

Feynman Integrals, Symbology and Twistor geometries



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2203.16112; 2206.04609 w. 张扬,马柔柔,吴子昊; 何颂, 李振杰;

2207.13482 w.何颂,唐一朝,刘家昊;

2212.09762 w. Matthias Wilhelm, Anne Spiering, Roger Morales and 张驰 from NBI

Feynman Integrals : Mathematical structures and calculation

- Bottom-up: direct integration, reduction, canonical differential equation...

- Up-bottom: physical singularities for certain integrals can be predicted before results are known

(cluster algebras, kinematics quivers, Schubert analysis...)

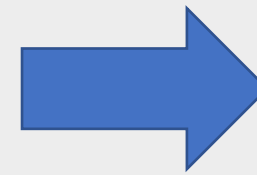
Planar, massless
propagators in 4D

- → bootstrapping at symbol/function level

$$dG^{(w)} = \sum_i G_i^{(w-1)} d \log x_i \implies \mathcal{S}[G^{(w)}] = \sum_i \mathcal{S}[G_i^{(w-1)}] \otimes \log x_i$$

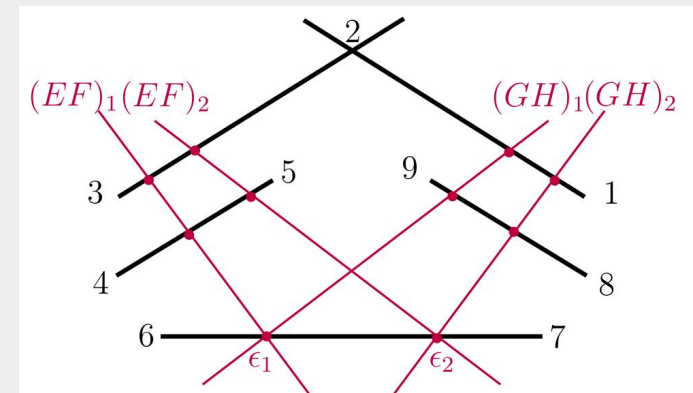
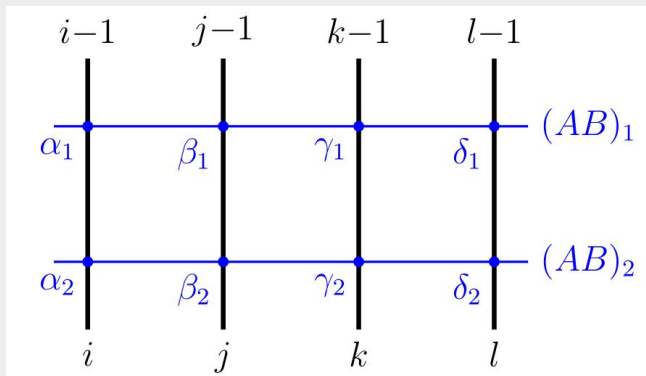
Twistor geometries associated to Feynman Integrals

- Three parts:
- 1. geometries (lines) from kinematics
- 2. geometries (lines) from Leading Singularities
- 3. Intersection points



**geometrical invariants
In twistor space**

Covering all possible physical singularities for individual integrals!



1. Twistor geometries from kinematics

Momentum twistors:

$$\lambda_a \tilde{\lambda}_a = p_a = x_{a+1} - x_a$$

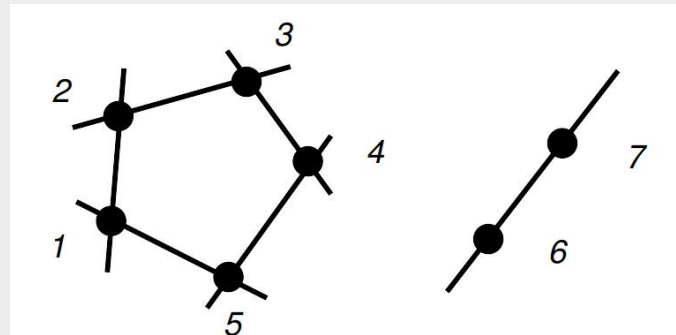
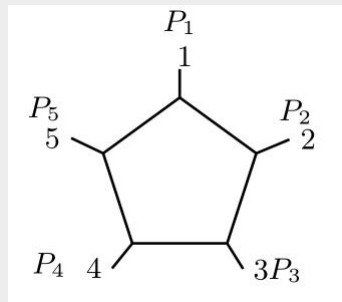
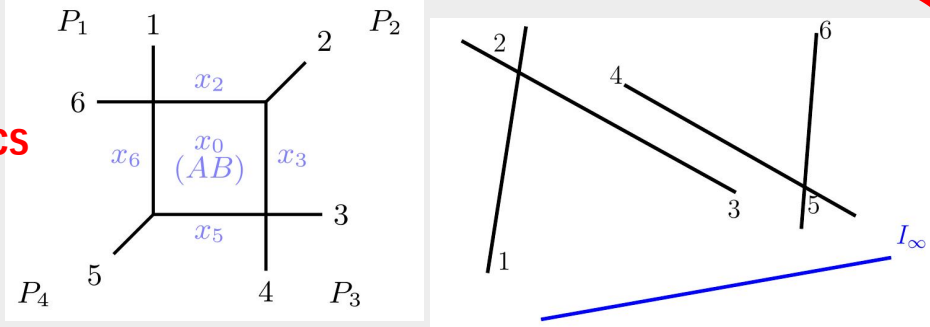
$$Z_a^I = (\lambda_a^\alpha, x_{a,\alpha\dot{\alpha}} \lambda_a^\alpha) \in \mathbb{P}^3, \quad x_a \sim Z_{a-1} \wedge Z_a$$

$$(x_a - x_b)^2 = \frac{\langle a-1ab-1b \rangle}{\langle a-1aI_\infty \rangle \langle b-1bI_\infty \rangle}$$

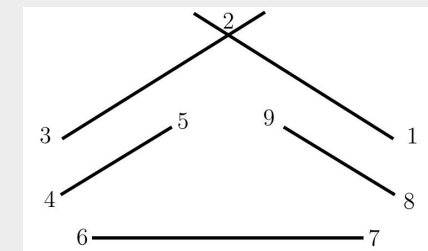
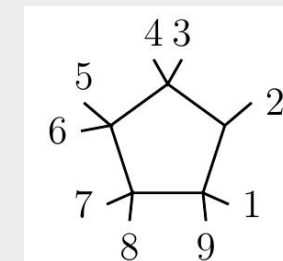
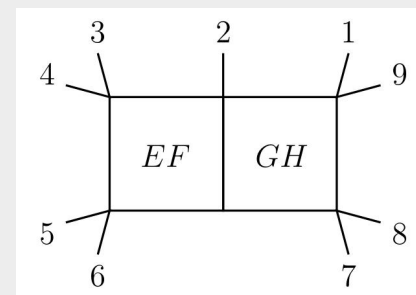
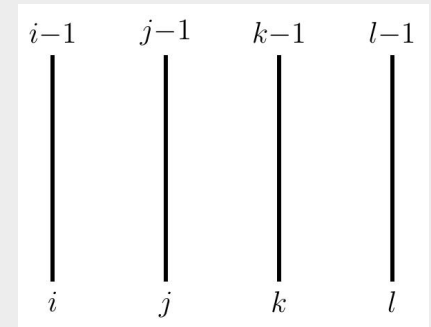
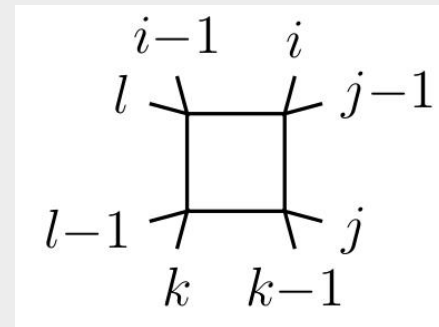
Dual momenta: lines in \mathbb{P}^3

Kinematics: several intersected lines, together with reference I_∞

General kinematics

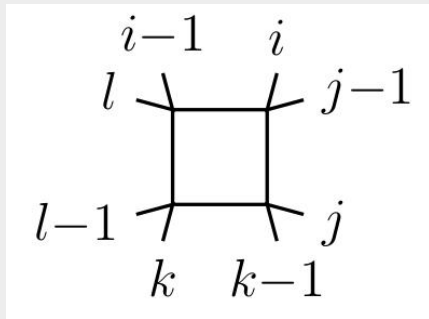


Dual Conformal Invariant



2. Twistor geometries from Leading Singularities (LS)

One-loop leading singularities: residue of the integrand under 4 on-shell conditions of propagators:



$$:= \int d^4\ell \frac{(x_i - x_k)^2 (x_j - x_l)^2}{(\ell - x_i)^2 (\ell - x_j)^2 (\ell - x_k)^2 (\ell - x_l)^2}$$

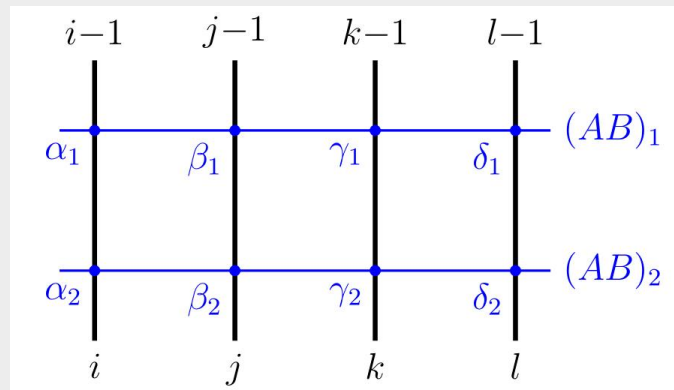
$$= \int_{AB} \frac{\langle i-1 | k-1 \rangle \langle j-1 | l-1 \rangle}{\langle AB | i-1 \rangle \langle AB | j-1 \rangle \langle AB | k-1 \rangle \langle AB | l-1 \rangle}$$

Two lines (AB) and (l-1 l) intersects

$$\langle AB | i-1 \rangle = \langle AB | j-1 \rangle = \langle AB | k-1 \rangle = \langle AB | l-1 \rangle = 0.$$

- D=4, loop momentum is fixed by the four conditions
- AB intersects with 4 external lines simultaneously (exactly two solutions in \mathbb{P}^3)
[Schubert, 19th century; N. Arkani-Hamed et al, 1012.6032]

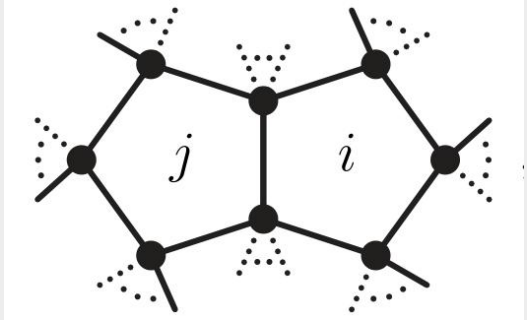
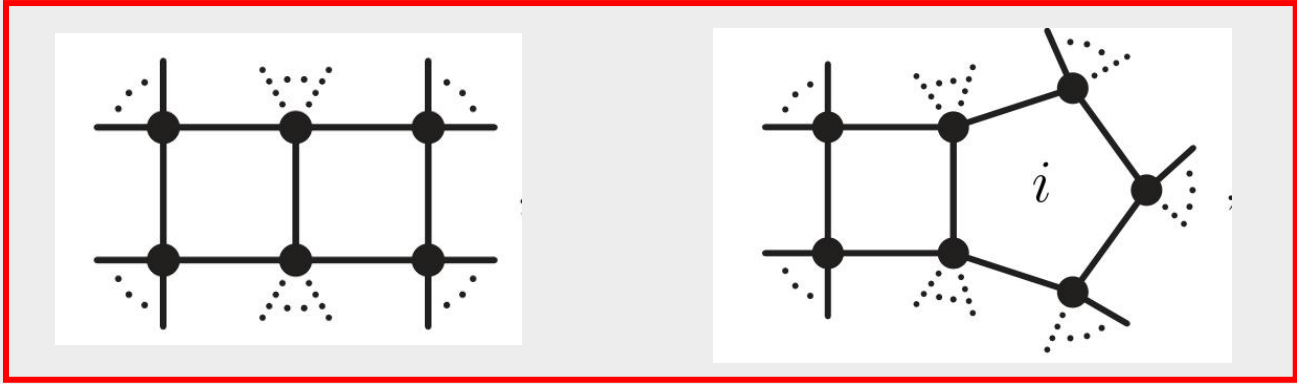
Schubert problem!



$$LS \propto \frac{1}{\Delta_{i,j,k,l}}$$

Two-loop Leading Singularities [1505.05886 by J.Bourjaily et al]

- Three types: composite, complete, overcomplete

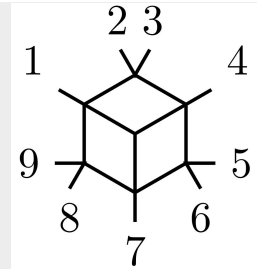
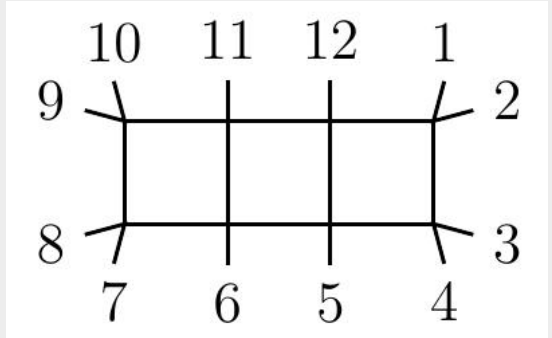
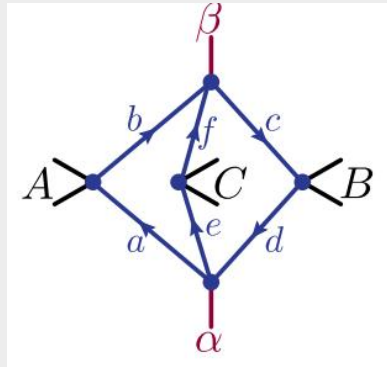
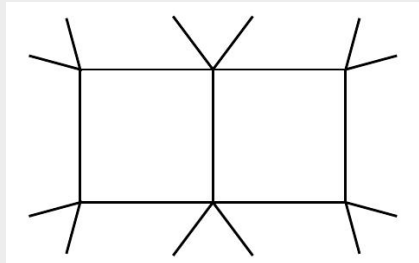
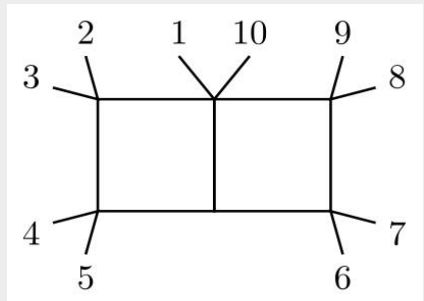
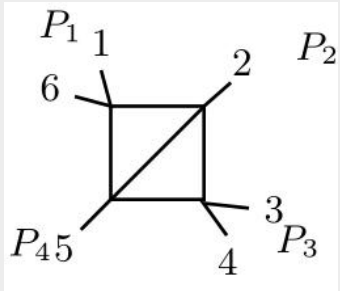


#propagator

$< 4L$

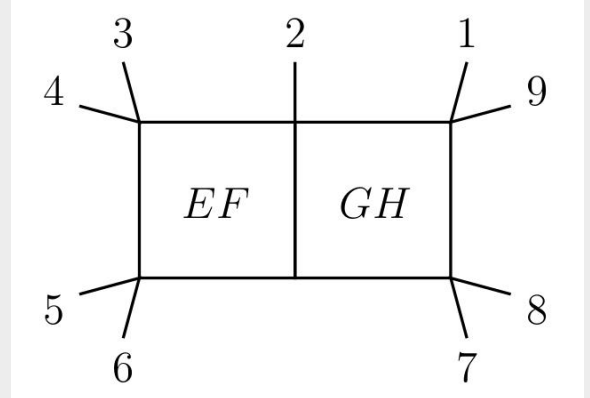
$= 4L$

$> 4L$



Composite LS

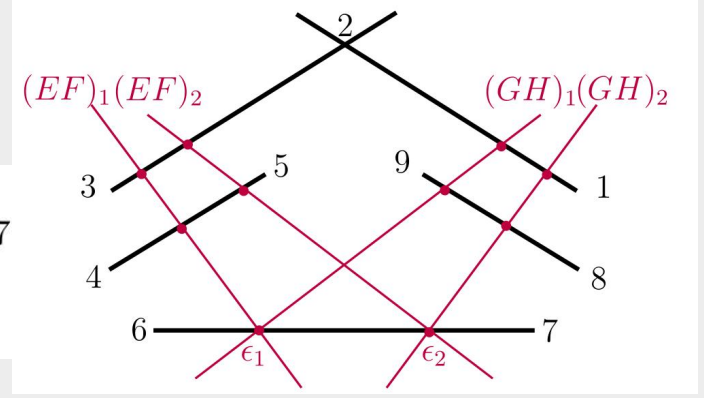
- # propagators < 4L: L loop momenta cannot be fixed by on-shell conditions (Jacobian)



Parametrizing two loop momenta with GL(2) gauge-fixing

$$E = \alpha_1 Z_5 + \beta_1 Z_6 + Z_4, \quad F = \alpha_2 Z_5 + \beta_2 Z_6 + Z_7$$

$$G = \gamma_1 Z_7 + \delta_1 Z_8 + Z_9, \quad H = \gamma_2 Z_7 + \delta_2 Z_8 + Z_6$$



Solving 7 on-shell conditions (linear in parameters)

$$\langle 7 (89) \cap (612) (645) \cap (623) \rangle - \langle 6 (89) \cap (712) (745) \cap (723) \rangle \gamma_2^2$$

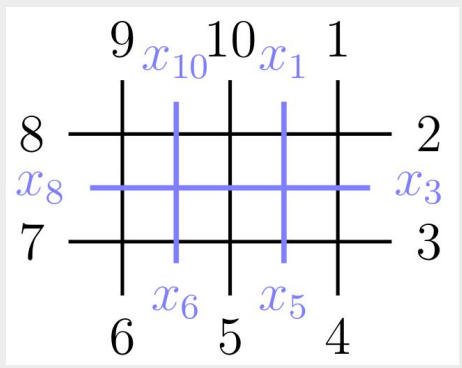
$$+ (\langle 6 (45) \cap (236) (789) \cap (712) \rangle - \langle 7 (45) \cap (237) (689) \cap (612) \rangle) \gamma_2$$

Jacobian factor

Discriminant of Jacobian: Δ_9 (two-loop square root)

$$\mathcal{J} = 0$$

$$LS \propto \frac{1}{\Delta_9}$$



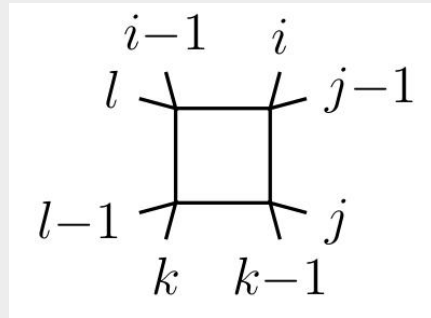
$\mathcal{J} = \sqrt{P(x)}$, with $P(x)$ a quartic curve

Elliptic function!

Elliptic LS: [2012.14438] by J.Bourjaily et al.

3. Intersections & letters from Schubert analysis [N. Arkani-Hamed, 21', QY 2203.16113]

After lines are determined by kinematics & LS, **intersections** are produced, from which we can construct **symbol letters**.



$$= \int_{AB} \frac{\langle i-1ik-1k \rangle \langle j-1jl-1l \rangle}{\langle ABi-1i \rangle \langle ABj-1j \rangle \langle ABk-1k \rangle \langle ABl-1l \rangle}$$

$$\Delta_{i,j,k,l} = \sqrt{(1-u-v)^2 - 4uv}$$

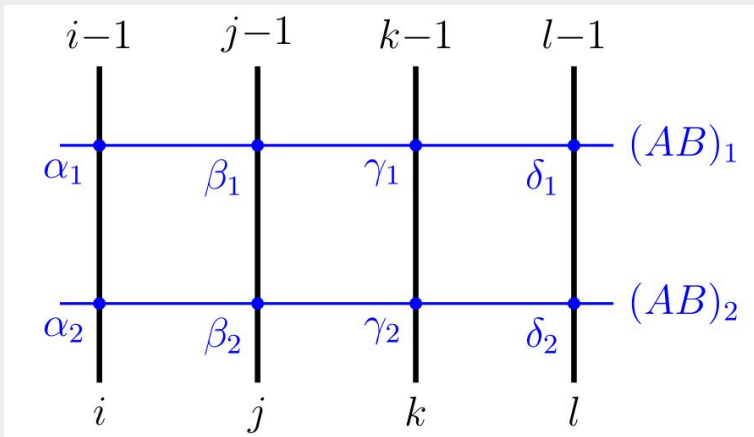
$$\frac{1}{2\Delta_{i,j,k,l}} \left(v \otimes \frac{z_{i,j,k,l}}{\bar{z}_{i,j,k,l}} + u \otimes \frac{1 - \bar{z}_{i,j,k,l}}{1 - z_{i,j,k,l}} \right)$$

One-loop square root

$$z_{i,j,k,l} \bar{z}_{i,j,k,l} = u, \quad (1 - z_{i,j,k,l})(1 - \bar{z}_{i,j,k,l}) = v.$$

$$u = \frac{\langle i-1ij-1j \rangle \langle k-1kl-1l \rangle}{\langle i-1ik-1k \rangle \langle j-1jl-1l \rangle}, \quad v = \frac{\langle i-1il-1l \rangle \langle j-1jk-1k \rangle}{\langle i-1ik-1k \rangle \langle j-1jl-1l \rangle}$$

- From internal lines of the box, we consider the **cross-ratios** formed by the intersections



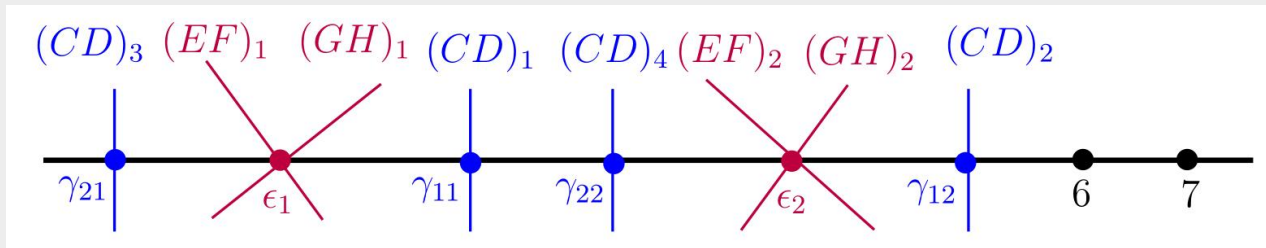
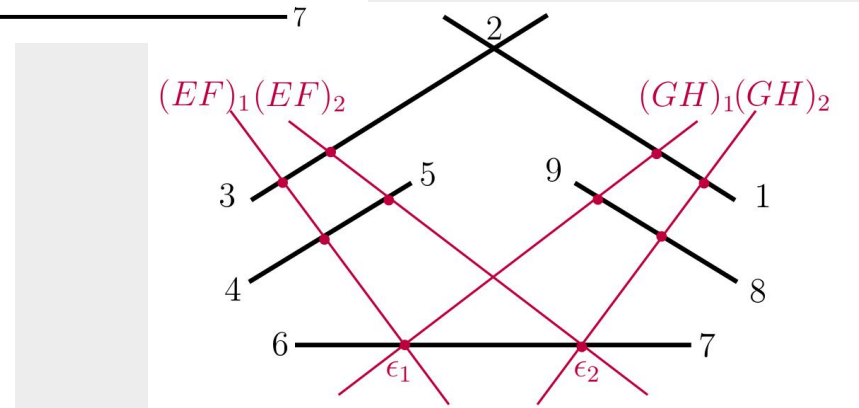
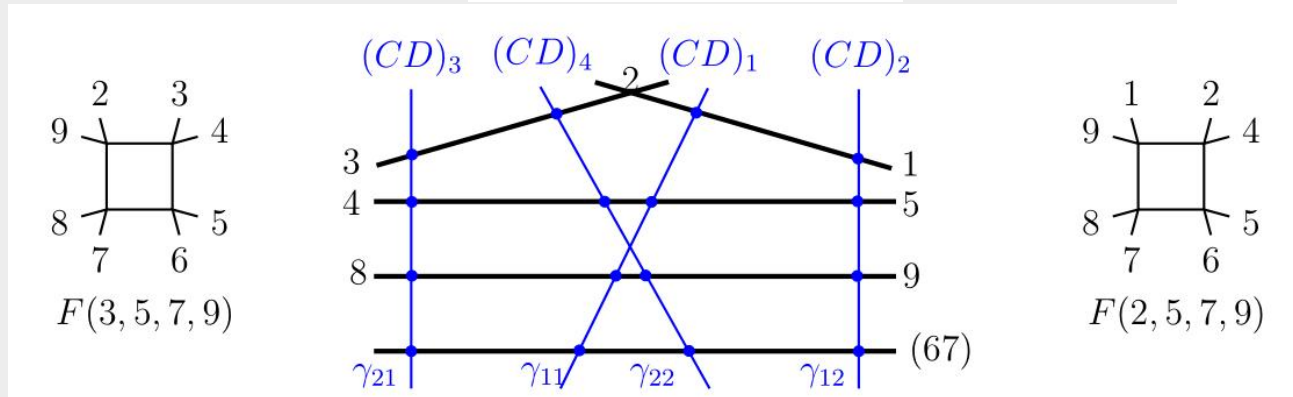
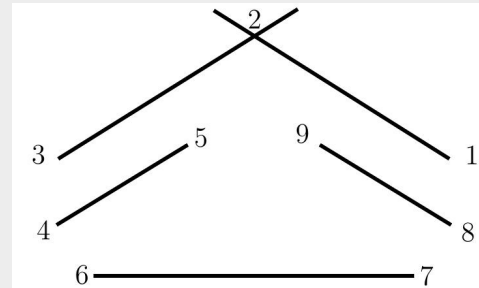
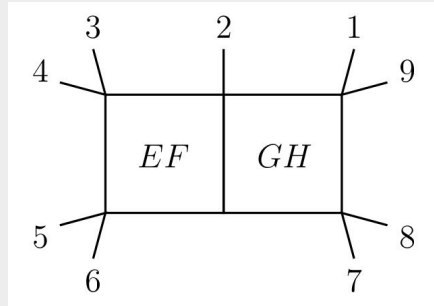
$$\frac{(\alpha_1, \beta_1)(\gamma_1, \delta_1)}{(\alpha_1, \gamma_1)(\beta_1, \delta_1)} = z_{i,j,k,l}, \quad \frac{(\alpha_1, \delta_1)(\gamma_1, \beta_1)}{(\alpha_1, \gamma_1)(\beta_1, \delta_1)} = 1 - z_{i,j,k,l}$$

$$\frac{(X_1, X_3)(X_2, X_4)}{(X_1, X_4)(X_2, X_3)} := \frac{\langle X_1 X_3 I \rangle \langle X_2 X_4 I \rangle}{\langle X_1 X_4 I \rangle \langle X_2 X_3 I \rangle}$$

$$\frac{(\alpha_2, \beta_2)(\gamma_2, \delta_2)}{(\alpha_2, \gamma_2)(\beta_2, \delta_2)} = \bar{z}_{i,j,k,l}, \quad \frac{(\alpha_2, \delta_2)(\gamma_2, \beta_2)}{(\alpha_2, \gamma_2)(\beta_2, \delta_2)} = 1 - \bar{z}_{i,j,k,l}$$

I: arbitrary reference line

Two-loop example 1: 9pt double-box integral



$$\frac{(\gamma_{11}, \gamma_{21})(\gamma_{12}, \gamma_{22})}{(\gamma_{11}, \gamma_{22})(\gamma_{21}, \gamma_{12})} = \frac{(z_{2,5,7,9} - z_{3,5,7,9})(\bar{z}_{2,5,7,9} - \bar{z}_{3,5,7,9})}{(z_{2,5,7,9} - \bar{z}_{3,5,7,9})(\bar{z}_{2,5,7,9} - z_{3,5,7,9})}$$

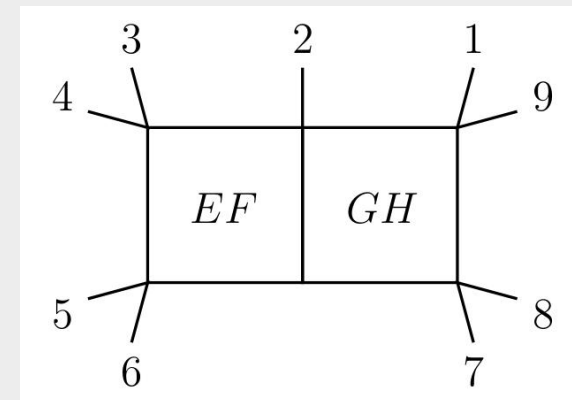
$$\frac{(\epsilon_1, \gamma_{11})(\epsilon_2, \gamma_{12})}{(\epsilon_2, \gamma_{11})(\epsilon_1, \gamma_{12})} = \frac{(1 + az_{2,5,7,9})(1 + b\bar{z}_{2,5,7,9})}{(1 + a\bar{z}_{2,5,7,9})(1 + bz_{2,5,7,9})}$$

- We call these configurations: combinations of Schubert problems on external lines

Two-loop example 1: 9pt double-box integral

9-point double-box integral (weight-4 MPL)

$$S(I_9) = \sum A_1 \otimes A_2 \otimes B \otimes C$$

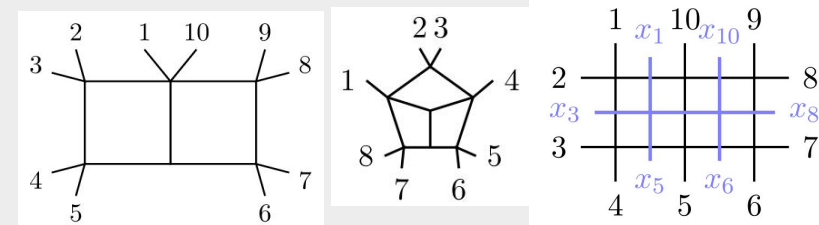


$A_1 \otimes A_2$ (first-two-entries) : letters from one-loop intersections on internal lines

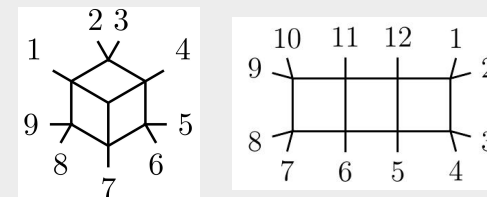
B (third entries): letters from one-loop intersections on external lines

C (last entries): letters from two-loop intersections combined with one-loop intersections

General structure for integrals with composite/complete LS



Guidance for symbology of higher loops/rigidity integrals



Two-loop example 2: 2-mass-easy box kinematics and its two-loop alphabet

[2207.13482,w.S.He et al]

$$D = 4 - 2\epsilon$$

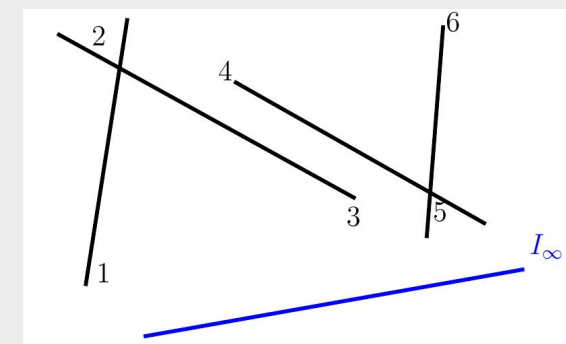
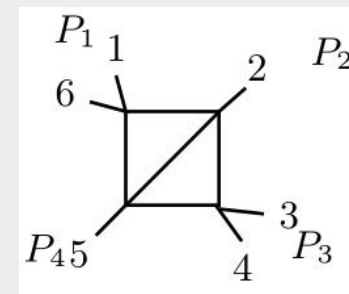
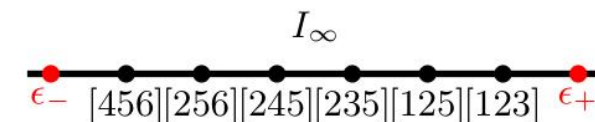
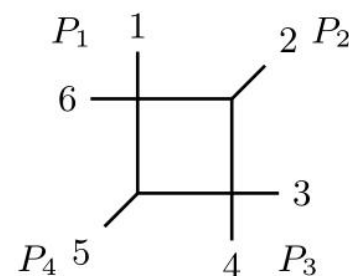
Black intersections: 4 one-loop triangles (A3)

Red intersections: LS of two-loop slashed-box

$$m_1^2, m_3^2, s, t \quad s - m_1^2, t - m_1^2, s - m_3^2, t - m_3^2, m_1^2 + m_3^2 - s - t, st - m_1^2 m_3^2.$$

$$L_1 = \frac{s + t + \Delta_{nc}}{s + t - \Delta_{nc}}, L_2 = \frac{s - t + \Delta_{nc}}{s - t - \Delta_{nc}}, L_3 = \frac{-2m_1^2 + s + t + \Delta_{nc}}{-2m_1^2 + s + t - \Delta_{nc}}$$

$$\Delta_{nc} = \sqrt{(s + t)^2 - 4m_1^2 m_3^2}$$



Following this procedure, we recover two-loop alphabets for non-DCI kinematics case by case, by going through all possible integrals...

(boxes, zero-, one-mass pentagon) (selection)

Schubert analysis beyond MPL

e.g. 10-pt double-box integral [\[2106.14902 by M.Wilhelm et al\]](#)

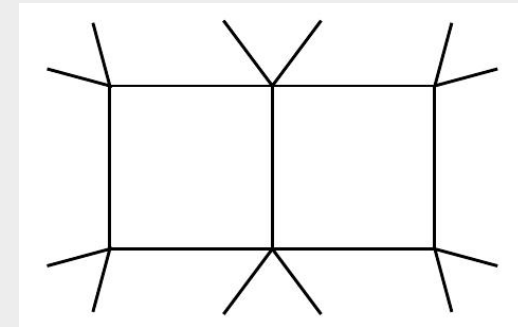
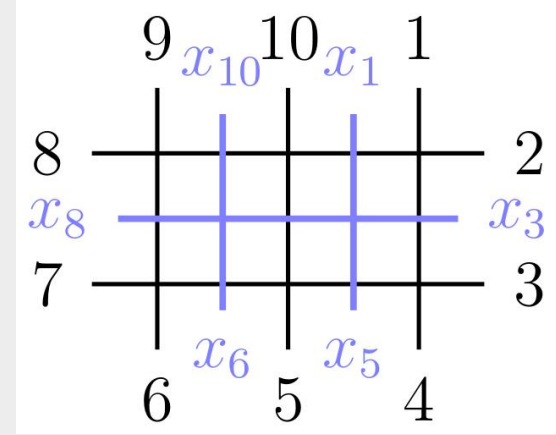
Elliptic “cross-ratios:”

$$\frac{1}{\omega_1} \int_a^b \frac{dx}{\sqrt{P(x)}}$$

Elliptic
curve in \mathbb{P}^3

a & b: two intersections from
another Schubert problem
 ω_1 : elliptic period

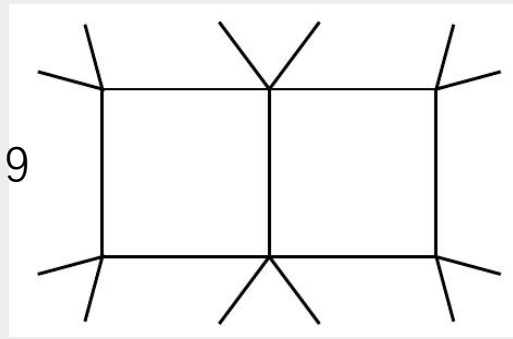
$$S(I_{10}) = \sum A_1 \otimes A_2 \otimes B \otimes D$$



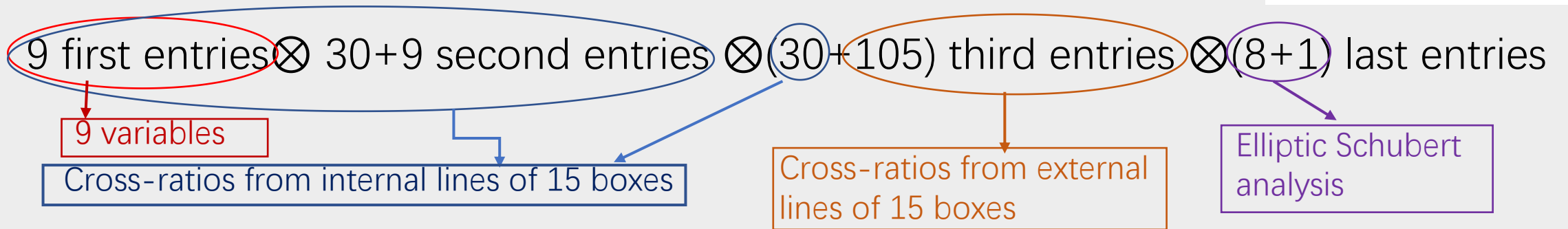
Successfully reproduce the last entries of 10-point db integrals and
give a predictions of 12-point db integrals → bootstrapping

Bootstrap 12-point double-box integral [w.M.Wilhelm et al 2212.09762]

d=9



- Alphabet from Schubert analysis:

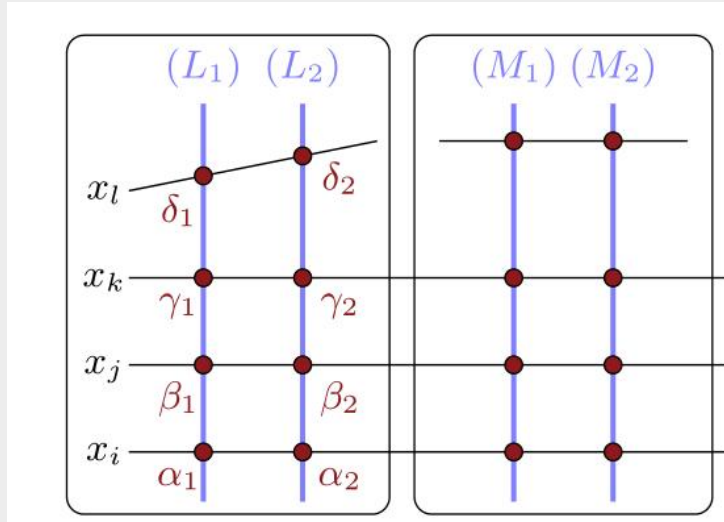


Imposing **integrability** and **Steinmann relations** (of first-two entries), determining one single symbol \Rightarrow **Symbol of the Result!**

Finally: 9 first entries \otimes 39 second entries \otimes 104 third entries \otimes (8+1) last entries

Selection Rule

- Truth: 1. for two boxes sharing three external lines, cross-ratios from each external line (A_{-1}) are **the same**
- **2. These cross-ratios account for all required third entries**



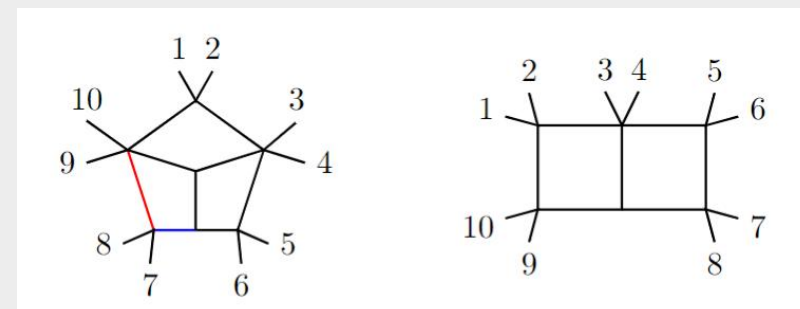
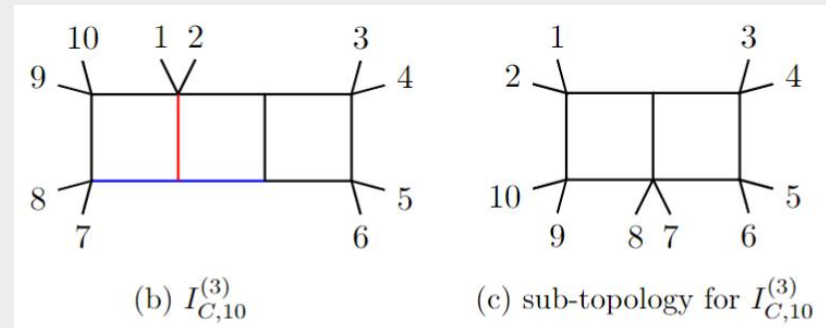
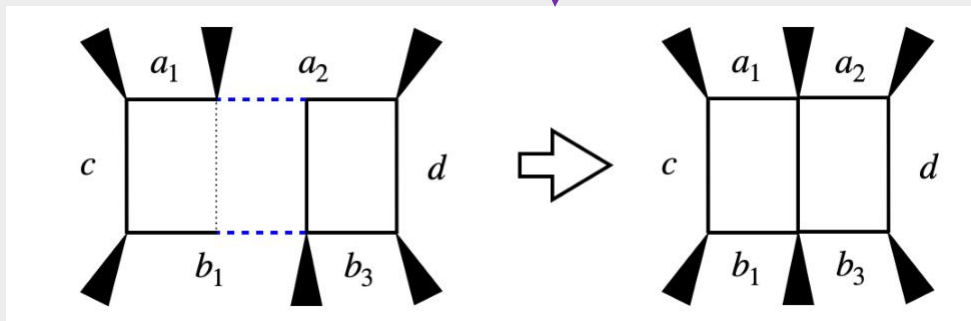
$$L([ad], [bc]) := \int_a^d d \log \frac{(xc)}{(xb)} = \log \mathcal{U} = L([bc], [ad])$$

- Rules: we only consider those combinations of n Schubert problems, $A_{\{2n-3\}}$ configurations from each shared external line are equivalent
- (checked for all known results/ one-loop letters of non-DCI kinematics...)

May still be redundant

Higher-loop analysis

- Generally speaking, for an L-loop integral with composite LS, its alphabet consists of:
 1. letters from all (L-1)-loop **sub-topologies**
 2. letters from all (legal) combinations of two Schubert problems from two sub-topologies (**after selections**)
 3. letters from all (legal) combinations of Schubert problem from LS with problems from one sub-topology (after selections, **last entries**)



Exercise: a three-loop example

Weight-6
pure MPL

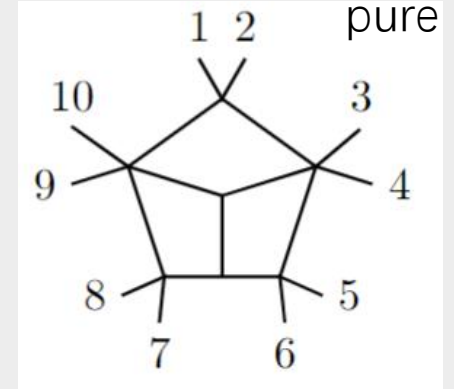
Let us construct its alphabet by Schubert analysis:

1. One-loop letters:

It has five one-loop sub-topologies (five 4m boxes)

5 **1st-entries** + 15 **2nd entries** from internal lines (u,v,z)

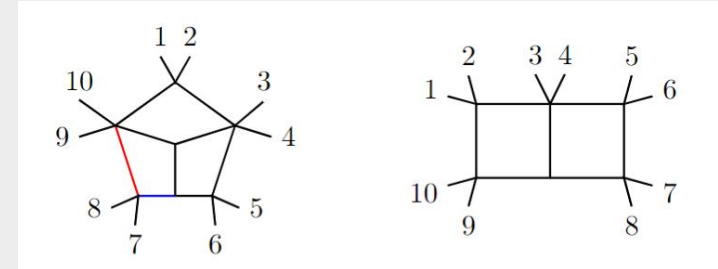
By combinations: 5 Deltas, 1 Gram and 10 mixed-algebraic letters from external lines (**3rd entries**)



2. Two-loop letters:

It has 2 two-loop sub-topologies (10-point double-box)

Each contributes 5 new algebraic letters and 1 Delta (**4th&5th entries**)



3. Three-loop letters (**last entries**)

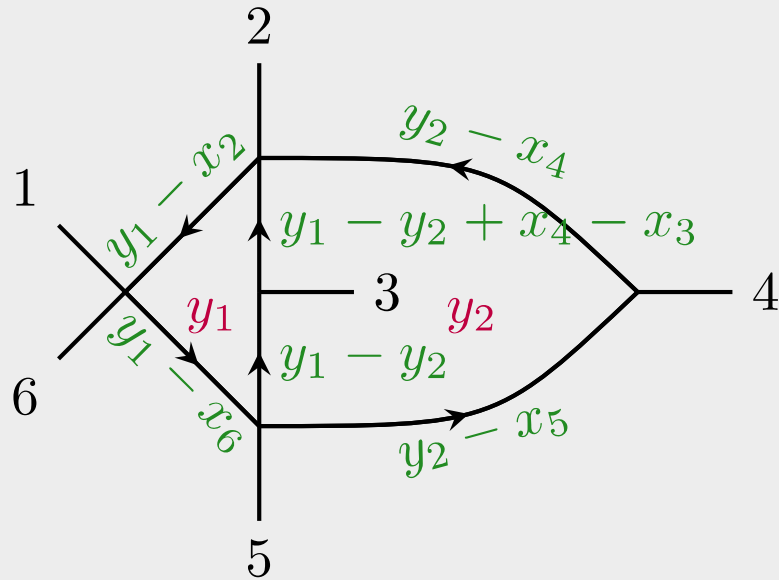
Five algebraic letters from combinations of three-loop intersections with one-loop intersections (and 1 delta)

Other cross-ratios have been ruled out by our selection

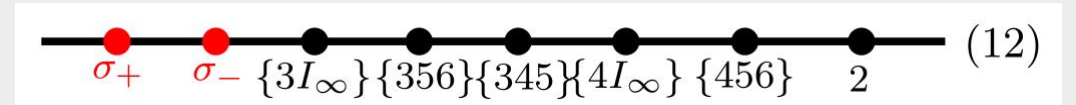
Integrability: only five possible functions
physical limits imposed:
determined!

Non-planar Schubert analysis

- By expanding propagators on $(y-x)^2$, we can rewrite the integrand by twistors



[2107.14180] by
S. Abreu et al



$$\Sigma_5 = (s_{12}s_{15} - s_{12}s_{23} - s_{15}s_{45} + s_{34}s_{45} + s_{23}s_{34})^2 - 4s_{23}s_{34}s_{45}(s_{34} - s_{12} - s_{15}).$$

$$\begin{aligned} (y_1 - y_2 + x_4 - x_3)^2 &= (y_1 - y_2)^2 + (x_4 - x_3)^2 \\ &\quad + (y_1 - x_3)^2 - (y_1 - x_4)^2 \\ &\quad - (y_2 - x_4)^2 + (y_2 - x_3)^2 \end{aligned}$$

Constructing cross-ratios by two red points and any two black points, we recover the 5-dim space for algebraic letters with Σ_5

$$\langle AB12 \rangle = \langle AB56 \rangle = \langle CD34 \rangle = \langle CD45 \rangle = \langle ABCD \rangle = 0 \ \& \ \langle ABI_\infty \rangle \langle CD \bar{3} \cap (3I_\infty) \rangle - (AB \leftrightarrow CD) = 0$$

Summary

- Twistor geometries provide a geometric explanations for the symbol letters of L-loop integrals with massless propagators in QCD, and especially for $N=4$ SYM
- Through this way, we can predict symbol letters for individual integrals quite precisely, and finally bootstrap them instead of direct computations
- This method can be generalized to elliptic cases, and proves to be useful for 10-/12-pt double-box integral
- More precise description for Integrals with overcomplete LSs ? non-planar integrals? Massive propagators? $D=6$?...
- Relation to Landau equations? Cut integrals? Why it works?
- K3 integrals?...

Thanks!