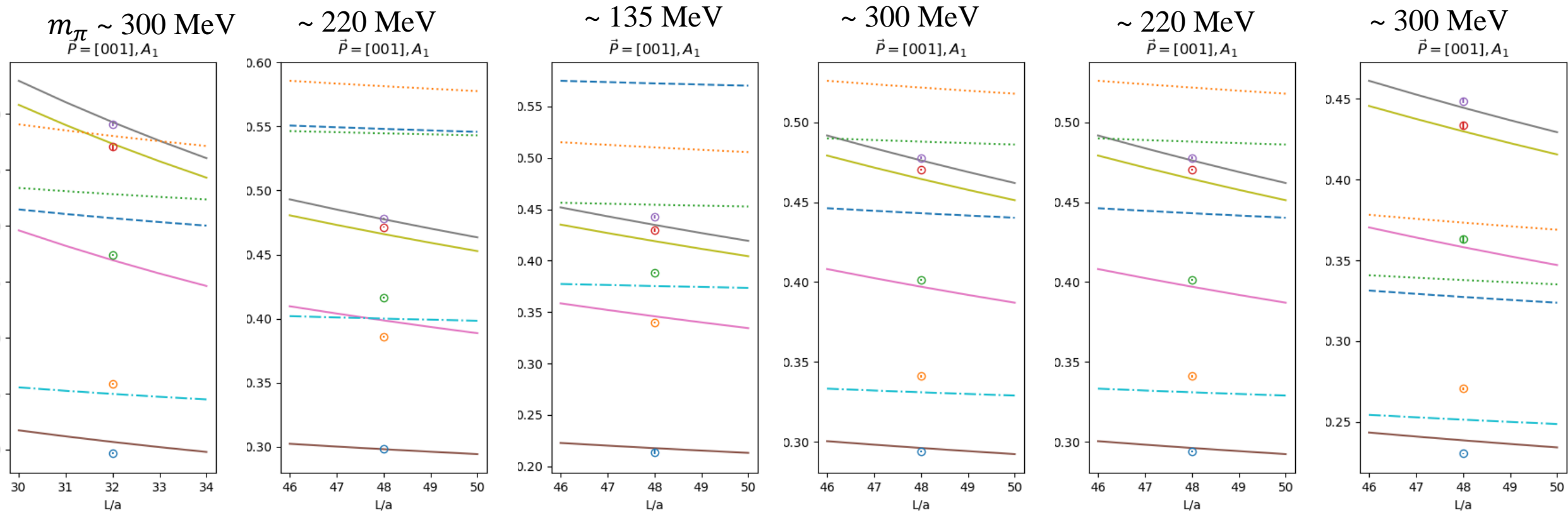


ρ meson spectra on LQCD -- Spectra



ρ meson spectra on LQCD -- Spectra

C32P29 C48P23 C48P14 F48P30 F48P21 H48P32



$a \sim 0.105 \text{ fm}$

$a \sim 0.077 \text{ fm}$

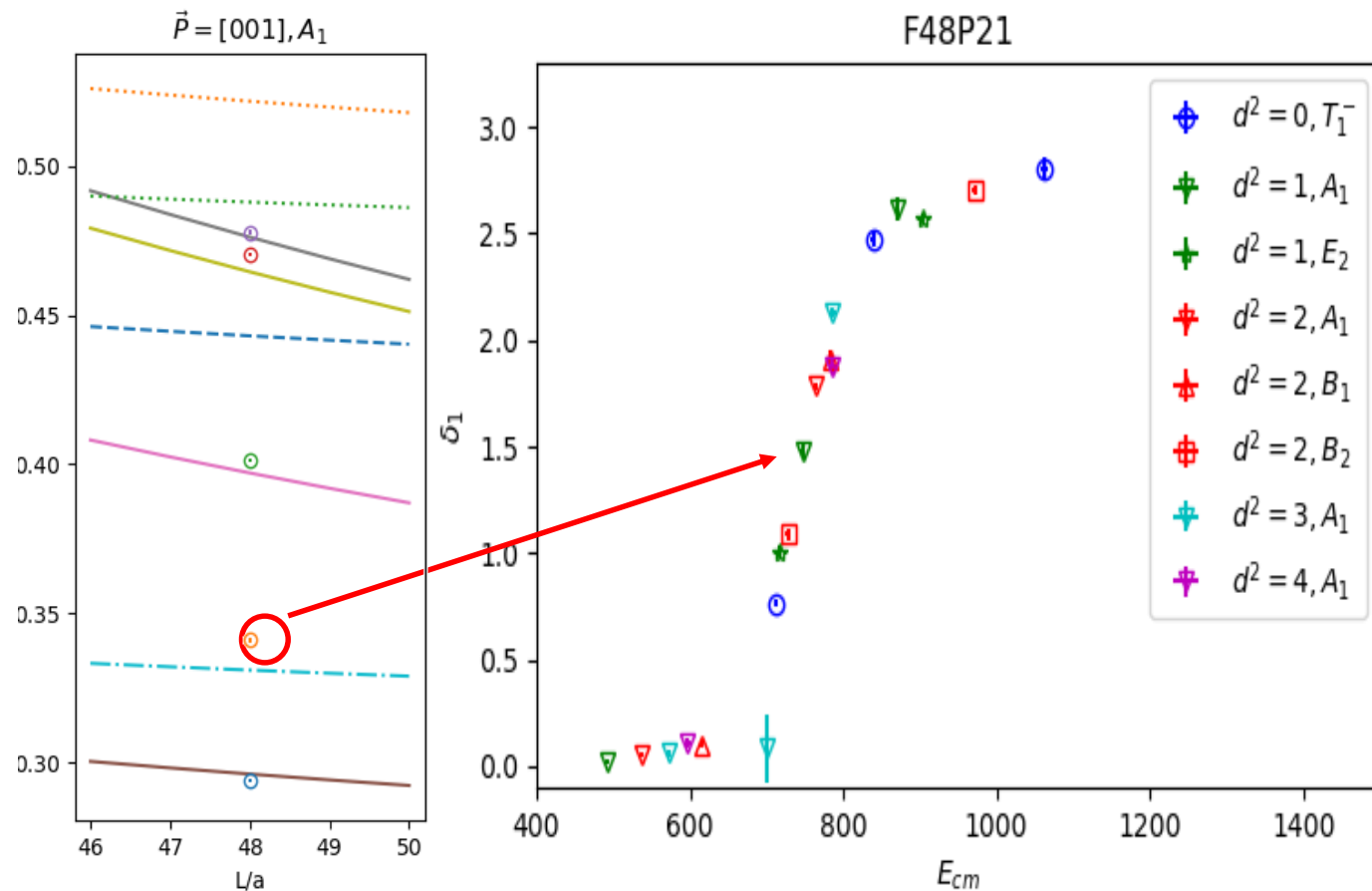
$a \sim 0.052 \text{ fm}$

ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

d^2	Λ, μ	$M_{11,11}^{(\vec{P}, \Lambda, \mu)}$
0	$T_1^-, 2$	$w_{0,0}$
1	$A_1, 1$	$w_{0,0} + 2w_{2,0}$
1	$E_2, 1$	$w_{0,0} - w_{2,0}$
1	$E_2, 2$	$w_{0,0} - w_{2,0}$
2	$A_1, 1$	$w_{0,0} + \frac{1}{2}w_{2,0} + i\sqrt{6}w_{2,1} - \frac{\sqrt{6}}{2}w_{2,2}$
2	$B_1, 1$	$w_{0,0} + \frac{1}{2}w_{2,0} - i\sqrt{6}w_{2,1} - \frac{\sqrt{6}}{2}w_{2,2}$
2	$B_2, 1$	$w_{0,0} - w_{2,0} + \sqrt{6}w_{2,2}$
3	$A_1, 1$	$w_{0,0} - i2\sqrt{6}w_{2,2}$

$$\det(M_{ln, \nu n'}(k) - \delta_{ll'} \delta_{nn'} \cot(\delta_l)) = 0, \quad \delta_1 = \text{arccot} M_{11,11}^{(\vec{P}, \Lambda, \mu)},$$

$$w_{j,s} = \frac{\mathcal{Z}_{js}(1, q^2)}{\pi^{3/2} \sqrt{2j+1} \gamma q^{j+1}}, \quad q = \frac{kL}{2\pi}, \quad \delta_{l \geq 3} \text{ are negligible}$$

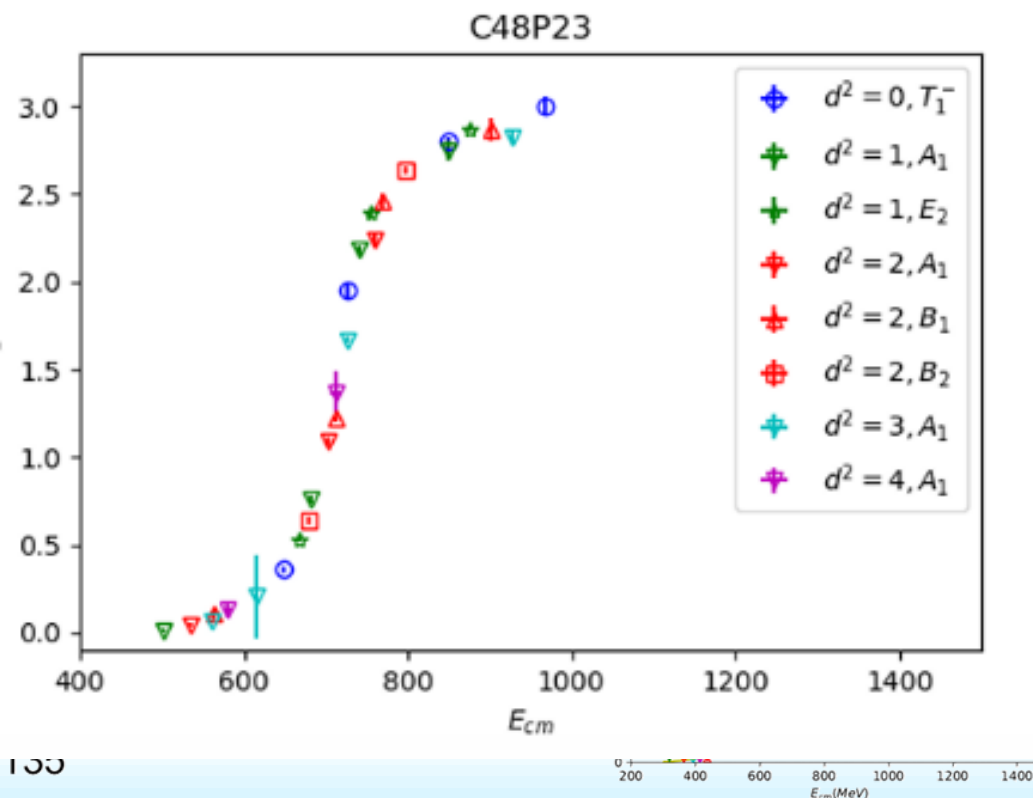
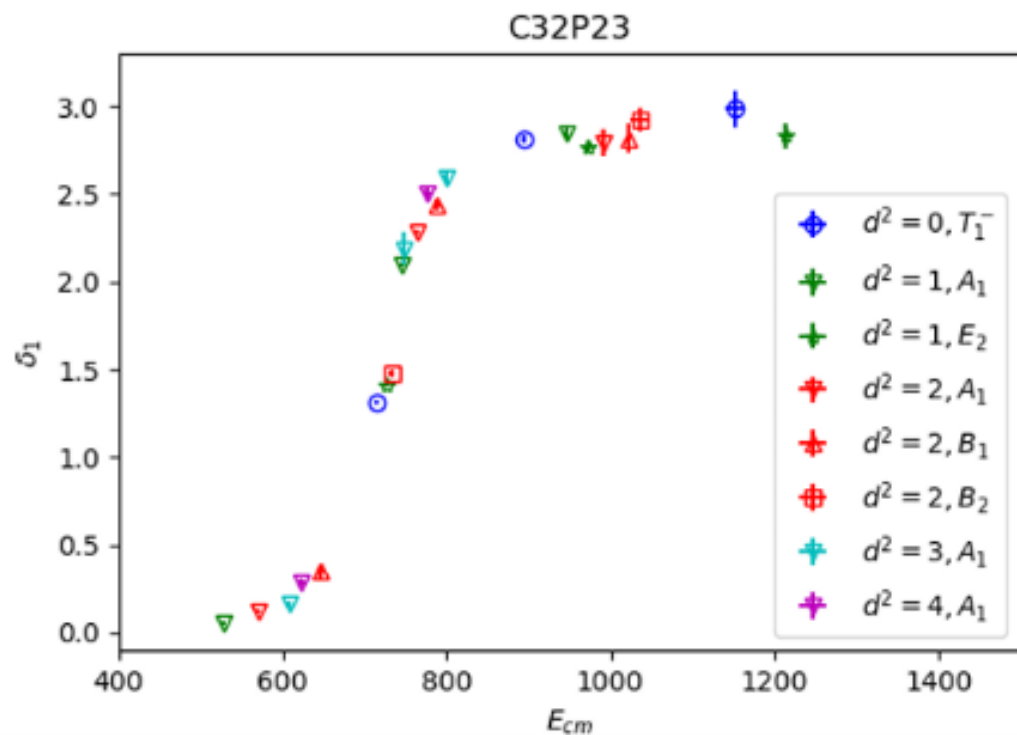


$$E_{cm} = \sqrt{E_{001}^2 - 4\pi^2/L^2} = \sqrt{(0.34/a)^2 - 4\pi^2/(48a)^2} = 771 \text{ MeV}$$

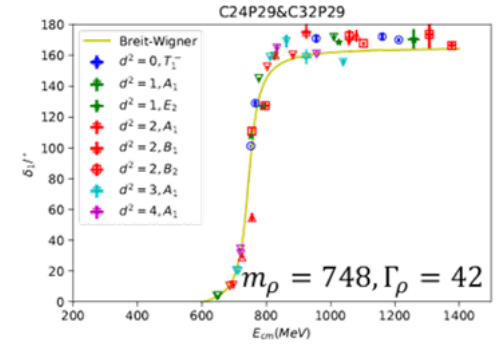
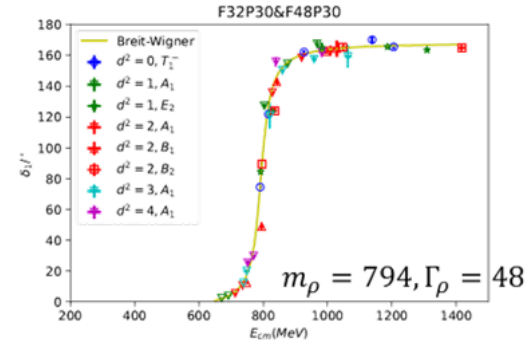
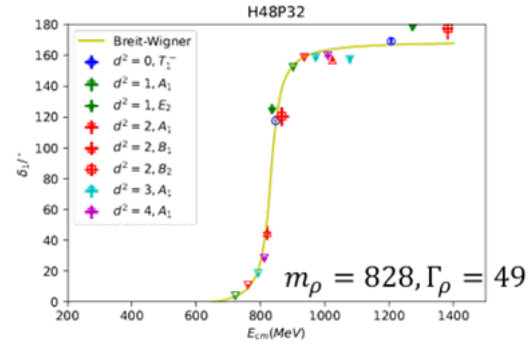


ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

Name	Volume	Spacing	β	m_π/MeV	$m_\pi L$
C32P23	$32^3 \times 64$	0.10530fm	6.20	228	3.91
C48P23	$48^3 \times 96$			225	5.79

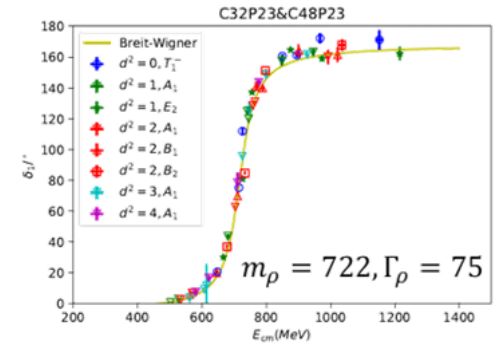
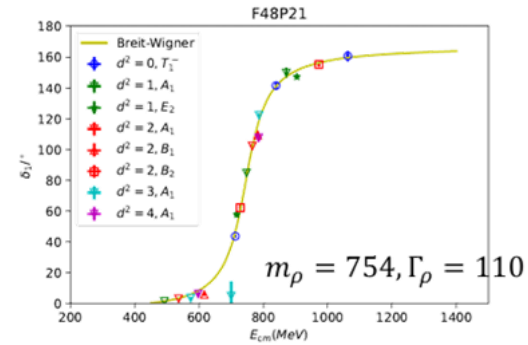


ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

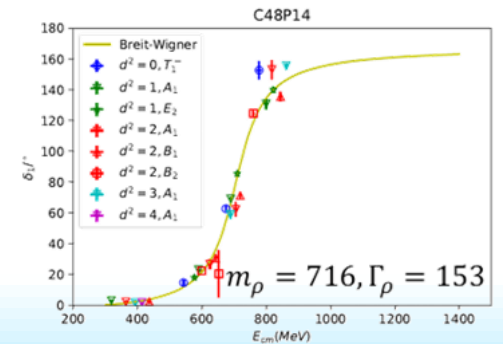
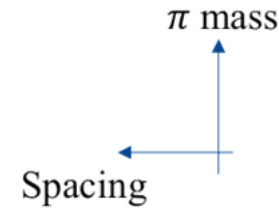
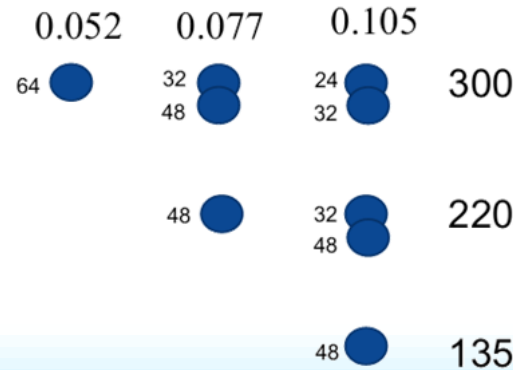


$$a_1 = \frac{-\sqrt{s}\Gamma(s)}{s - m_\rho^2 + i\sqrt{s}\Gamma(s)} = e^{i\delta(s)} \sin \delta(s),$$

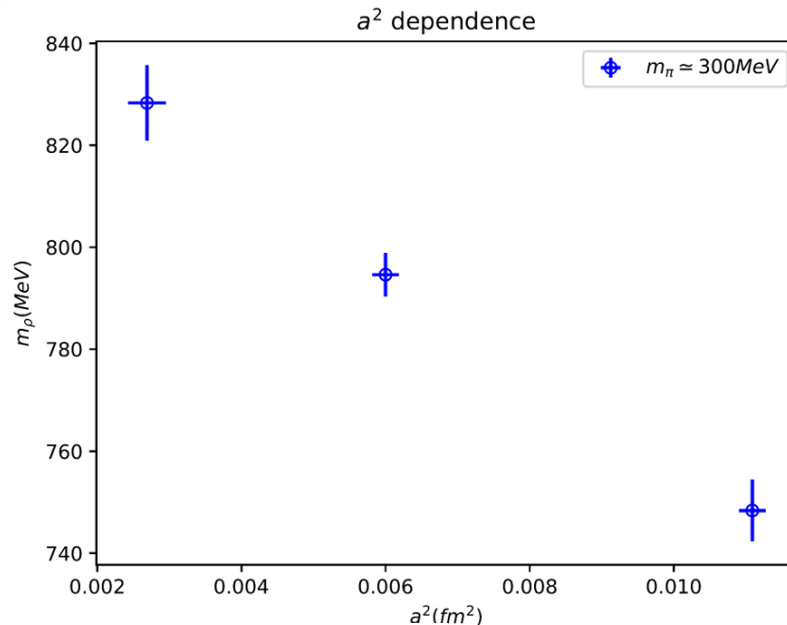
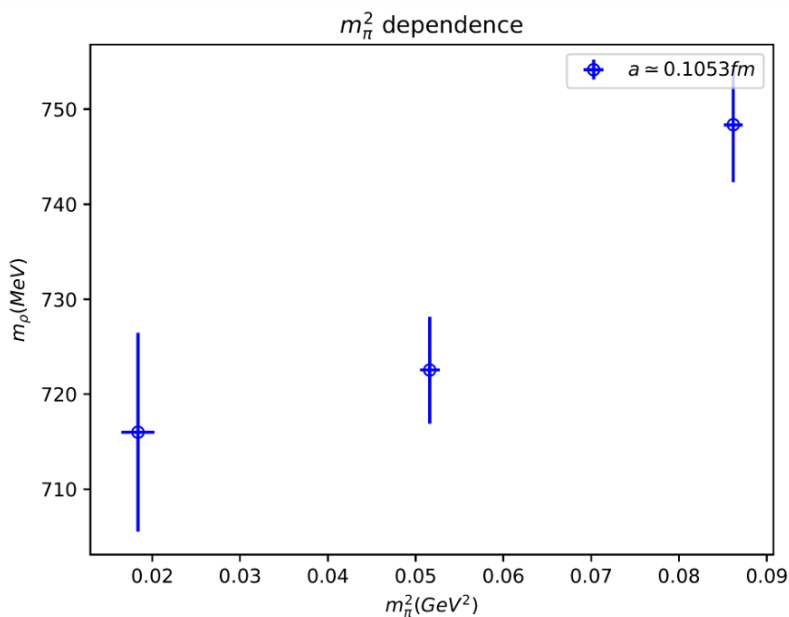
$$\Gamma(s) = \frac{p^{*3} g_{\rho\pi\pi}^2}{s 6\pi}.$$



1. The lineshape of the phase shift from large lattice size are even more smooth since $m_\pi L$ is larger
2. There exists the jump of phase shift from 0 degree to 180 degrees



ρ meson spectra on LQCD -- M & Γ of ρ meson



0.052	0.077	0.105	
828/49	794/48	748/42	300
	754/110	722/75	220
		716/153	135

$$m_\rho = c_0 + c_1 m_\pi^2 + c_2 a^2$$

$$c_1 = 0.7611(85) \text{ GeV},$$

$$c_2 = 0.807(81) \text{ GeV}^{-1},$$

$$c_3 = -7.13(79) \text{ GeV} \cdot \text{fm}^{-2},$$

$$\chi^2/d.o.f = 0.76$$

$$g = \tilde{c}_0 + \tilde{c}_1 m_\pi^2 + \tilde{c}_2 a^2$$

$$\tilde{c}_1 = 6.37(11),$$

$$\tilde{c}_2 = -3.35(87) \text{ GeV}^{-2},$$

$$\tilde{c}_3 = -16.8(6.9) \text{ fm}^{-2},$$

$$\chi^2/d.o.f = 0.26$$

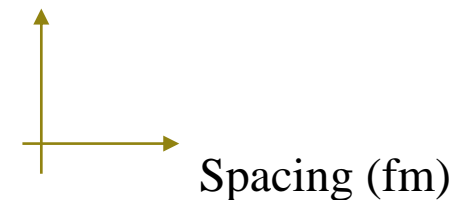
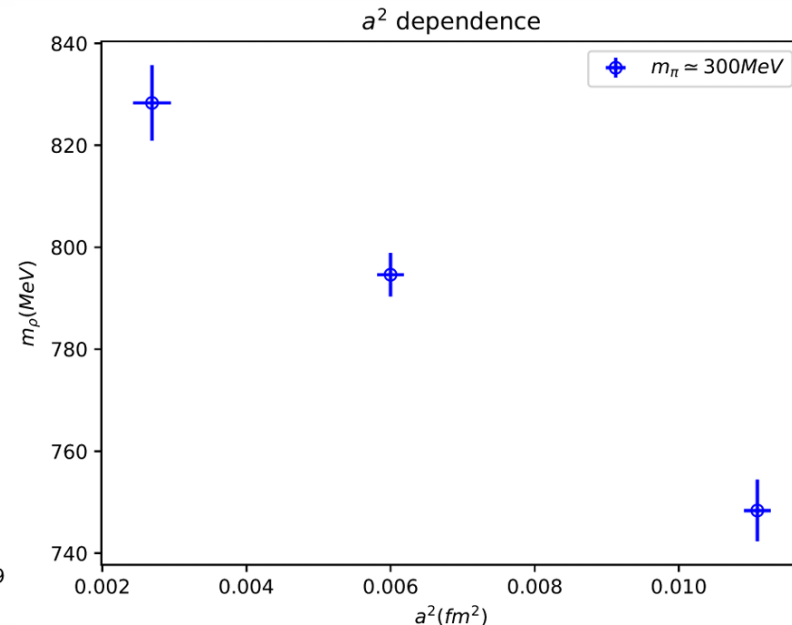
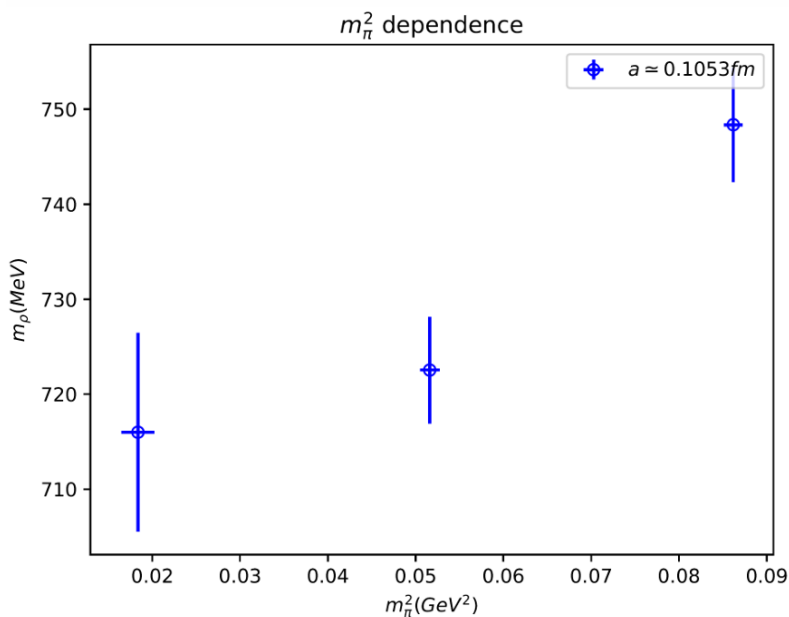


Table 4: Fit results for m_ρ , $g_{\rho\pi\pi}$ and Z_{pole}

	C24P29 & C32P29	C32P23 & C48P23	C48P14	F32P30 & F48P30	F48P21	H48P32
m_ρ	748.38(6.0)	722.55(5.6)	716.0(10.46)	794.62(4.24)	751.3(2.98)	828.3(7.4)
$g_{\rho\pi\pi}$	5.249(79)	5.32(14)	6.04(44)	6.09(58)	6.80(38)	6.44(45)
Z_{pole}	744.94(6.0) - 19.89(1.5)i	715.23(5.6) - 38.47(2.8)i	703.89(10.4) - 70.77(8.5)i	789.65(4.3) - 24.71(1.0)i	743.53(3.0) - 52.46(1.7)i	824.51(7.4) - 22.93(0.9)i



ρ meson spectra on LQCD -- M & Γ of ρ meson



$$m_\rho = c_0 + c_1 m_\pi^2 + c_2 a^2 \quad \text{Here we take } m_{\pi\text{-phy}} = m_{\pi^0} \sim 135 \text{ MeV}$$

$$m_\rho = 775.8(8.5) \text{ MeV}, \quad \Gamma_\rho(m_\rho) = 160(10) \text{ MeV},$$

$$Z_{\text{pole}} = 766.7(8.6) - i76.5(4.4) \text{ MeV}.$$

$$\rho(770) \text{ T-MATRIX POLE } \sqrt{s} \quad (761 - 765) - i(71 - 74) \text{ MeV}$$

$\rho(770)$ MASS

$$\text{NEUTRAL ONLY, } e^+e^- \quad 775.26 \pm 0.23 \text{ MeV}$$

$$\text{CHARGED ONLY, } \tau \text{ DECAYS and } e^+e^- \quad 775.11 \pm 0.34 \text{ MeV}$$

$$\text{MIXED CHARGES, OTHER REACTIONS} \quad 763.0 \pm 1.2 \text{ MeV}$$

$$\text{CHARGED ONLY, HADROPRODUCED} \quad 766.5 \pm 1.1 \text{ MeV}$$

$$\text{NEUTRAL ONLY, PHOTOPRODUCED} \quad 769.2 \pm 0.9 \text{ MeV}$$

$$\text{NEUTRAL ONLY, OTHER REACTIONS} \quad 769.0 \pm 0.9 \text{ MeV (S = 1.4)}$$

$$m_{\rho(770)^0} - m_{\rho(770)^\pm} \quad -0.7 \pm 0.8 \text{ MeV (S = 1.5)}$$

$$m_{\rho(770)^+} - m_{\rho(770)^-}$$

$$\rho(770) \text{ RANGE PARAMETER} \quad 5.3^{+0.9}_{-0.7} \text{ GeV}^{-1}$$

$\rho(770)$ WIDTH

$$\text{NEUTRAL ONLY, } e^+e^- \quad 147.4 \pm 0.8 \text{ MeV (S = 2.0)}$$

$$\text{CHARGED ONLY, } \tau \text{ DECAYS and } e^+e^- \quad 149.1 \pm 0.8 \text{ MeV}$$

$$\text{MIXED CHARGES, OTHER REACTIONS} \quad 149.5 \pm 1.3 \text{ MeV}$$

$$\text{CHARGED ONLY, HADROPRODUCED} \quad 150.2 \pm 2.4 \text{ MeV}$$

$$\text{NEUTRAL ONLY, PHOTOPRODUCED} \quad 151.5^{+1.9}_{-2.1} \text{ MeV}$$

$$\text{NEUTRAL ONLY, OTHER REACTIONS} \quad 150.9 \pm 1.7 \text{ MeV (S = 1.1)}$$

$$\Gamma_{\rho(770)^0} - \Gamma_{\rho(770)^\pm} \quad 0.3 \pm 1.3 \text{ MeV (S = 1.4)}$$

$$\Gamma_{\rho(770)^+} - \Gamma_{\rho(770)^-} \quad 1.8 \pm 2.1$$



ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

$$T_{\pi\pi \rightarrow \pi\pi}^{l=1}(z) = \frac{|V_{\rho\pi\pi}(\bar{k}(z))|^2}{z - m_\rho^B - \Sigma(z)},$$

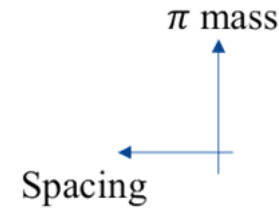
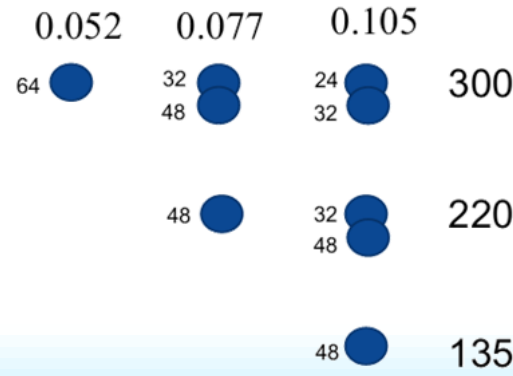
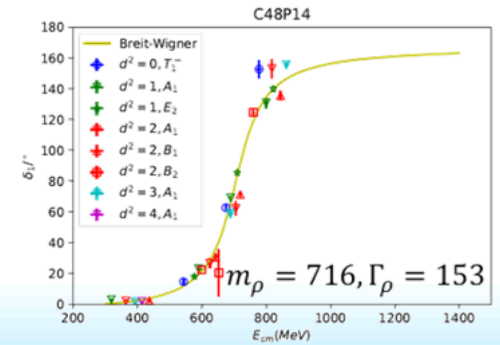
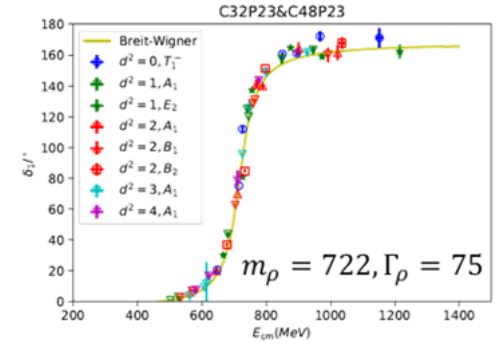
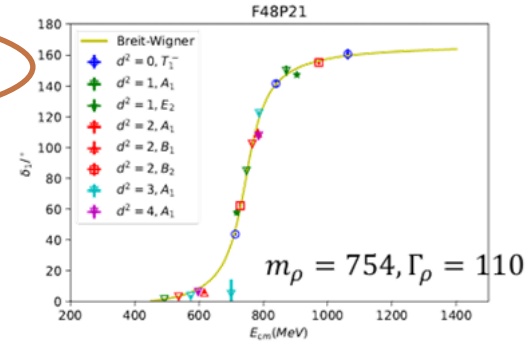
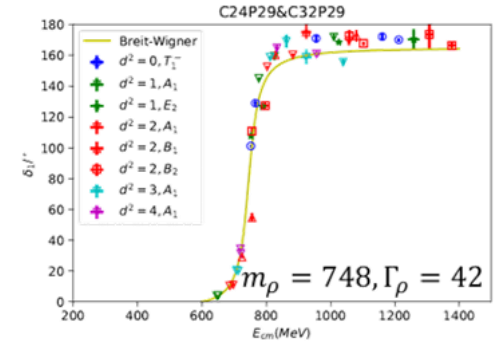
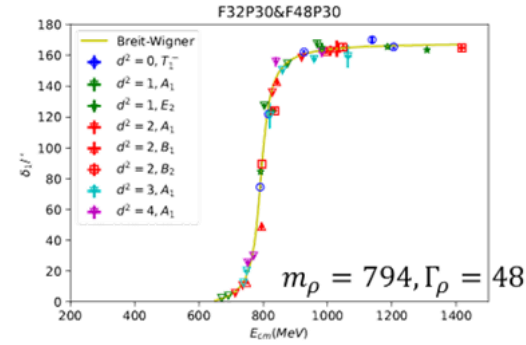
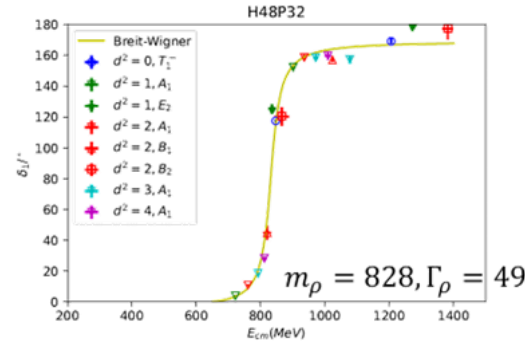
$$\Sigma(z) = \int q^2 dq \frac{|V_{\rho\pi\pi}(q)|^2}{z - 2\sqrt{q^2 + m_\pi^2} + i\epsilon}.$$

$$a_1 = \frac{-\sqrt{s}\Gamma(s)}{s - m_\rho^2 + i\sqrt{s}\Gamma(s)} = e^{i\delta(s)} \sin \delta(s),$$

$$\Gamma(s) = \frac{p^{*3} g_{\rho\pi\pi}^2}{s 6\pi}.$$

HEFT

1. The lineshape of the phase shift from large lattice size are even more smooth since $m_\pi L$ is larger
2. There exists the jump of phase shift from 0 degree to 180 degrees



Introduction of HEFT

J. M. M. Hall *etc.* PRD 87(2013), 094510
 J.-j. Wu *etc.* PRC90 (2014), 055206
 Y. Li *etc.* PRD 101(2020), 114501
 PRD 103(2021), 094518

$$H = H_0 + H_I$$

$$H_0 = \sum_{i=1,n} |B_i\rangle m_i \langle B_i| + \sum_{\alpha} |\alpha(k_{\alpha})\rangle \left[\sqrt{m_{\alpha 1}^2 + k_{\alpha}^2} + \sqrt{m_{\alpha 2}^2 + k_{\alpha}^2} \right] \langle \alpha(k_{\alpha})|$$

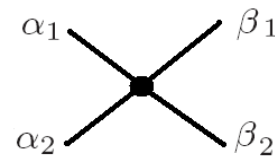
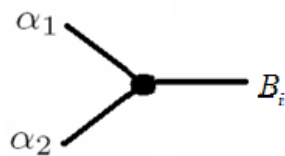
$|B_i\rangle$ bare state, bare mass m_i

$|\alpha(k_{\alpha})\rangle$ non-interaction channels

$$H_I = \hat{g} + \hat{v}$$

$$\hat{g} = \sum_{\alpha} \sum_{i=1,n} \left[|\alpha(k_{\alpha})\rangle g_{i,\alpha}^+ \langle B_i| + |B_i\rangle g_{i,\alpha} \langle \alpha(k_{\alpha})| \right]$$

$$\hat{v} = \sum_{\alpha,\beta} |\alpha(k_{\alpha})\rangle v_{\alpha,\beta} \langle \beta(k_{\beta})|$$



T matrix
 (Phase Shifts,
 inelasticity)

Resonance
 (Mass, Width, Pole position, Coupling)

HEFT

**Lattice
 Spectrum**

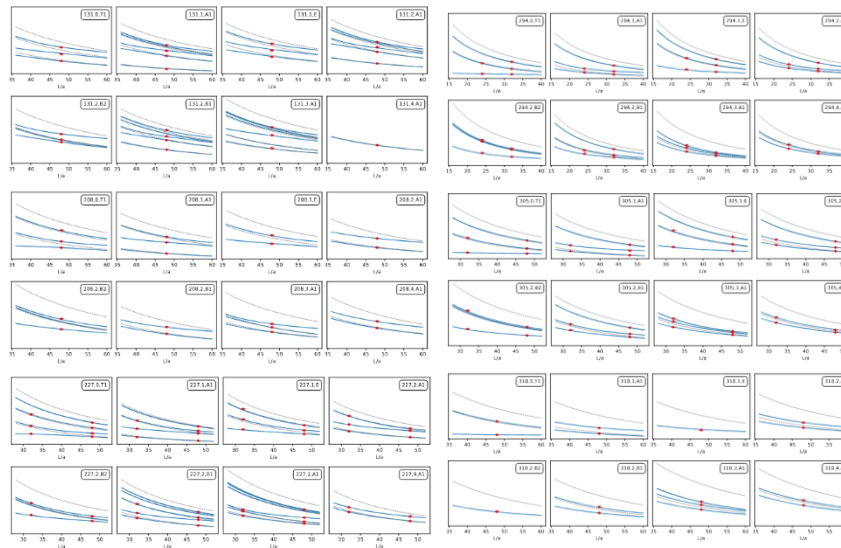


ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

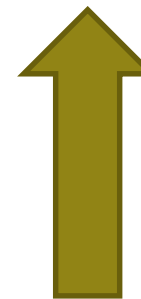
$$T_{\pi\pi \rightarrow \pi\pi}^{l=1}(z) = \frac{|V_{\rho\pi\pi}(\bar{k}(z))|^2}{z - m_\rho^B - \Sigma(z)},$$

$$\Sigma(z) = \int q^2 dq \frac{|V_{\rho\pi\pi}(q)|^2}{z - 2\sqrt{q^2 + m_\pi^2} + i\epsilon}.$$

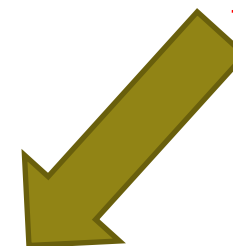
$$V_{\rho\pi\pi}(k) = \frac{g_{\rho\pi\pi}k}{\sqrt{m_\rho^B} \sqrt{k^2 + m_\pi^2}} \left(\frac{\Lambda_{\rho\pi\pi}^2}{k^2 + \Lambda_{\rho\pi\pi}^2} \right)^2$$



Resonance
(Mass, Width, Pole position, Coupling)



HEFT



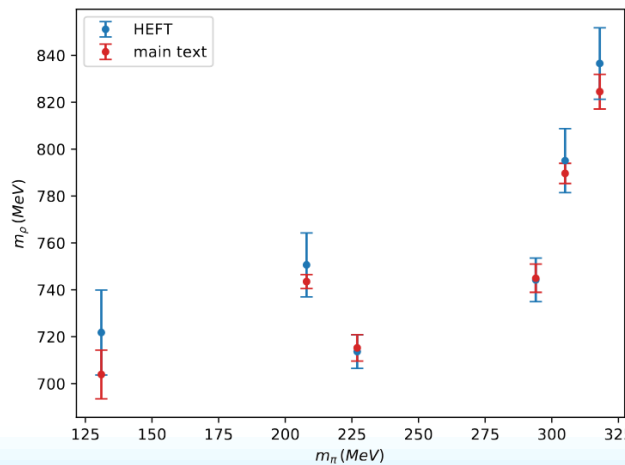
T matrix
(Phase Shifts,
inelasticity)



**Lattice
Spectrum**

- Three free parameters, $m_\rho^B, g_{\rho\pi\pi}, \Lambda_{\rho\pi\pi}$
- Here we take $g_{\rho\pi\pi}, \Lambda_{\rho\pi\pi}$ as constant for different m_π and a .
- But m_ρ^B is pion mass dependence.

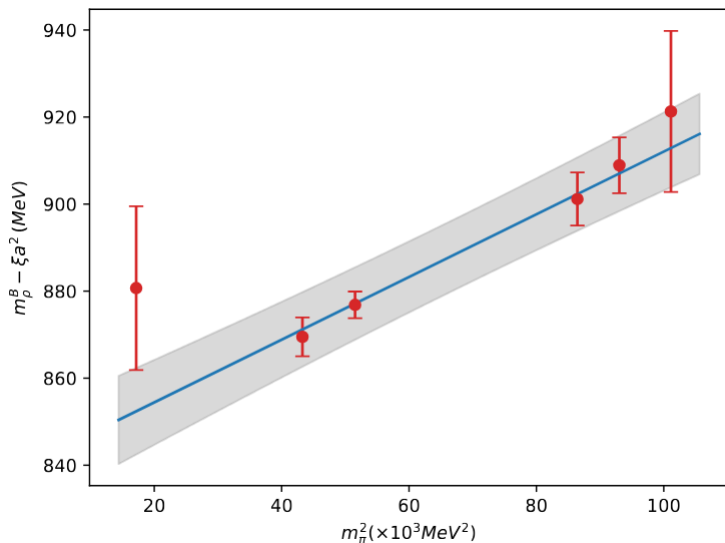
$$m_\rho^B(m_\pi, a) = c_0 + c_1 m_\pi^2 + \xi a^2,$$



ρ meson spectra on LQCD -- Phase shift of $\pi\pi$

- Three free parameters, $m_\rho^B, g_{\rho\pi\pi}, \Lambda_{\rho\pi\pi}$
- Here we take $g_{\rho\pi\pi}, \Lambda_{\rho\pi\pi}$ as constant for different m_π and a .
- But m_ρ^B is pion mass dependence.

$$m_\rho^B(m_\pi, a) = c_0 + c_1 m_\pi^2 + \xi a^2,$$



$$(c_0, c_1, \xi) = (840, 0.72 \times 10^{-3}, -7912),$$

$$\text{cov}(c_0, c_1, \xi) = \begin{pmatrix} 125 & -0.88 \times 10^{-3} & -7474 \\ & 0.013 \times 10^{-6} & 11.7 \times 10^{-3} \\ & & 7.39 \times 10^5 \end{pmatrix},$$

$$\begin{aligned} c_1 &= 0.7611(85) \text{ GeV}, \\ c_2 &= 0.807(81) \text{ GeV}^{-1}, \\ c_3 &= -7.13(79) \text{ GeV} \cdot \text{fm}^{-2}, \end{aligned}$$

$$m_\rho^B(m_\pi^{\text{phy}}, 0) = 853 \pm 10 \text{ MeV}$$

$$m_\rho^{\text{pole, ext}} = (792 \pm 11 \pm 4) - i(64 \pm 4 \pm 4),$$

$$Z_{\text{pole}} = 766.7(8.6) - i76.5(4.4) \text{ MeV}.$$

$\rho(770)$ T-MATRIX POLE \sqrt{s}

$(761 - 765) - i(71 - 74) \text{ MeV}$



Summary

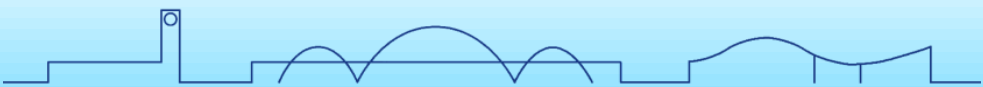
We using CLQCD configuration measure the rho meson mass and width at the physical pion mass and continuum limit with the highest precision up to now!

$$m_\rho = 775.8(8.5)\text{MeV}, \quad \Gamma_\rho(m_\rho) = 160(10)\text{MeV},$$

$$Z_{\text{pole}} = 766.7(8.6) - i76.5(4.4)\text{MeV}.$$



Thanks for attention!



中国科学院大学
University of Chinese Academy of Sciences

