#### Peng Huanwu Center for Fundamental Physics, USTC April 9, 2024



# **PULSAR POLARIZATION ARRAY**

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#### **Astronomical Clock**





- Extraordinary astronomical clocks
- Misalignment between spin and magnetic axes => Send pulses with extremely high regularity







#### **Pulsars to Be Observed**



[Rauber Beck (2009)]

• Rich in Milk Way: to be observed (~20000, grey points) vs. observation (~3000, yellow points)



### Pulsar Timing Array (PTA)



- A network of widely distributed and well-timed millisecond pulsars (MSPs)
- A galactic timing interferometer to detect ~nanoHz gravitational waves (GWs) [S. L. Detweiler, Atrophy's. J. 234 (1979)]



### Leading PTA Programs

- Parkes Pulsar Timing Array (PPTA; Australia)
  - Parkes radio-telescope
- North American Nanohertz Observatory for Gravitational Waves (NANOGrav; Canada and USA)
  - Arecibo and Green Bank radio telescopes
- European Pulsar Timing Array (EPTA; Europe)
  - Westerbork Synthesis Radio Telescope, the Effelsberg Radio Telescope, the Lovell Telescope, the Nançay Radio Telescope and the Sardinia Radio Telescope.
- Indian Pulsar Timing Array Experiment (InPTA; India and Japan)
  - Upgraded Giant Meterwave Radio Telescope
- Chinese Pulsar Timing Array (CPTA; China)
  - Five hundred meter Aperture Spherical Telescope
  - QTT 110m full-band radio telescope at Qitai, Xinjiang
  - JRT 120m low-band telescope in Jindong, Yunnan
- SKA-PTA (e.g., MeerTime PTA as a precursor)
  - Square Kilometer Array telescope

Over 80 MSPs have been monitored by the global PTA network over a timespan of years. Future observations (SKA/FAST) can increase this number to O(100), with higher timing precision





### Why Arrayed Pulsars?

THE ASTROPHYSICAL JOURNAL, **265**:L39–L42, 1983 February 15 © 1983. The American Astronomical Society. All rights reserved. Printed in U.S.A.

#### UPPER LIMITS ON THE ISOTROPIC GRAVITATIONAL RADIATION BACKGROUND FROM PULSAR TIMING ANALYSIS<sup>1</sup>

R. W. HELLINGS AND G. S. DOWNS Jet Propulsion Laboratory, California Institute of Technology Received 1982 October 1; accepted 1982 October 20



per limit of about 10° s for the periods to which a are sensitive. It should also be noted from h(t) that data from any pulsar contain informabut h(t) at the time and place of reception (i.e., a) and about the value h(t) had at the pulsar at of emission of the signal. Thus, data from any will have a gravitational wave signal in common other pulsars (though with an amplitude scaled  $\cos \theta$ ) as well as a component of the signal

which will be independent of the others due to the long light times between pulsars compared with the 12 yr data span. When data from several pulsars are cross-correlated, this common signal will allow one to dig into the pulsar noise to detect a possible common gravitational wave signal.

#### b) Cross-Correlation

The fractional frequency shift observed in the data on pulsar number *i* may be written

Hellings-Downs Curve

Encodes exactly the cross-correlation of pulsar timing data that would indicate a common GW signal.



#### **A Milestone**







#### **Astronomical Linear Polarizer**



- Astronomical linear polarizers
- Polarization is often measured for calibrating pulsar observation

![](_page_7_Picture_6.jpeg)

![](_page_8_Picture_0.jpeg)

#### **Astronomical Linear Polarizer**

Cata	ATNF Pulsar Catalogue										
1 2 3 4 5	J0002+6216 J0006+1834 J0007+7303 J0011+08 J0014+4746	<u>cwp+17</u> <u>cnt96</u> <u>aaa+09c</u> <u>dsm+16</u> dtb78	8.6682478274 1.4414462816 3.165827392 0.391716 0.805997239145	1 3 3 0 7	<u>cwp+17</u> <u>cn95</u> <u>awd+12</u> <u>dsm+16</u> blk+04	218.6 11.41 * 24.9 30.405	6 55 0 13	wcp+18 bkk+16 * dsm-16 bkk+16	* -20 * -15-56	0 * 3 <u>npn+20</u> 0 * 0 * 10 sbg+19	6.357 0.860 1.400 5.399 1.776
111 112 113 114 115	J0211-8159 J0212+5222 J0214+5222 J0215+6218 J0218+4232	lml+98 bck+13 slr+14 llc98 nbf+95	0.9282183663 2.65684439046 40.6912718043 1.821892452777 430.461054545748	7 1 4 9 15	dsb+98 lsk+18 slr+14 hlk+04 dcl+16	24.36 38.21 22.0354 84.00 61.252	3 34 5 5	dsb+98 lsk+18 slr+14 hlk+04 hlk+04	54 -13.68 -16.44 380.9 -61.40	9 hml+06 8 sbg+19 7 sbg+19 10 hmvd18 8 sbg+19	1.523 1.558 1.161 2.004 3.150
1906 1907 1908 1909 1910	J1810+1744 J1810-1820 J1810-2005 J1810-5338 J1811-0154	<u>hrm+11</u> mhl+02 clm+01 mlt+78 ebvb01	602.409639 6.5054918751 30.467142155106 3.8306868647 1.08114557226	0 3 7 4 5	hrm+11 mhl+02 jsb+10 lbs+20 ebvb01	39.7 452.2 241.0 45 148.1	0 25 3 2 2	<u>hrn+11</u> mhl+02 jsb+10 nmc81 mss-20	88.5 110.3 -15 58 46	1 <u>sbg+19</u> 81 <u>hmvd18</u> 14 <u>hmvd18</u> 3 <u>qmlg95</u> 11 <u>njkk08</u>	2.361 4.237 3.514 1.647 11.112
3276	J2257+5909 J2301+5852	<u>dls72</u> <u>fg81</u>	2.71557265035 0.14328554678 192.5919636477142	5 3 8	<u>hlk+04</u> <u>dk14</u> aab+21a	151.082 * 13.788120	6 0 0	<u>hlk+04</u> * aab+21a	-323.5 * 19.1	4 <u>fdr15</u> 0 * 16 <u>npn+20</u>	3.000 3.300 0.863

Can we cross-correlate pulsar polarization data, as done for the timing data, to give full play to its capability in exploring astrophysics and fundamental physics?

![](_page_9_Picture_0.jpeg)

#### PHYSICAL REVIEW LETTERS 130, 121401 (2023)

#### [arXiv:2111.10615]

#### **Pulsar Polarization Arrays**

Tao Liu,<sup>1,\*</sup> Xuzixiang Lou<sup>(D)</sup>,<sup>1,†</sup> and Jing Ren<sup>(D)</sup><sup>2,‡</sup>

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Received 21 February 2022; revised 31 December 2022; accepted 27 February 2023; published 23 March 2023)

Pulsar timing arrays (PTAs) consisting of widely distributed and well-timed millisecond pulsars can serve as a galactic interferometer to measure gravitational waves. With the same data acquired for PTAs, we propose to develop pulsar polarization arrays (PPAs), to explore astrophysics and fundamental physics. As in the case of PTAs, PPAs are best suited to reveal temporal and spatial correlations at large scales that are hard to mimic by local noise. To demonstrate the physical potential of PPAs, we consider detection of ultralight axionlike dark matter (ALDM), through cosmic birefringence induced by its Chern-Simons coupling. Because of its tiny mass, the ultralight ALDM can be generated as a Bose-Einstein condensate, characterized by a strong wave nature. Incorporating both temporal and spatial correlations of the signal, we show that PPAs have a potential to probe the Chern-Simons coupling up to  $\sim 10^{-14} - 10^{-17}$  GeV<sup>-1</sup>, with a mass range  $\sim 10^{-27} - 10^{-21}$  eV.

DOI: 10.1103/PhysRevLett.130.121401

PTAs: suited for revealing physics with a common correlated timing signal PPAs: suited for revealing physics with a common correlated polarization signal

![](_page_9_Picture_11.jpeg)

![](_page_10_Picture_0.jpeg)

### **Big Questions for Particle Physicists**

![](_page_10_Figure_2.jpeg)

#### [Working group report (particle physics) for Snowmass 2013]

As one scientific case, we consider the detection of axion-like wave Dark Matter

![](_page_10_Picture_5.jpeg)

### Wave Dark Matter (WDM)

![](_page_11_Figure_1.jpeg)

Wave Dark Matter: Bosonic and ma << 30 eV

=> Large occupation number per de Broglie volume (NdB >> 1) in a Milky-Way-like environment

=> Formation of a coherent state with strong wave nature within galaxies

12

![](_page_12_Picture_0.jpeg)

### **Axion-like WDM**

![](_page_12_Figure_2.jpeg)

Phys. Rev. D 95 (2017)]

Axion-like particles are probably the most important WDM candidate. Their coherent state usually oscillates as

$$a(\mathbf{x},t) = a_0(\mathbf{x},t)\cos(m_a t + m_a \mathbf{v} \cdot \mathbf{x} + \phi)$$

- Period is determined by the axion mass in temporal direction and its momentum in spatial direction;
- Amplitude is determined by energy density of DM halo

$$\rho_{\rm DM}(\mathbf{x},t) = \frac{1}{2}m_a^2 a_0^2(\mathbf{x},t) + \mathcal{O}(v^2)$$

 Stochastic - essentially an uncorrelated superposition of particle waves

$$a(\mathbf{x},t) \approx \frac{\sqrt{\rho(\mathbf{x})}}{m_a} \sum_{\mathbf{v}\in\Omega} C_{\mathbf{v}} \cos[m_a(t-\mathbf{v}\cdot\mathbf{x})+\psi_{\mathbf{v}}]_{\mathbf{x}}$$

![](_page_12_Picture_10.jpeg)

![](_page_13_Figure_0.jpeg)

Fuzzy DM [Hu, et. al., Phys.Rev.Lett. 85 (2000)]: ma ~ 10^-21 - 10^-22 eV (oscillation period 0  $2^{i}$  pi/ma ~ 1 yr, with a dB wavelength ~ O(100) pc), where small-scale problems on astronomical structure could be addressed.

AUX 2014 HOUTE N

![](_page_13_Figure_2.jpeg)

![](_page_14_Picture_0.jpeg)

![](_page_14_Picture_1.jpeg)

## A joint Fermilab/SLAC publication

follow +

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Illustration by Sandbox Studio, Chicago with Ariel Davis Is dark matter the most powerful wave in the universe?

04/04/23 | By Kimberly Hickok

Dark matter could consist of particles so ultralight, they behave more like waves.

![](_page_14_Picture_6.jpeg)

![](_page_15_Picture_0.jpeg)

Axion-like WDM can affect pulsar polarization via an effect known as ``cosmological birefringence''

![](_page_15_Figure_3.jpeg)

![](_page_16_Figure_0.jpeg)

$$\omega_{\pm} \simeq k \pm g \left( \frac{\partial a}{\partial t} + \nabla a \cdot \frac{\mathbf{k}}{k} \right) \xrightarrow[\text{relativistic}]{\text{non-}} k \pm g \frac{\partial a}{\partial t}$$

Parity-violating Chern-Simons term => Different dispersion relations for left- and right-circular polarized light => Position angle rotated for the linearly polarized light traveling across an axion field (including axion-like WDM halo)

![](_page_16_Picture_4.jpeg)

### **Comparison with Faraday Rotation (FR)**

![](_page_17_Figure_1.jpeg)

- CB: determined by the difference of axion field profile between two endpoints of the light path due to the topological nature of Chern-Simons coupling. VS FR: relies on path length directly.
- CB: no frequency dependence. VS FR: increases with wavelength square.
- CB: features oscillation with a period of 2\*pi/ma. VS FR: no characteristic time dependence is expected.

![](_page_17_Picture_5.jpeg)

#### [Lue, Wang, KamionKowski, PRL(1999)]

#### **CMB-Based Detection**

![](_page_18_Figure_2.jpeg)

PA rotation - Determined by the difference of axion field between at recombination and for the Universe today

NASA/WMAP Science Test

![](_page_19_Picture_0.jpeg)

#### **CMB-Based Detection**

## **Testing a Universal Symmetry**

Searching for *CPT* Violation with Cosmic Microwave Background Data from WMAP and BOOMERANG

Bo Feng, Mingzhe Li, Jun-Qing Xia, Xuelei Chen, and Xinmin Zhang

Phys. Rev. Lett. 96, 221302 (2006)

Published June 7, 2006

June 12, 2006 • Phys. Rev. Focus 17, 21

The cosmic microwave background that fills the Universe provides a test for asymmetries in the laws of physics.

![](_page_19_Figure_9.jpeg)

NASA/WMAP Science Team

**Physics test.** An analysis of the polarization (white lines) of the cosmic microwave background measured across the entire sky can test for violations of the fundamental symmetry known as *CPT*.

![](_page_19_Picture_12.jpeg)

#### **CMB-Based Detection**

A&A 596, A110 (2016) DOI: 10.1051/0004-6361/201629018 © ESO 2016

#### Astronomy Astrophysics

#### **Planck** intermediate results

#### XLIX. Parity-violation constraints from polarization data

Planck Collaboration: N. Aghanim<sup>47</sup>, M. Ashdown<sup>57,4</sup>, J. Aumont<sup>47</sup>, C. Baccigalupi<sup>67</sup>, M. Ballardini<sup>23, 38, 41</sup>, A. J. Banday<sup>77,7</sup>, R. B. Barreiro<sup>52</sup>, N. Bartolo<sup>22, 53</sup>, S. Basak<sup>67</sup>, K. Benabed<sup>48, 76</sup>, J.-P. Bernard<sup>77, 7</sup>, M. Bersanelli<sup>26, 39</sup>, P. Bielewicz<sup>65, 7, 67</sup>, L. Bonavera<sup>12</sup>, J. R. Bond<sup>6</sup>, J. Borrill<sup>9, 73</sup>, F. R. Bouchet<sup>48,72</sup>, C. Burigana<sup>38,24,41</sup>, E. Calabrese<sup>74</sup>, J.-F. Cardoso<sup>60,1,48</sup>, J. Carron<sup>17</sup>, H. C. Chiang<sup>19,5</sup>, L. P. L. Colombo<sup>15,54</sup>, B. Comis<sup>61</sup>, D. Contreras<sup>14</sup>, F. Couchot<sup>58</sup>, A. Coulais<sup>59</sup>, B. P. Crill<sup>54,8</sup>, A. Curto<sup>52,4,57</sup>, F. Cuttaia<sup>38</sup>, P. de Bernardis<sup>25</sup>, A. de Rosa<sup>38</sup>, G. de Zotti<sup>35,67</sup>, J. Delabrouille<sup>1</sup>, F.-X. Désert<sup>45</sup>, E. Di Valentino<sup>48,72</sup>, C. Dickinson<sup>55</sup>, J. M. Diego<sup>52</sup>, O. Doré<sup>54,8</sup>, A. Ducout<sup>48,46</sup>, X. Dupac<sup>30</sup>, S. Dusini<sup>53</sup>, F. Elsner<sup>16,48,76</sup>, T. A. Enßlin<sup>63</sup>, H. K. Eriksen<sup>50</sup>, Y. Fantaye<sup>29</sup>, F. Finelli<sup>38,41</sup>, F. Forastieri<sup>24,42</sup>, M. Frailis<sup>37</sup>, E. Franceschi<sup>38</sup>, A. Frolov<sup>71</sup>, S. Galeotta<sup>37</sup>, S. Galli<sup>56</sup>, K. Ganga<sup>1</sup>, R. T. Génova-Santos<sup>51,11</sup>, M. Gerbino<sup>75,66,25</sup>, Y. Giraud-Héraud<sup>1</sup>, J. González-Nuevo<sup>12,52</sup>, K. M. Górski<sup>54,79</sup>, A. Gruppuso<sup>38,41,\*</sup>, J. E. Gudmundsson<sup>75,66,19</sup>, F. K. Hansen<sup>50</sup>, S. Henrot-Versillé<sup>58</sup>, D. Herranz<sup>52</sup>, E. Hivon<sup>48,76</sup>, Z. Huang<sup>6</sup>, A. H. Jaffe<sup>46</sup>, W. C. Jones<sup>19</sup>, E. Keihänen<sup>18</sup>, R. Keskitalo<sup>9</sup>, K. Kiiveri<sup>18, 34</sup>, N. Krachmalnicoff<sup>26</sup>, M. Kunz<sup>10, 47, 2</sup>, H. Kurki-Suonio<sup>18, 34</sup>, J.-M. Lamarre<sup>59</sup>, M. Langer<sup>47</sup>, A. Lasenby<sup>4, 57</sup>, M. Lattanzi<sup>24, 42</sup>, C. R. Lawrence<sup>54</sup>, M. Le Jeune<sup>1</sup>, J. P. Leahy<sup>55</sup>, F. Levrier<sup>59</sup>, M. Liguori<sup>22, 53</sup>, P. B. Lilje<sup>50</sup>, V. Lindholm<sup>18,34</sup>, M. López-Caniego<sup>30</sup>, Y.-Z. Ma<sup>55,68</sup>, J. F. Macías-Pérez<sup>61</sup>, G. Maggio<sup>37</sup>, D. Maino<sup>26,39</sup>, N. Mandolesi<sup>38,24</sup>, M. Maris<sup>37</sup>, P. G. Martin<sup>6</sup>, E. Martínez-González<sup>52</sup>, S. Matarrese<sup>22, 53, 32</sup>, N. Mauri<sup>41</sup>, J. D. McEwen<sup>64</sup>, P. R. Meinhold<sup>20</sup>, A. Melchiorri<sup>25, 43</sup>, A. Mennella<sup>26, 39</sup> M. Migliaccio<sup>49,57</sup>, M.-A. Miville-Deschênes<sup>47,6</sup>, D. Molinari<sup>24,38,42</sup>, A. Moneti<sup>48</sup>, G. Morgante<sup>38</sup>, A. Moss<sup>70</sup>, P. Natoli<sup>24,3,42</sup>, L. Pagano<sup>25,43</sup>, D. Paoletti<sup>38,41</sup>, G. Patanchon<sup>1</sup>, L. Patrizii<sup>41</sup>, L. Perotto<sup>61</sup>, V. Pettorino<sup>33</sup>, F. Piacentini<sup>25</sup>, L. Polastri<sup>24,42</sup>, G. Polenta<sup>3,36</sup>, J. P. Rachen<sup>13,63</sup>, B. Racine<sup>1</sup>, M. Reinecke<sup>63</sup>, M. Remazeilles<sup>55,47,1</sup>, A. Renzi<sup>29,44</sup>, G. Rocha<sup>54,8</sup>, C. Rosset<sup>1</sup>, M. Rossetti<sup>26,39</sup>, G. Roudier<sup>1,59,54</sup>, J. A. Rubiño-Martín<sup>51,11</sup>, B. Ruiz-Granados<sup>78</sup>, M. Sandri<sup>38</sup>, M. Savelainen<sup>18,34</sup>, D. Scott<sup>14</sup>, C. Sirignano<sup>22,53</sup>, G. Sirri<sup>41</sup>, L. D. Spencer<sup>69</sup>, A.-S. Suur-Uski<sup>18,34</sup>, J. A. Tauber<sup>31</sup>, D. Tavagnacco<sup>37,27</sup>, M. Tenti<sup>40</sup>, L. Toffolatti<sup>12,52,38</sup>, M. Tomasi<sup>26,39</sup>, M. Tristram<sup>58</sup>, T. Trombetti<sup>38,24</sup>, J. Valiviita<sup>18,34</sup>, F. Van Tent<sup>62</sup>, P. Vielva<sup>52</sup>, F. Villa<sup>38</sup>, N. Vittorio<sup>28</sup>, B. D. Wandelt<sup>48,76,21</sup>, I. K. Wehus<sup>54,50</sup>, A. Zacchei<sup>37</sup>, and A. Zonca<sup>20</sup>

Becomes a standard task for the CMB missions today

![](_page_20_Picture_7.jpeg)

![](_page_21_Picture_0.jpeg)

### [TL, G. Smoot, Y. Zhao, arXiv:1901.10981] Pulsar Light-Based Detection

PA rotation - Determined by the difference of axion field between near the PSR at the time of photon emission and around the Earth at the moment of photon receiving.

![](_page_21_Picture_3.jpeg)

![](_page_22_Picture_0.jpeg)

#### **Detecting Axion-like WDM with PPAs**

WDM Two important facts on the axion-like WDM It can influence the polarization of pulsar light while it travels across galactic DM halo • As it is wave-like, its influence for pulsars in the Milky Way is common and correlated => PPAs are especially suited for its detection !

### Cross-correlation: GWs VS WDM

![](_page_23_Figure_1.jpeg)

Pulsar term

Timing residue caused by stochastic GWs

$$\Delta T(t) = \int_{-\infty}^{\infty} df \frac{1}{2} u^a u^b h_{ab}(f, \hat{n}) \frac{1}{i2\pi f} \frac{1}{1 + \hat{n} \cdot \hat{u}} \left[ e^{i2\pi f(t_2 + \hat{n} \cdot \vec{r}_2/c)} - e^{i2\pi f(t_1 + \hat{n} \cdot \vec{r}_1/c)} \right]$$

PA rotation caused by stochastic axion-like WDM

$$\Delta \theta_p(t) = \frac{g}{m_a} \int \alpha_{\mathbf{v}} \left\{ \sqrt{\rho_p f_p(\mathbf{v})} \cos[m_a(t - L_p - \mathbf{v} \cdot \mathbf{x}_p) + \phi_{\mathbf{v}}] - \sqrt{\rho_e f_e(\mathbf{v})} \cos(m_a t + \phi_{\mathbf{v}}) \right\} d^3 \mathbf{v}$$

	SGWB (PTAs)	Axion-like WDM (PPAs)
Earth-Earth Term	quadrupolar correlation (Hellings-Downs curve)	monopolar correlation
Pulsar-Pulsar Term	spatial correlation degrades quickly (L»ldв~1/w)	spatial correlation degrades much slower (L»ldB»1/ma), enhanced at galactic center

![](_page_23_Picture_8.jpeg)

![](_page_24_Picture_0.jpeg)

#### **NANOGrav Anomaly**

Search...

![](_page_24_Figure_2.jpeg)

NANOGrav reported anomalous sinusoidal trends of PA for a set of pulsars with a period ~ 1-2 years

$$\Delta PA(t) = \Delta RM(t) * \lambda^2$$

Such an annual sinusoidal trend for individual pulsars can be wellexplained by the CB of axion-like WDM (for ma ~ 10^-22 eV ~ O(1)yr^-1)

$$\Delta PA(t) = \Delta \theta(t)$$

Question: How to exclude or confirm the possibility of axion-like WDM?

![](_page_24_Picture_8.jpeg)

![](_page_25_Picture_0.jpeg)

![](_page_25_Figure_2.jpeg)

Question: How to exclude or confirm the possibility of axion-like WDM?

- The data for two individual MSPs indicates an anomalous peak for ma ~ 0.7 \*10^-22 eV
- However, strongly disfavored by the cross-correlation analysis

![](_page_25_Picture_6.jpeg)

The anomalous NANOGrav sinusoidal trends in polarization data are very unlikely to be caused by axion-like WDM.

![](_page_25_Picture_8.jpeg)

![](_page_26_Picture_0.jpeg)

#### **Sensitivity Projection for Benchmark PPAs**

![](_page_26_Figure_2.jpeg)

- NPPA: 100 MSPs around the Earth; 10 years' observation with a cadence 10/ yr; noise variance - 1deg<sup>2</sup>
- FPPA: 100 MSPs near the galactic center; 10 years' observation with a cadence 10/yr; noise variance - 1deg<sup>2</sup>
- OPPA: 100 MSPs following the ATNF pulsar distribution; 30 years' observation with a cadence 1/week; noise variance - 1deg<sup>2</sup>

- The projected PPA limits form a complementarity with the existing CMB bounds
- With noise variance ~ (0.1 deg)^2, the limits can be improved by one more order of magnitude

![](_page_26_Picture_8.jpeg)

![](_page_27_Picture_0.jpeg)

### **Global Efforts of Searching for Ultralight Axions**

![](_page_27_Figure_2.jpeg)

![](_page_27_Picture_3.jpeg)

### **Global Efforts of Searching for Ultralight Axions**

![](_page_28_Figure_2.jpeg)

"..... The ultimate reach of the PTA approach is to think about pulsar <u>polarisation</u> arrays, in which the <u>polarisation</u> information on different pulsars in the network is spatially and temporally correlated rather than considered separately as a measure of Faraday rotation."

Ciaran A. J. O'Hare [arXiv:2403.17697 [hep-ph]]

![](_page_28_Picture_6.jpeg)

![](_page_29_Picture_0.jpeg)

## High-F band detection

#### A Landscape

![](_page_29_Picture_3.jpeg)

![](_page_29_Picture_4.jpeg)

Radio band detection

![](_page_29_Picture_6.jpeg)

![](_page_29_Picture_7.jpeg)

PPA detection (non-gravitational)

PTA detection (gravitational)

![](_page_29_Picture_10.jpeg)

![](_page_30_Picture_0.jpeg)

### **PTA Detection**

### $\exists \mathbf{r} \times \mathbf{i} \mathbf{V} > astro-ph > arXiv:1309.5888$

Search...

Help | Advance

#### Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 23 Sep 2013]

#### Pulsar timing signal from ultralight scalar dark matter

#### Andrei Khmelnitsky, Valery Rubakov

An ultralight free scalar field with mass around  $10^{-23} - 10^{-22}$  eV is a viable dark mater candidate, which can help to resolve some of the issues of the cold dark matter on subgalactic scales. We consider the gravitational field of the galactic halo composed out of such dark matter. The scalar field has oscillating in time pressure, which induces oscillations of gravitational potential with amplitude of the order of  $10^{-15}$  and frequency in the nanohertz range. This frequency is in the range of pulsar timing array observations. We estimate the magnitude of the pulse arrival time residuals induced by the oscillating gravitational potential. We find that for a range of dark matter masses, the scalar field dark matter signal is comparable to the stochastic gravitational wave signal and can be detected by the planned SKA pulsar timing array experiment.

Oscillating halo density => Oscillating metric => Oscillating timing residual

![](_page_30_Picture_11.jpeg)

![](_page_31_Picture_0.jpeg)

#### **First Parkes PTA Measurement**

![](_page_31_Figure_2.jpeg)

Parkes Pulsar Timing Array constraints on ultralight scalar-field dark matter

Nataliya K. Porayko *et al.* (PPTA Collaboration) Phys. Rev. D **98**, 102002 – Published 5 November 2018

![](_page_32_Picture_0.jpeg)

### Gamma-Ray PTA

Science	Current Issue	First rele	First release papers				More 🗸		
HOME > SCIENCE > VOL. 376, NO. 6592 > A GAMM	IA-RAY PULSAR TIMING A	RRAY CONSTRA	INS TH	E NAN(	OHERT	Z GRAV	ITA <del>.</del>		
REPORT GRAVITATIONAL WAVES		f	y	in	Ð	<b>FQ</b>	$\bowtie$		
A gamma-ray pulsar timing array constrains the nanohertz gravitational wave background									
THE FERMI-LAT COLLABORATION Authors Info & Affiliat	ions								
SCIENCE • 7 Apr 2022 • Vol 376, Issue 6592 • pp. 521-52	23 • DOI: 10.1126/science.	.abm3231							

While being limited by statistics, such a high-frequency PTA can benefit from a suppression of intrinsic red noise of dispersion measure variance

![](_page_32_Picture_4.jpeg)

![](_page_33_Picture_0.jpeg)

![](_page_33_Figure_2.jpeg)

![](_page_34_Picture_0.jpeg)

The PPA and PTA can be correlated to further strengthen their capability in identifying the nature of signals.

![](_page_34_Figure_3.jpeg)

- PTA detection is essentially gravitational (measuring energy density of DM halo), while the PPA detection is non-gravitational (measuring axion Chern-Simons coupling)
- Non-gravitational detection is highly important for exploring the DM properties. Our ignorance on the DM properties to a great extent is due to the limitation of our main knowledge source to gravitational detections

![](_page_34_Picture_6.jpeg)

![](_page_35_Picture_0.jpeg)

- To fully extend the physical reach of pulsars as a precision astronomical tool, we have developed the methodology of Pulsar Polarization Array
- As one scientific case, we demonstrated that the PPA can be applied to detect the axion-like WDM as a common correlated signal. The first results based on real data are expected to come soon.
- In view of its non-gravitational nature, PPA (polarization data) can be correlated with PTA (timing data) to further strengthen their capability in identifying the nature of signals
- Stay tuned ... ...

![](_page_35_Picture_6.jpeg)

![](_page_36_Picture_0.jpeg)

The second location

Mill Sil

#### CRF under Grant No. C6017-20G

![](_page_36_Picture_2.jpeg)

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