

Testing Quantum Gravity by Entanglements and Decoherence

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Plan:

Quantum gravity

- Approaches
- From quantum information

Are cosmological perturbations quantum?

- Generation mechanism
- Decoherence (new: with boundary term)

Observational tests?

- Bell inequalities?
- Quantum noise?
- Null tests?

What's quantum gravity?

Claims from strongest to weakest:

- What is *the* quantum gravity theory?
 - Not only QG but also *new physics*
 - Strings: what is a/the quantum gravity theory?
 - Other possible unifications of QM & GR?
- Check if gravity is quantized (known physics)
 - Seeking for gravitons
- Check gravity has some quantum natures (known physics)
 - Superposition, entanglements & decoherence

The quantum information regime of quantum gravity:

- Can matter superposition be affected by gravity?
- Can matter entanglement be caused by gravity?
- Can gravitational superposition affect matter?
- Other ideas?

The quantum information regime of quantum gravity:

- Can matter superposition be affected by gravity?

Yes (COW 1975), but does not really quantum gravity

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PHYSICAL REVIEW LETTERS

9 JUNE 1975

Observation of Gravitationally Induced Quantum Interference*

R. Colella and A. W. Overhauser

Department of Physics, Purdue University, West Lafayette, Indiana 47907

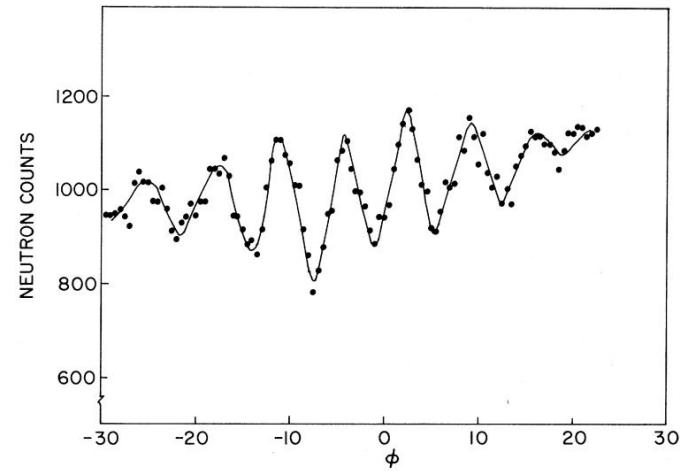
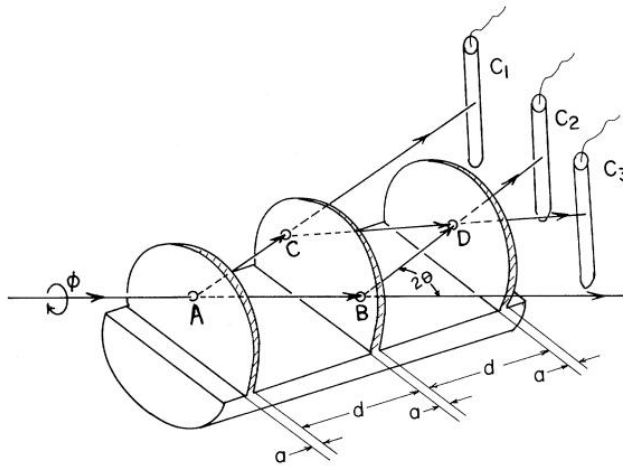
and

S. A. Werner

Scientific Research Staff, Ford Motor Company, Dearborn, Michigan 48121

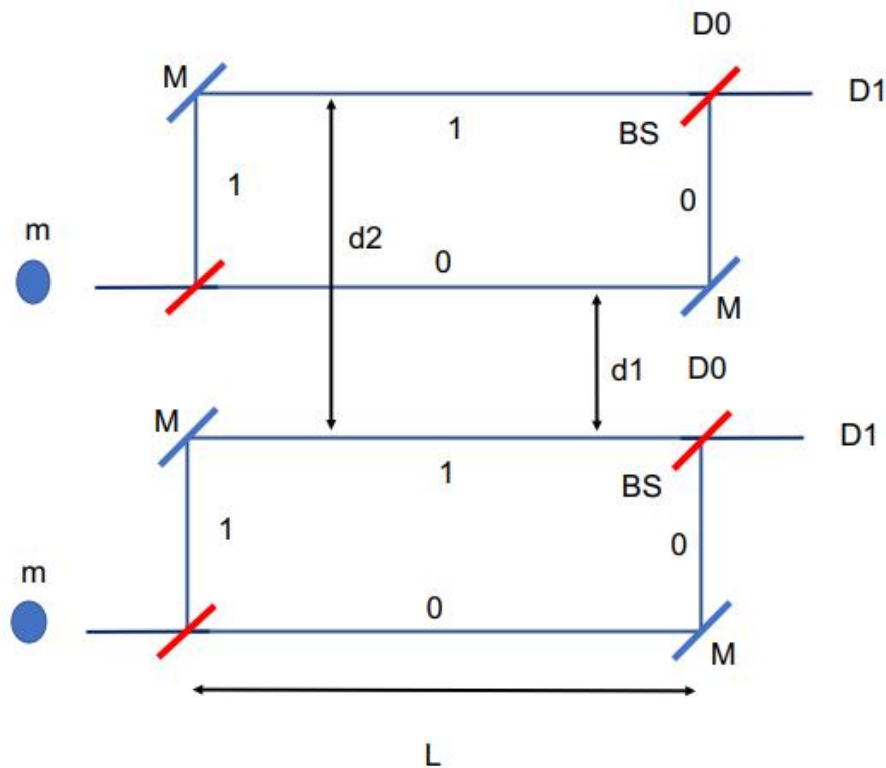
(Received 14 April 1975)

We have used a neutron interferometer to observe the quantum-mechanical phase shift of neutrons caused by their interaction with Earth's gravitational field.



The quantum information regime of quantum gravity:

- Can matter superposition be affected by gravity?
- Can matter entanglement be caused by gravity?



Bose et al (2017), Marletto & Vedral (2017)

Matter entanglement \Rightarrow quantum gravity:

Qubit 1 --- creation of entanglement with C --- Qubit 2

Then C cannot have only one observable (classical)

Proof: If initially state separable: $|Q_1(0)\rangle|Q_2(0)\rangle$ Marletto & Vedral 2017

and if Hamiltonian H_1 acts on $|Q_1(0)\rangle$, H_2 acts on $|Q_2(0)\rangle$,

gravity is reflected by time-dependent coefficients in H_1 & H_2

Then the final state $|Q_1(t)\rangle|Q_2(t)\rangle$ is also separable

Challenge: is the
constraint part "local"?

More generally:

Separated systems cannot be entangled by

Local Operations and Classical Communication (LOCC)

See, e.g. R. Horodecki, P. Horodecki, M. Horodecki, K. Horodecki

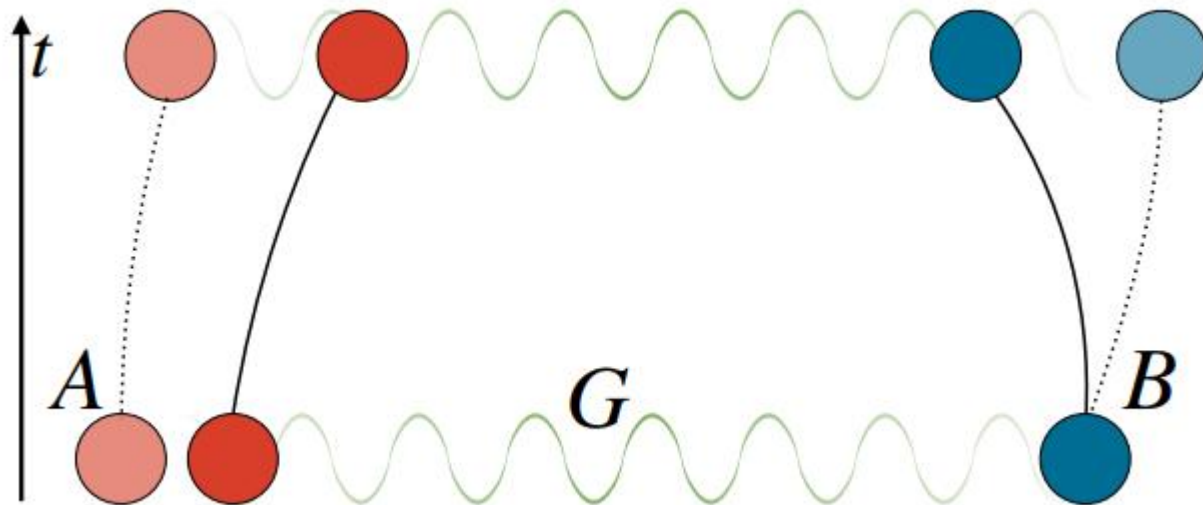
Generalizing LOCC to Generalized Probabilistic Theories (GPT)

A no-go theorem on the nature of the gravitational field beyond quantum theory

Galley, Giacomini, Selby (2020)

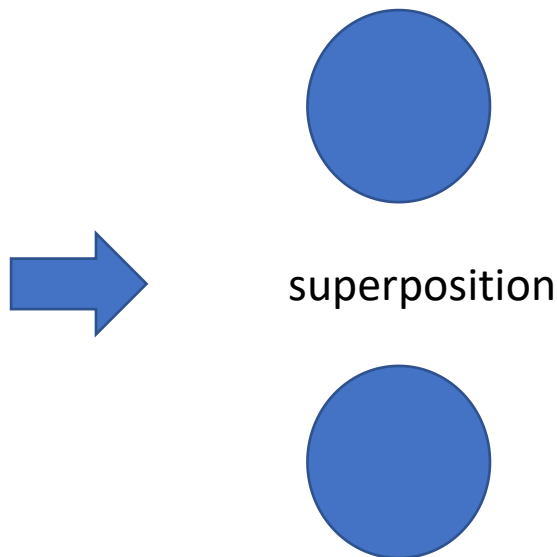
Theorem III.1. We consider two non-classical systems A and B, initially in a separable state, and some unknown field G. If, after some time t , entanglement between the systems A and B is observed, then the following statements are incompatible:

1. Subsystem independence of A and B;
2. A and B interact locally via the mediator G;
3. G is classical.



The quantum information regime of quantum gravity:

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Matter:

- Attracted by one (superposition)? $G_{\mu\nu} \propto T_{\mu\nu}$
- Attracted by the average? $G_{\mu\nu} \propto \langle T_{\mu\nu} \rangle$

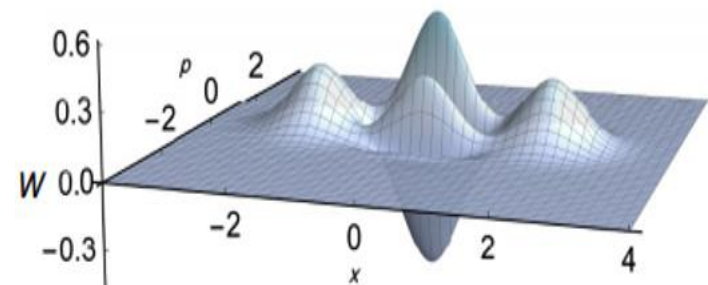
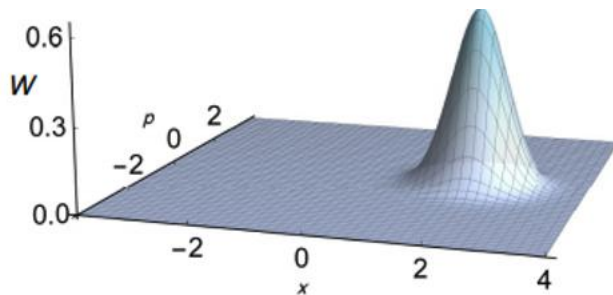
(Schrodinger-Newton semi-classical gravity)

The quantum information regime of quantum gravity:

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- Other ideas?

Non-Gaussianity as a Signature of a Quantum Theory of Gravity

Richard Howl, Vlatko Vedral, Devang Naik, Marios Christodoulou, Carlo Rovelli, and Aditya Iyer
PRX Quantum **2**, 010325 – Published 17 February 2021



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- Decoherence (detailed)

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Coherent, squeezed and thermal states

- Coherent:

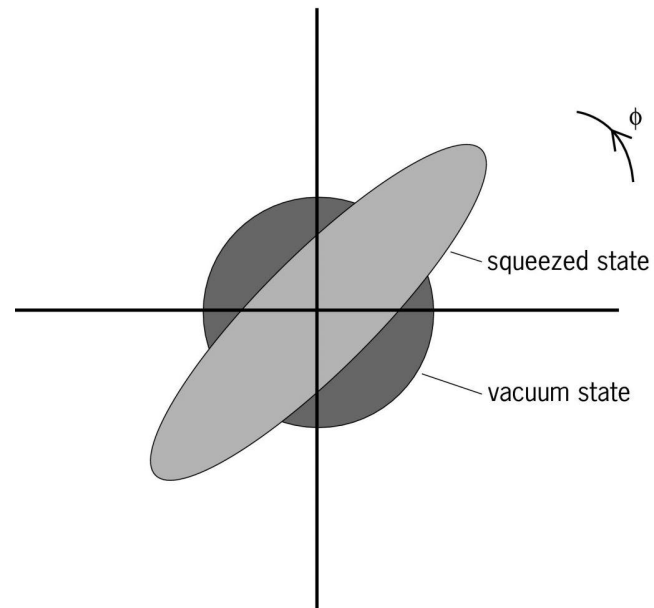
- Eigenstate of annihilation operator: $a|\alpha\rangle = \alpha|\alpha\rangle$, $|\alpha\rangle = e^{\alpha a^\dagger - \alpha^* a} |0\rangle$
- Gaussian wave function with minimal uncertainty
- Oscillatory behavior of $\langle x \rangle$ for a harmonic oscillator (semi-classical)

- Squeezed:

- $e^{\frac{1}{2}(za^{\dagger 2} + z^*a^2)} |0\rangle$

- Thermal:

$$\rho_\omega^{\text{th}} = \frac{1}{Z} \sum_{n=0}^{\infty} e^{-\frac{\hbar\omega(n+1/2)}{k_B T}} |n\rangle\langle n|$$



The standard cosmological perturbation is "squeezed"

$$S'_{\text{pert}} = \frac{1}{2} \int d^4x [(v')^2 - c_s^2 (v_{;i})^2 - 2 \frac{z'}{z} v v' + \left(\frac{z'}{z}\right)^2 v^2]$$

$$\hat{\mathcal{H}}_{\vec{k}} = \hat{\mathcal{H}}_{\vec{k}}^{(0)} + \hat{\mathcal{H}}_{\vec{k}}^{(I)} = \Omega_{\vec{k}} (a_{\vec{k}}^\dagger a_{\vec{k}} + a_{-\vec{k}}^\dagger a_{-\vec{k}} + 1) + i \lambda_{\vec{k}} (e^{-2i\varphi_{\vec{k}}} a_{\vec{k}}^\dagger a_{-\vec{k}} - \text{h.c.})$$

$$\Omega_{\vec{k}} = \frac{k}{2} (1 + c_s^2)$$

$$\lambda_{\vec{k}} = \left[\left(\frac{k}{2} (1 - c_s^2) \right)^2 + \left(\frac{z'}{z} \right)^2 \right]^{\frac{1}{2}}$$

$$\varphi_{\vec{k}} = -\frac{\pi}{2} + \frac{1}{2} \arctan \left(\frac{kz}{2z'} (1 - c_s^2) \right)$$

The time evolution operator of a state can be written as

$$\mathcal{U}_{\mathcal{H}_{\vec{k}}}(\eta, \eta_0) = \mathcal{S}[R_{\vec{k}}, \Phi_{\vec{k}}] \mathcal{R}[\Theta_{\vec{k}}] \quad \mathcal{R}[\Theta_{\vec{k}}] = \exp[-i\Theta_{\vec{k}} (a_{\vec{k}}^\dagger a_{\vec{k}} + a_{-\vec{k}}^\dagger a_{-\vec{k}} + 1)] \quad \text{Two-mode rotation operator}$$

$$\mathcal{S}[R_{\vec{k}}, \Phi_{\vec{k}}] = \exp \left[\frac{R_{\vec{k}}}{2} (e^{-2i\Phi_{\vec{k}}} a_{-\vec{k}} a_{\vec{k}} - \text{h.c.}) \right] \quad \text{Two-mode squeeze operator}$$

$$|SS_{\vec{k}}\rangle = \mathcal{S}[R_{\vec{k}}, \Phi_{\vec{k}}] |0\rangle_{\text{in}} = \sum_{n=0}^{\infty} \frac{1}{\cosh R_{\vec{k}}} (-e^{2i\Phi_{\vec{k}}} \tanh R_{\vec{k}})^n |n, \vec{k}; n, -\vec{k}\rangle$$

$$|n, \vec{k}; n, -\vec{k}\rangle = \sum_{n=0}^{\infty} \frac{1}{n!} (a_{\vec{k}}^\dagger a_{-\vec{k}}^\dagger)^n |0\rangle_{\text{in}}$$

Albrecht, Ferreira, Joyce, Prokopec, 1993

See also Polarski, Starobinsky, 1995

But can the inflationary fluctuation be

- Semiclassical (coherent)?
- Thermal? (thermal inflation, warm inflation)

And decoherence by interaction?

Decoherence of cosmological perturbations:

Horizon exit: Two-mode squeezed state has large $\langle N \rangle$,
but not yet classical

Interaction: Decoherence. Consider wave function of ζ

$$\langle \mathcal{E}, \xi | \Psi(\tau) \rangle = \exp \left(\int_{\mathbf{k}, \mathbf{k}', \mathbf{q}} \mathcal{F}_{\mathbf{k}, \mathbf{k}', \mathbf{q}} \mathcal{E}_{\mathbf{k}} \mathcal{E}_{\mathbf{k}'} \xi_{\mathbf{q}} \right) \Psi_G(\mathcal{E}, \xi)$$

↓ System
↓
↓ Gaussian part

↓ Environment
↓
↓

$$\mathcal{F}_{\mathbf{k}, \mathbf{k}', \mathbf{q}} \approx i \int_{\tau_i}^{\tau} \frac{d\tau'}{H\tau'} \tilde{H}_{\mathbf{k}, \mathbf{k}', \mathbf{q}}^{\text{int}}(\tau') \frac{u_{\mathbf{k}}(\tau') u_{\mathbf{k}'}(\tau') u_{\mathbf{q}}(\tau')}{u_{\mathbf{k}}(\tau) u_{\mathbf{k}'}(\tau) u_{\mathbf{q}}(\tau)} + \mathcal{O} \left((\tilde{H}^{\text{int}})^2 \right)$$

$$\rho_R(\xi, \tilde{\xi}) = \Psi_G^{(\xi)}(\xi) \left[\Psi_G^{(\xi)}(\tilde{\xi}) \right]^* \langle \exp(X) \rangle_{\mathcal{E}}$$

$$\langle \dots \rangle_{\mathcal{E}} \equiv \int D\mathcal{E} \left| \Psi_G^{(\mathcal{E})} \right|^2 (\dots) \quad X \equiv \int_{\mathbf{k}, \mathbf{k}', \mathbf{q}} \mathcal{E}_{\mathbf{k}} \mathcal{E}_{\mathbf{k}'} \left(\xi_{\mathbf{q}} \mathcal{F}_{\mathbf{k}, \mathbf{k}', \mathbf{q}} + \tilde{\xi}_{\mathbf{q}} \mathcal{F}_{\mathbf{k}, \mathbf{k}', \mathbf{q}}^* \right)$$

$$D(\xi, \tilde{\xi}) \equiv \left| \frac{\rho_R(\xi, \tilde{\xi})}{\sqrt{\rho_R(\xi, \xi) \rho_R(\tilde{\xi}, \tilde{\xi})}} \right|$$

Nelson, 1601.03734

Liu, Sou, YW, 1608.07909

Decoherence of cosmological perturbations:

Horizon exit: Two-mode squeezed state has large $\langle N \rangle$,
but not yet classical

Interaction: Decoherence. Consider wave function of ζ

[Importance of the "wave function" of perturbations:

It's the fundamental object,

applications: locality, analyticity, bootstrap, decoherence, ...]

Decoherence of cosmological perturbations:

Horizon exit: Two-mode squeezed state has large $\langle N \rangle$,
but not yet classical

Interaction: Decoherence. Consider wave function of ζ

$$\mathcal{F}_{\mathbf{k},\mathbf{k}',\mathbf{q}} \approx i \int_{\tau_i}^{\tau} \frac{d\tau'}{H\tau'} \tilde{H}_{\mathbf{k},\mathbf{k}',\mathbf{q}}^{\text{int}}(\tau') \frac{u_{\mathbf{k}}(\tau')u_{\mathbf{k}'}(\tau')u_{\mathbf{q}}(\tau')}{u_{\mathbf{k}}(\tau)u_{\mathbf{k}'}(\tau)u_{\mathbf{q}}(\tau)} + \mathcal{O}\left((\tilde{H}^{\text{int}})^2\right)$$



Which interaction? Nelson, 1601.03734: $\mathcal{L}_{\text{int}} = -\frac{M_p^2}{2}\epsilon(\epsilon + \eta)a(t)\zeta^2\partial^2\zeta$

We noted: this boundary term is more important:

$$\mathcal{L}_{\text{bd}} = \partial_t \left(-2a^3 H M_p^2 e^{3\zeta} \right)$$

Order ϵ^{-2} faster decoherence!

A note about the boundary term:

- Fast-varying phase
- Will not change:

$$\begin{aligned} & \int_{\mathcal{M}} d^4x \sqrt{-g} \left[\frac{M_p^2}{2} R_{\text{ADM}} + \mathcal{L}_\phi \right] \\ &= \int_{\mathcal{M}} d^4x \sqrt{-g} \left[\frac{M_p^2}{2} R + \mathcal{L}_\phi \right] + S_{\text{GHY}} \\ &= \int dt d^3x \partial_t (-2a^3 H M_p^2 e^{3\zeta}) + \dots \end{aligned}$$

$$\text{Correlation functions } \langle \zeta^n \rangle = \langle \bar{T} e^{i \int H_{\text{int}} d\tau} \zeta^n T e^{-i \int H_{\text{int}} d\tau} \rangle$$

Isocurvature fluctuations

- But will change:

$\langle \dot{\zeta}^n \rangle$ Do not commute on the boundary,

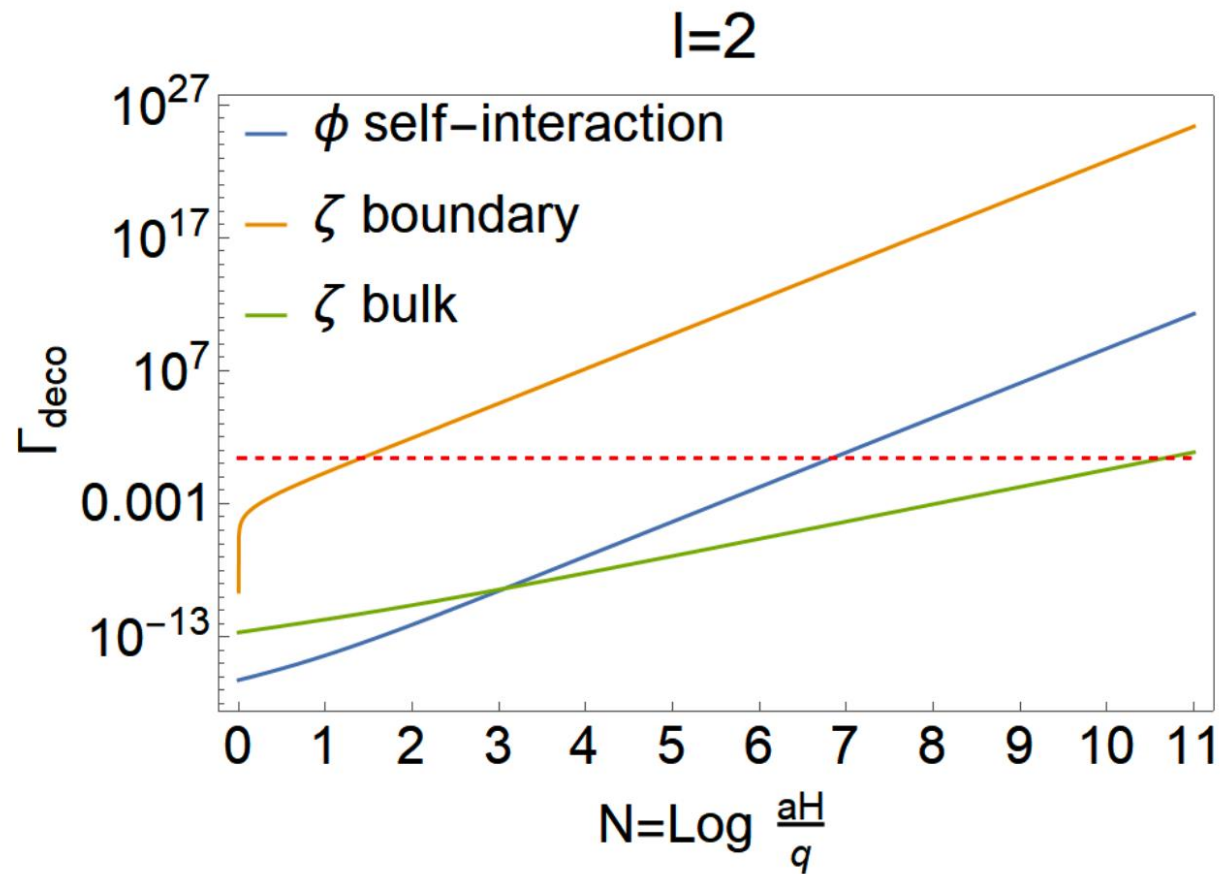
But vanishing for inflation [Bounce? Ultra-slow-roll?]

Decoherence rate

As decoration of wave function

$\delta\varphi$ is more quantum than ζ

Gravity-induced decoherence – Quantum gravity effect!



Decoherence induced by gravity

- Indicator of quantum gravity
- Comparison: Schrodinger-Newton gravity $G_{\mu\nu} = 8\pi G \langle T_{\mu\nu} \rangle$

See Galley, Giacomini, Selby (2020) for other classical gravity models

- But can decoherence rate be observable?

Generalizations:

- Isocurvature
- Gravitational waves
- Quantum creation of the universe

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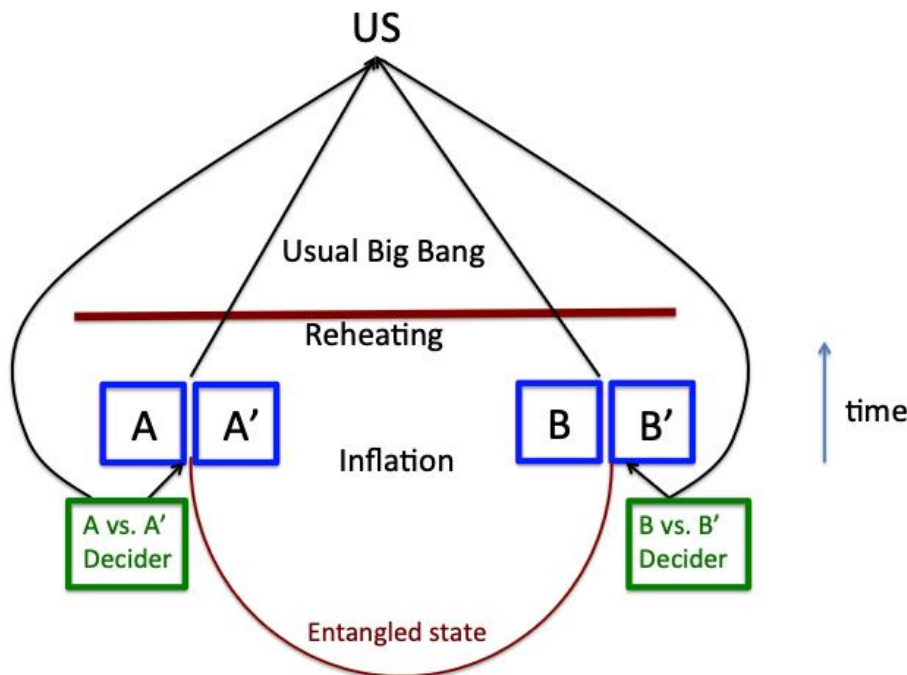
Observational tests?

- Bell inequalities?
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What's the difference between quantum / classical?

1. Bell inequality? Violation of

$$| \langle AB \rangle + \langle AB' \rangle + \langle A'B \rangle - \langle A'B' \rangle | \leq 2$$



Difficulty: non-communicate
observables are rare

Maldacena 1508.01082

See also Kanno & Soda, 1705.06199

What's the difference between quantum / classical?

1. Bell inequality?
2. Quantum noise?

$$S = \frac{1}{16\pi G} \int d^4x \sqrt{-g} R - M_0 \int dt \sqrt{-\dot{X}^2} - m_0 \int dt \sqrt{-\dot{Y}^2}$$

$$S_\omega = \int dt \left(\frac{1}{2} m (\dot{q}^2 - \omega^2 q^2) + \frac{1}{2} m_0 \dot{\xi}^2 - g \dot{q} \dot{\xi} \xi \right)$$

$$\ddot{\xi} = \frac{1}{2} \left(\ddot{h} + \ddot{N}_\Psi - \frac{m_0 G}{c^5} \frac{d^5}{dt^5} \xi^2 \right) \xi$$

Parikh, Wilczek, Zahariade, 2005.07211

Kanno, Soda, Tokuda, 2007.09838

Squeezing \Rightarrow Large quantum noise

Can quantum noise be reformulated to fit the general framework of cosmic perturbations?

Is it a matter of operator ordering?

- Linear perturbations: equivalent (Chen, Namjoo, YW 2015)
- Nonlinear level?

What's the difference between quantum / classical?

1. Bell inequality?
2. Quantum noise?
3. Position uncertainty of dark matter particles?

- Without decoherence:

$$s(t) = \sqrt{\frac{s_0^4 + (\hbar t/m)^2}{s_0^2}} \quad s_{\min}(t) = \sqrt{\frac{2\hbar t}{m}} \quad (10^{11} \text{ GeV} \Leftrightarrow 1 \text{ m})$$

- Decoherence: $m < 10^7 \text{ kg}$

(Lindblad equation)

Neppoleon, Iyer, Vedral, YW, 2110.13438

See also Allali, Hertzberg 2012.12903 for ALPS

What's the difference between quantum / classical?

1. Bell inequality?
2. Quantum noise?
3. Position uncertainty of dark matter particles?
4. What about we see nothing – null test

Decoherence, if happened gravitationally
also indicates quantum gravity!

Conclusion:

- Squeezed state & decoherence
- Test quantum gravity? A null test may also be a test.