

# 4D differential equations for the leading divergences of dimensionally-regulated loop integrals

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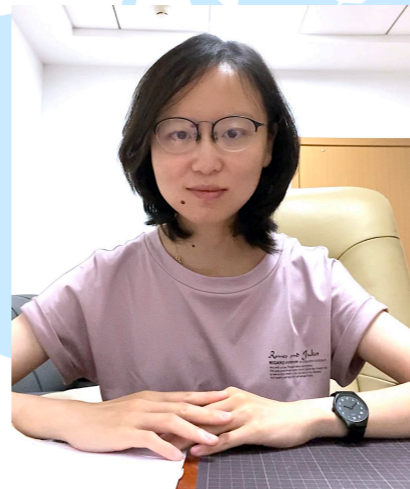


This talk is about our recent project JHEP 03(2023)162 with

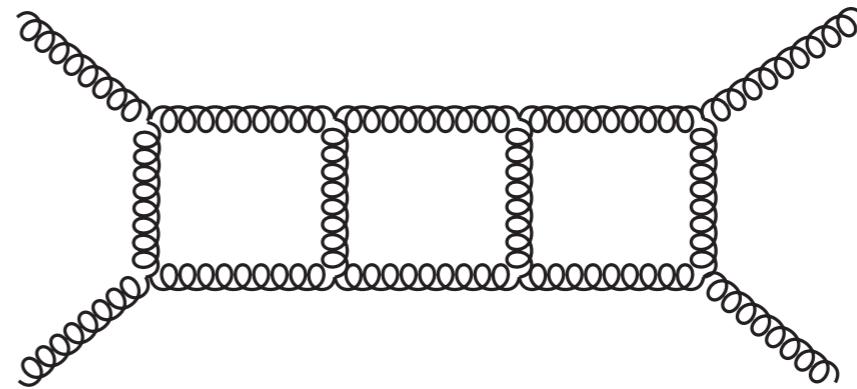
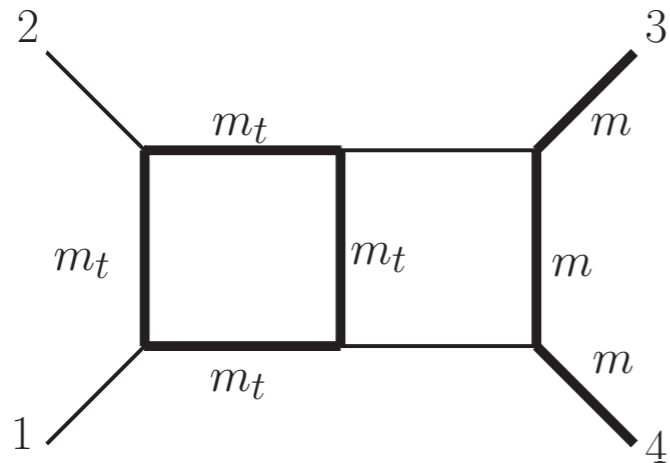
Johannes Henn, *Max-Planck Institute for Physics, Munich*

Kai Yan (颜开), *Shanghai Jiao Tong University*

Yang Zhang (张扬), *University of Science and Technology of China*



# Feynman integral evaluation



Multi-loop Feynman integrals are the **hardcore** objects for a perturbative QFT computation

- Important for high-energy phenomenology
- Theoretically important

# Mainstream Analytic Methods

## Integration by parts (IBP)

[Analytic Tool for Feynman Integrals] *Smirnov*

For the integrals  $F(a) = \int \frac{d^d k}{(k^2 - m^2)^a}$

The IBP identity  $\int d^d k \frac{\partial}{\partial k_\mu} \cdot k_\mu \frac{1}{(k^2 - m^2)^a} = 0$

$$(d - 2a)F(a) - 2am^2 F(a + 1) = 0$$

Use this set of relations between Feynman integrals in order to solve the reduction problem



# Mainstream Analytic Methods

Differential equation for Feynman integrals

$$d = 4 - 2\epsilon \quad \vec{f}(\bar{x}, \epsilon) \text{ the set of convenient basis integral}$$

**Canonical** Differential equation

$$d \vec{f}(\bar{x}, \epsilon) = \epsilon [d\tilde{A}(\bar{x})] \vec{f}(\bar{x}, \epsilon)$$

$$d\tilde{A}(\bar{x}) = \sum_k a_k d\log(\alpha_k) \quad (\text{Henn 2013})$$

certain constant matrices



letters  
alphabet

$$\vec{f}(\bar{x}) = P \exp \left( \epsilon \int d\tilde{A} \right) \vec{f}(\bar{x}_0)$$

Chen's (陈国才) iterated integrals, to all orders in the  $\epsilon$  expansion.

It may take a lot of efforts to find a **complete** UT basis...

# Motivation

For many applications, it is **not necessary** to evaluate all MIs or all order of  $\epsilon$ .

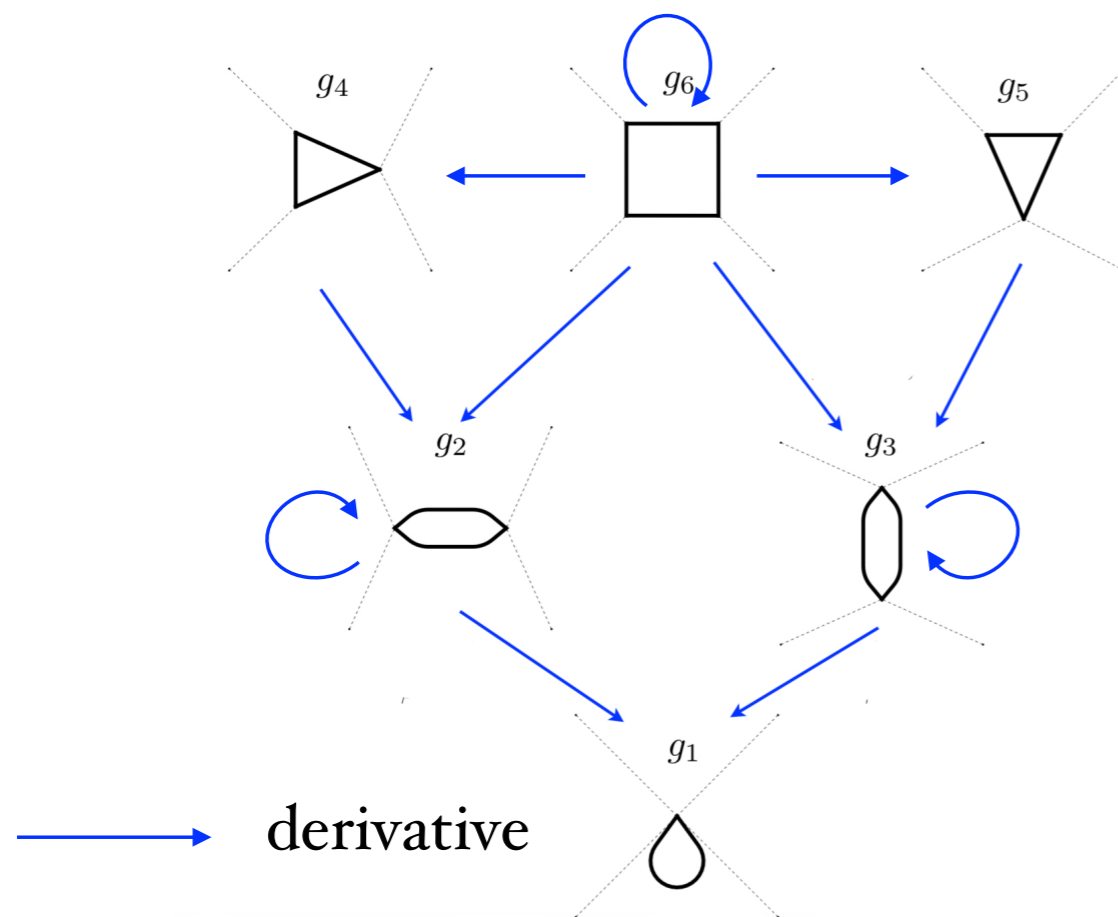
Can we simplify the DE in these cases?

# Motivation

For many applications, it is **not necessary** to evaluate all MIs or all order of  $\epsilon$ .

Can we simplify the DE in these cases?

Example: canonical DE for Finite integrals



Simon Caron-Huot, Johannes Henn 2014

$$g_1 = 2m^2 G[0, 0, 0, 3] / \epsilon^2$$

$$g_2 = -\sqrt{s(s - 4m^2)} G[1, 0, 2, 0] / \epsilon$$

$$g_3 = -\sqrt{t(t - 4m^2)} G[0, 1, 0, 2] / \epsilon$$

$$g_4 = s G[1, 1, 1, 0]$$

$$g_5 = t G[1, 1, 0, 1]$$

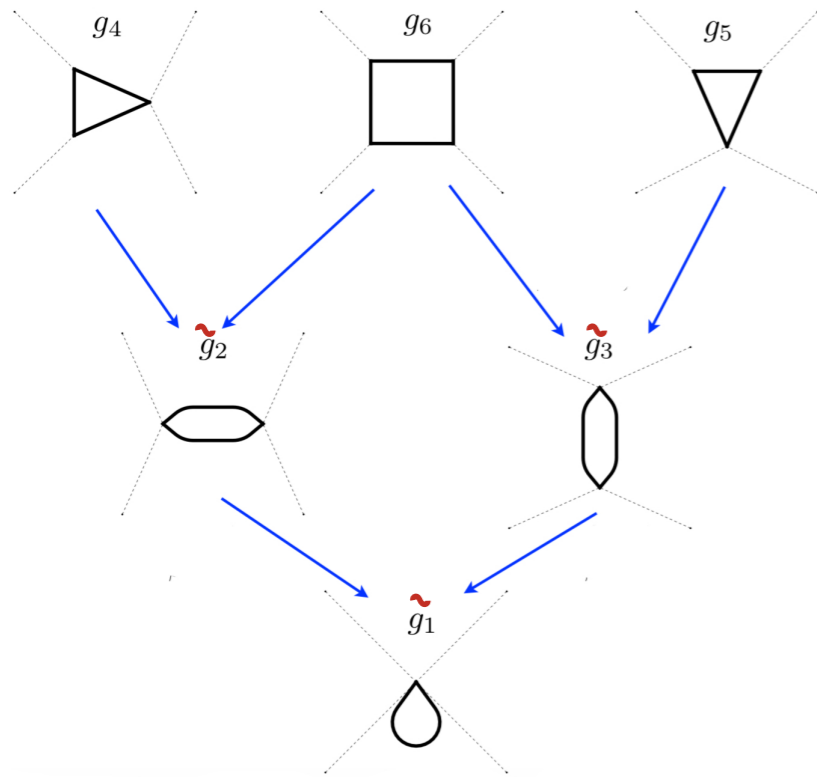
$$g_6 = \frac{1}{2} \sqrt{st(st - 4m^2(s + t))} G[1, 1, 1, 1]$$

$$d\vec{g} = \epsilon(dA)\vec{g}$$

# canonical DE

$$\mathcal{E} \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ \log\left(\frac{\beta_u-1}{\beta_u+1}\right) & \log\left(\frac{u}{1+u}\right) & 0 & 0 & 0 & 0 \\ \log\left(\frac{\beta_v-1}{\beta_v+1}\right) & 0 & \log\left(\frac{v}{1+v}\right) & 0 & 0 & 0 \\ 0 & -\log\left(\frac{\beta_u-1}{\beta_u+1}\right) & 0 & 0 & 0 & 0 \\ 0 & 0 & -\log\left(\frac{\beta_v-1}{\beta_v+1}\right) & 0 & 0 & 0 \\ 0 & \log\left(\frac{\beta_{uv}-\beta_u}{\beta_{uv}+\beta_u}\right) & \log\left(\frac{\beta_{uv}-\beta_v}{\beta_{uv}+\beta_v}\right) & \log\left(\frac{\beta_{uv}-1}{\beta_{uv}+1}\right) & \log\left(\frac{\beta_{uv}-1}{\beta_{uv}+1}\right) & \log\left(\frac{u+v}{1+u+v}\right) \end{pmatrix}$$

# canonical DE for the finite integrals



$$\tilde{g}_1 = 2m^2 G[0, 0, 0, 3] / \epsilon^2$$

$$\tilde{g}_2 = -\sqrt{s(s-4m^2)} G[1, 0, 2, 0] / \epsilon$$

$$\tilde{g}_3 = -\sqrt{t(t-4m^2)} G[0, 1, 0, 2] / \epsilon$$

$$g_4 = sG[1, 1, 1, 0]$$

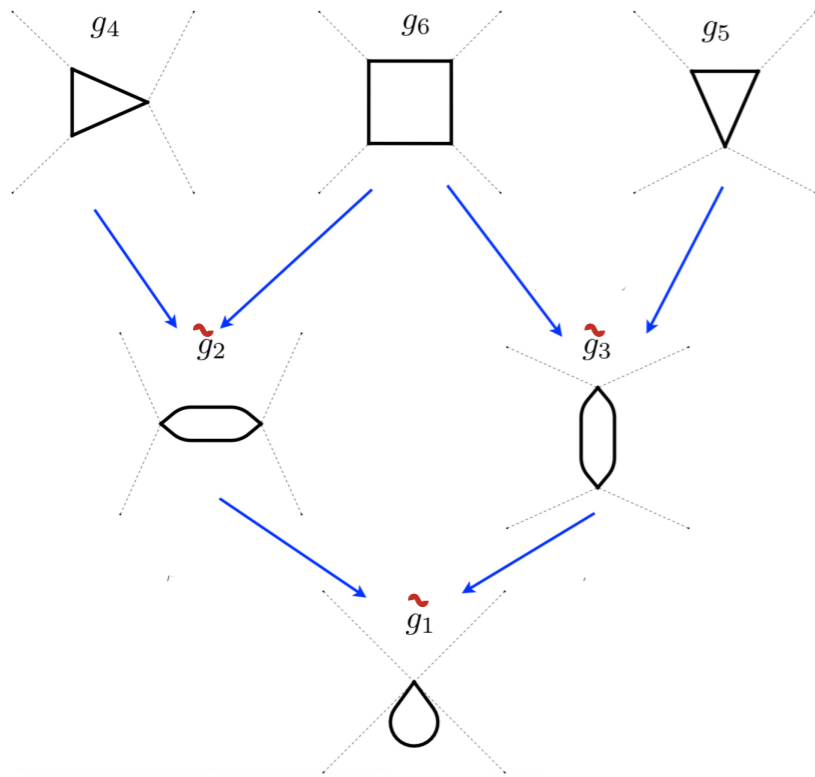
$$g_5 = tG[1, 1, 0, 1]$$

$$g_6 = \frac{1}{2} \sqrt{st(st-4m^2(s+t))} G[1, 1, 1, 1]$$

# canonical DE

$$\mathfrak{E} \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \log\left(\frac{\beta_u-1}{\beta_u+1}\right) & \log\left(\frac{u}{1+u}\right) & 0 & 0 & 0 & 0 & 0 \\ \log\left(\frac{\beta_v-1}{\beta_v+1}\right) & 0 & \log\left(\frac{v}{1+v}\right) & 0 & 0 & 0 & 0 \\ 0 & -\log\left(\frac{\beta_u-1}{\beta_u+1}\right) & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -\log\left(\frac{\beta_v-1}{\beta_v+1}\right) & 0 & 0 & 0 & 0 \\ 0 & \log\left(\frac{\beta_{uv}-\beta_u}{\beta_{uv}+\beta_u}\right) & \log\left(\frac{\beta_{uv}-\beta_v}{\beta_{uv}+\beta_v}\right) & \log\left(\frac{\beta_{uv}-1}{\beta_{uv}+1}\right) & \log\left(\frac{\beta_{uv}-1}{\beta_{uv}+1}\right) & \log\left(\frac{\beta_{uv}-1}{\beta_{uv}+1}\right) & \log\left(\frac{u+v}{1+u+v}\right) \end{pmatrix}$$

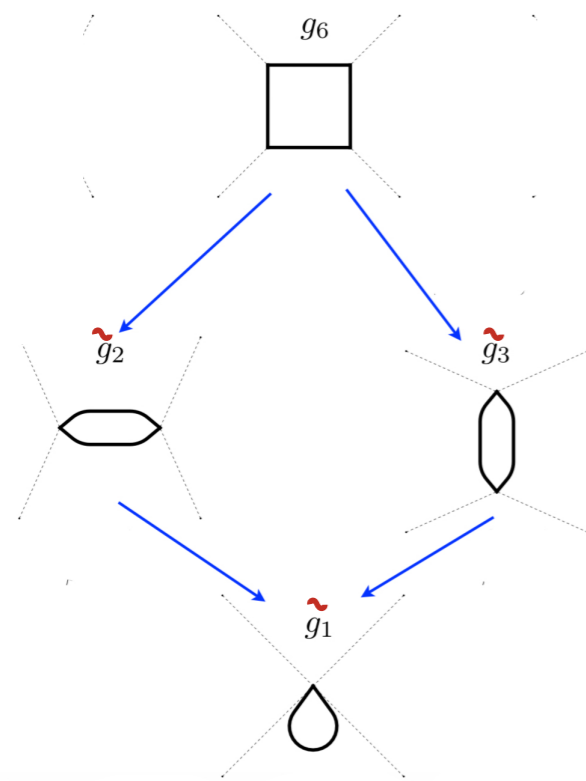
# canonical DE for the finite integrals



# canonical DE

$$\mathcal{E} = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \log\left(\frac{\beta_u-1}{\beta_u+1}\right) & \log\left(\frac{u}{1+u}\right) & 0 & 0 & 0 & 0 & 0 \\ \log\left(\frac{\beta_v-1}{\beta_v+1}\right) & 0 & \log\left(\frac{v}{1+v}\right) & 0 & 0 & 0 & 0 \\ 0 & -\log\left(\frac{\beta_u-1}{\beta_u+1}\right) & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -\log\left(\frac{\beta_v-1}{\beta_v+1}\right) & 0 & 0 & 0 & 0 \\ 0 & \log\left(\frac{\beta_{uv}-\beta_u}{\beta_{uv}+\beta_u}\right) & \log\left(\frac{\beta_{uv}-\beta_v}{\beta_{uv}+\beta_v}\right) & \log\left(\frac{\beta_{uv}-1}{\beta_{uv}+1}\right) & \log\left(\frac{\beta_{uv}-1}{\beta_{uv}+1}\right) & \log\left(\frac{u+v}{1+u+v}\right) & 0 \end{pmatrix}$$

# canonical DE for the finite integrals



6 → 4

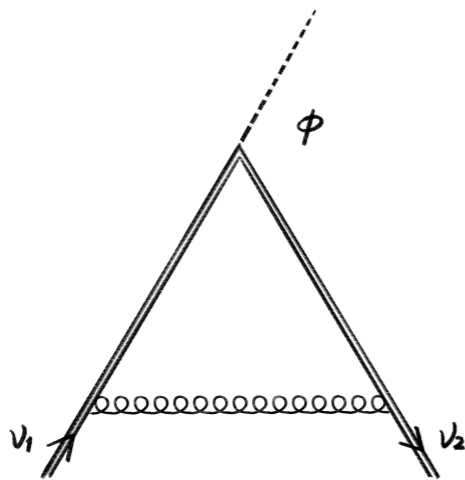
$$\tilde{A} = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ \log\left(\frac{\beta_u-1}{\beta_u+1}\right) & 0 & 0 & 0 & 0 & 0 \\ \log\left(\frac{\beta_v-1}{\beta_v+1}\right) & 0 & 0 & 0 & 0 & 0 \\ 0 & \log\left(\frac{\beta_u-1}{\beta_u+1}\right) & 0 & 0 & 0 & 0 \\ 0 & 0 & \log\left(\frac{\beta_v-1}{\beta_v+1}\right) & 0 & 0 & 0 \\ 0 & \log\left(\frac{\beta_{uv}-\beta_u}{\beta_{uv}+\beta_u}\right) & \log\left(\frac{\beta_{uv}-\beta_v}{\beta_{uv}+\beta_v}\right) & 0 & 0 & 0 \end{pmatrix}$$

All diagonal elements vanish  
and the DE can be solved recursively.

# Why divergent Feynman Integral

Consider the renormalization of quantum field theory,  
 $\beta$  and  $\gamma$  function are simply related to the coefficients of  
the divergent logarithms of counterterm

In HQET, the coefficient of the  $1/\epsilon$  pole is  
the angle-dependent cusp anomalous dimension



.....

$$I^{(1)}(\phi, \epsilon) \sim \frac{1}{\epsilon} \phi \cot \phi + \mathcal{O}(\epsilon^0)$$

$$\text{where } \cos \phi = \frac{v_1 \cdot v_2}{\sqrt{v_1^2 v_2^2}}$$

- have to deal with divergent Feynman integrals (UV/IR)
- have no general method to compute them using IBPs

Can we simplify the canonical DE of **divergent** integrals?

If we only need the divergent part...

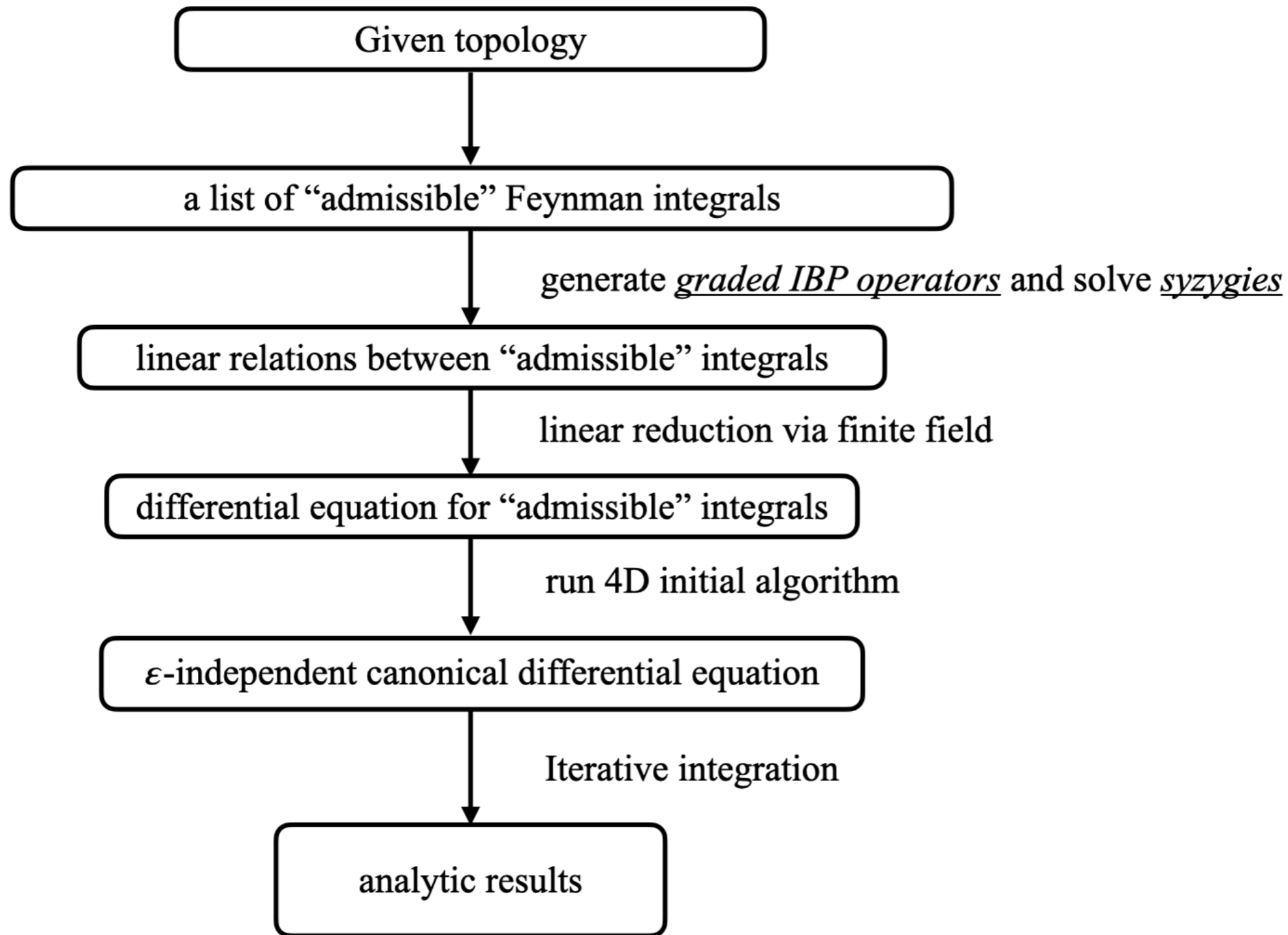


# Our New Method

Canonical DE of the divergent part

- Generate admissible integrals
- Search for an IBP system with only admissible integrals  
(with graded syzygy approach/cross loop order relations)
- 4D initial algorithm  
the original initial algorithm is proposed by Dlapa, Henn, Kai (2020)





# Example: 3-loop HQET

Admissible: only one region of divergence, overall UV

The coefficient of the  $1/\epsilon$  pole is the angle-dependent cusp anomalous dimension

$$D_1 = -2k_1 v_1 + \delta, \quad D_2 = -2k_2 v_1 + \delta, \quad D_3 = -2k_3 v_1 + \delta$$

$$D_4 = -2k_1 v_2 + \delta, \quad D_5 = -2k_2 v_2 + \delta, \quad D_6 = -2k_3 v_2 + \delta$$

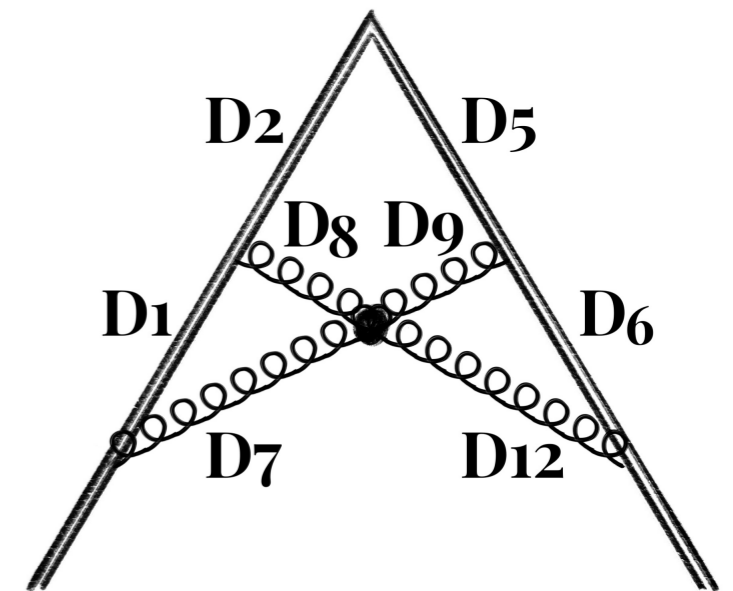
$$D_7 = -k_1^2, \quad D_8 = -(k_1 - k_2)^2, \quad D_9 = -(k_2 - k_3)^2$$

$$D_{10} = -(k_1 - k_3)^2, \quad D_{11} = -k_2^2, \quad D_{12} = -k_3^2$$

power counting zero

free of UV sub-divergences

free of soft divergence



divergence of  $k_3$

divergence of  $k_3$

admissible  $G[2, 1, \underline{-1}, 0, 1, 1, 1, 1, 1, 0, 0, 1] - G[2, 1, \underline{-1}, 0, 1, 1, 1, 1, 1, 1, 1, 0, 0]$

$$\frac{1}{k_3^2} - \frac{1}{(k_1 - k_3)^2} = \frac{k_1^2}{k_3^2 (k_1 - k_3)^2} - \frac{2k_1 k_3}{k_3^2 (k_1 - k_3)^2}$$

reducible SP

$$D_{13} \equiv k_1 \cdot k_2 \quad D_{14} \equiv k_1 \cdot k_3, \quad D_{15} \equiv k_2 \cdot k_3$$

0 superficial degree of divergence

$$a_1 + a_2 + a_3 + a_4 + a_5 + a_6 + 2a_7 + 2a_8 + 2a_9 + 2a_{10} \\ + 2a_{11} + 2a_{12} + 2a_{13} + 2a_{14} + 2a_{15} = 12$$

free of sub-divergences

$$a_1 + a_4 + 2a_7 + 2a_8 + 2a_{10} + a_{13} + a_{14} > 4,$$

$$a_2 + a_5 + 2a_8 + 2a_9 + 2a_{11} + a_{13} + a_{15} > 4,$$

$$a_3 + a_6 + 2a_9 + 2a_{10} + 2a_{12} + a_{14} + a_{15} > 4,$$

$$a_1 + a_2 + a_4 + a_5 + 2a_7 + 2a_8 + 2a_9 + 2a_{10} + a_{11} + 2a_{13} + a_{14} + a_{15} > 8,$$

$$a_1 + a_2 + a_4 + a_6 + 2a_7 + 2a_8 + 2a_9 + 2a_{10} + 2a_{12} + a_{13} + 2a_{14} + a_{15} > 8,$$

$$a_2 + a_3 + a_5 + a_6 + 2a_8 + 2a_9 + 2a_{10} + 2a_{11} + a_{12} + a_{13} + a_{14} + 2a_{15} > 8.$$

free of soft divergence

$$a_7, a_8, a_9, a_{10}, a_{11}, a_{12} < 2$$

it's easy to list admissible integrals by simple computer program

# IBP operator via Graded Syzygy

Graded ensure overall UV divergence, control UV sub-divergence

$$\text{deg}_i(k_i) = 1, \text{deg}_i\left(\frac{\partial}{\partial k_i}\right) = -1$$

each term have the same degree

Solved by  
linear algebra

Syzygy make sure free of soft divergence

$$O D_k = g_k D_k \quad g_k \text{ is a polynomial}$$

---

$$O \mathcal{F} \equiv \frac{\partial}{\partial k_1^\mu} \left( (v_1^\mu (k_1 \cdot k_2) - k_2^\mu (k_1 \cdot v_1) + k_1^\mu (k_2 \cdot v_1)) \mathcal{F} \right) \\ + \frac{\partial}{\partial k_2^\mu} \left( v_1^\mu (k_2 \cdot k_2) \mathcal{F} \right) + \frac{\partial}{\partial k_3^\mu} \left( (v_1^\mu (k_2 \cdot k_3) + k_3^\mu (k_2 \cdot v_1) - k_2^\mu (k_3 \cdot v_1)) \mathcal{F} \right)$$

$$O D_7 = 2(d + 2)(k_2 \cdot v_1) D_7$$

$$\begin{aligned}
O\mathcal{F} &\equiv \frac{\partial}{\partial k_1^\mu} \left( (v_1^\mu(k_1 \cdot k_2) - k_2^\mu(k_1 \cdot v_1) + k_1^\mu(k_2 \cdot v_1)) \mathcal{F} \right) \\
&+ \frac{\partial}{\partial k_2^\mu} \left( v_1^\mu(k_2 \cdot k_2) \mathcal{F} \right) + \frac{\partial}{\partial k_3^\mu} \left( (v_1^\mu(k_2 \cdot k_3) + k_3^\mu(k_2 \cdot v_1) - k_2^\mu(k_3 \cdot v_1)) \mathcal{F} \right)
\end{aligned}$$

$$O G[1, 1, 0, 0, 2, 1, 1, 1, 1, 0, 0, 1]$$



$$\begin{aligned}
0 = & -2 \left( x + \frac{1}{x} \right) G(1, 1, 0, 0, 3, 1, 1, 1, 1, 0, -1, 1) - \frac{1}{2} G(1, 0, 0, 0, 2, 1, 1, 1, 1, 0, 0, 1) \\
& - \frac{1}{2} G(1, 1, -1, 0, 1, 2, 1, 1, 1, 0, 0, 1) - 2G(1, 2, 0, 0, 2, 1, 1, 1, 1, 0, -1, 1) \\
& + 2g[2, 1, 0, 0, 2, 1, 1, 1, 1, 0, 0, 1, -1, 0, 0] \\
& + \left( x + \frac{1}{x} \right) g[1, 1, 0, 0, 2, 2, 1, 1, 1, 0, 0, 1, 0, 0, -1]
\end{aligned}$$

IBP with **only**  
admissible integrals

$$\begin{aligned}
O\mathcal{F} &\equiv \frac{\partial}{\partial k_1^\mu} \left( (v_1^\mu(k_1 \cdot k_2) - k_2^\mu(k_1 \cdot v_1) + k_1^\mu(k_2 \cdot v_1)) \mathcal{F} \right) \\
&+ \frac{\partial}{\partial k_2^\mu} \left( v_1^\mu(k_2 \cdot k_2) \mathcal{F} \right) + \frac{\partial}{\partial k_3^\mu} \left( (v_1^\mu(k_2 \cdot k_3) + k_3^\mu(k_2 \cdot v_1) - k_2^\mu(k_3 \cdot v_1)) \mathcal{F} \right)
\end{aligned}$$

$$\frac{\partial}{\partial k_1^\mu} (v_1^\mu(k_1 \cdot k_2)) \quad G[1, 1, 0, 0, 2, 1, 1, 1, 1, 0, 0, 1]$$



$$-2 \frac{\partial}{\partial k_1^\mu} (v_1^\mu(k_1 \cdot v_1)) \quad g[2, 1, 0, 0, 2, 1, 1, 1, 1, 0, 0, 1, -1, 0, 0]$$

zero overall degree

admissible integral

Get IBP with **only admissible** integrals

**safely set  $\varepsilon$  and  $\delta$  to zero!**

# “4D” Initial Algorithm

“One UT integral rules all”

Dlapa, Henn, Kai (2020)

$(f_1, \dots, f_n)$  is a master integral basis, where  $f_1$  is UT but the others are not.  
 Try to find a UT basis  $(f_1, g_2 \dots g_n)$  which is UT.

$$\begin{pmatrix} f_1 \\ f_1^{[1]} \\ \dots \\ f_1^{[n-1]} \end{pmatrix} = \underset{\substack{\uparrow \\ \text{known}}}{\Psi} \begin{pmatrix} f_1 \\ f_2 \\ \dots \\ f_n \end{pmatrix} \qquad \begin{pmatrix} f_1 \\ f_1^{[1]} \\ \dots \\ f_1^{[n-1]} \end{pmatrix} = \underset{\substack{\uparrow \\ \text{unknown}}}{\Phi} \begin{pmatrix} f_1 \\ g_2 \\ \dots \\ g_n \end{pmatrix}$$

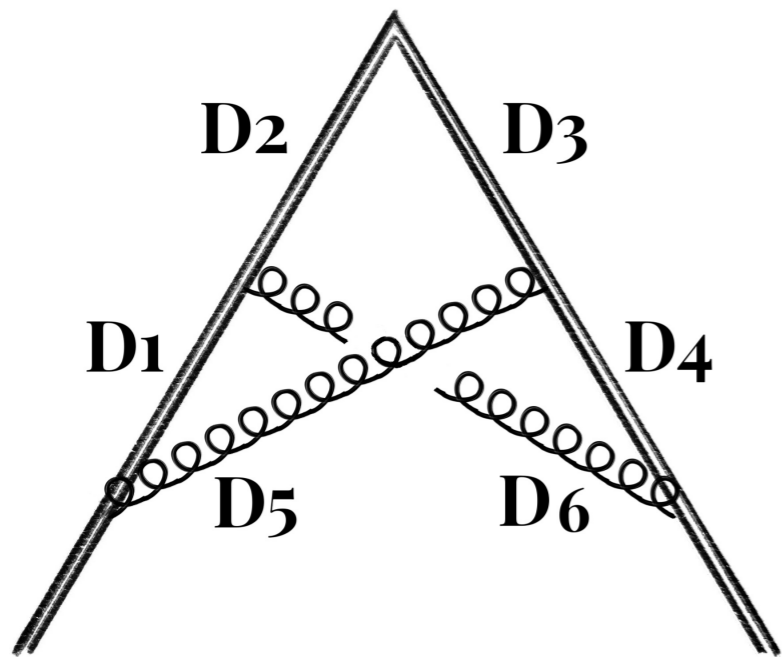
$\Phi$  is derived from a canonical DE with undermined coefficients

$$\sum_{i=1}^k a_i \partial_x \log(W_i) \qquad v_0 \Psi^{-1} \Phi = v_0, \quad v_0 = (1, 0 \dots 0)$$

For 4D, ansatz  $\{a_i\}$  is a strictly upper triangular matrix  
 with rational number entries



# Toy example: 2-loop HQET integral



$$2v_1 v_2 \equiv x + \frac{1}{x}, \quad v_1^2 = 1, \quad v_2^2 = 1$$

$$\mathbf{d} \equiv \frac{x^2 - 1}{2} \frac{d}{dx} \equiv \frac{(v_1 \cdot v_2) v_1^\mu - v_1^2 v_2^\mu}{\sqrt{v_1^2 v_2^2}} \frac{\partial}{\partial v_1^\mu}$$

$$T^{(2)}[a_1, \dots, a_7] = \int \frac{d^D k_1}{i\pi^{D/2}} \frac{d^D k_2}{i\pi^{D/2}} \frac{D_7^{-a_7}}{D_1^{a_1} \dots D_6^{a_6}}$$

$$D_1 = -2k_1 \cdot v_1 + \delta, \quad D_2 = -2(k_1 + k_2) \cdot v_1 + \delta, \quad D_3 = -2(k_1 + k_2) \cdot v_2 + \delta,$$

$$D_4 = -2k_2 \cdot v_2 + \delta, \quad D_5 = -k_1^2, \quad D_6 = -2k_2^2, \quad D_7 = k_1 \cdot k_2$$

reducible SP  $D_8 \equiv -2k_1 \cdot v_2 \quad D_9 \equiv -2k_2 \cdot v_1$

$$T^{(2)}[1, 1, 1, 1, 1, 1, 0] \text{ is admissible}$$

All its derivatives in  $x$  are admissible

# Toy example: 2-loop HQET integral

Admissible Integrals



Admissible IBP vectors

IBPs with **only admissible integrals**

$$A_1 = T^{(2)} [1, 1, 1, 1, 1, 1, 0], A_2 = t_{(2)} [2, 1, 1, 1, 1, 1, 0, -1, 0],$$

$$A_3 = t^{(2)} [1, 1, 1, 2, 1, 1, 0, 0, -1], B_1 = T^{(2)} [1, 2, 1, 0, 1, 1, 0],$$

$$B_2 = T(2) [1, 1, 2, 0, 1, 1, 0], B_3 = T^{(2)} [1, 2, 0, 1, 1, 1, 0],$$

$$C_1 = T^{(2)} [1, 0, 3, 0, 1, 1, 0], C_2 = T^{(2)} [1, 3, 0, 0, 1, 1, 0] \dots$$

$$\mathbf{P}_1 \equiv \partial_{1\mu} k_{1\nu},$$

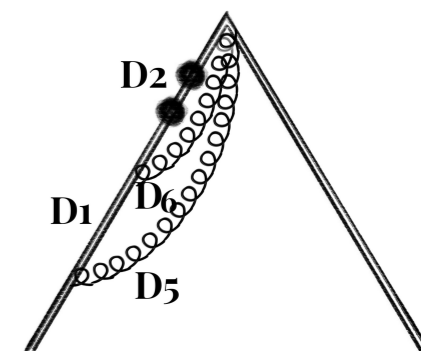
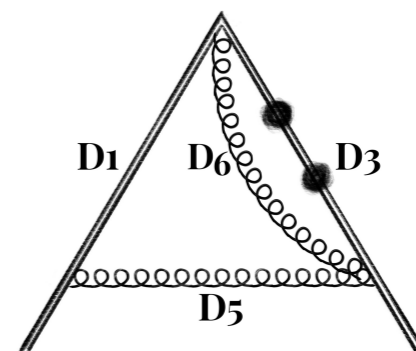
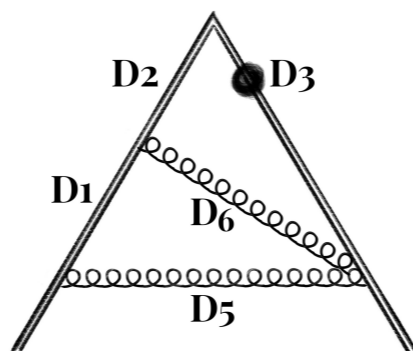
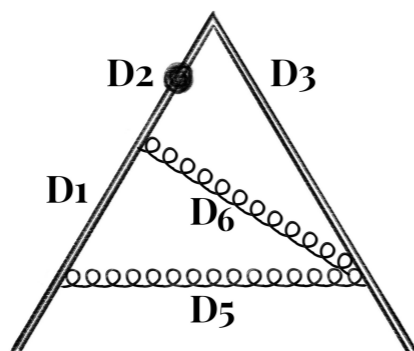
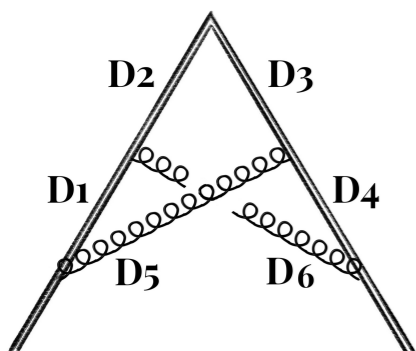
$$\mathbf{P}_2 \equiv \partial_{1\mu} k_{1\nu} v_1^{[\mu} v_2^{\nu]},$$

$$\mathbf{P}_3 \equiv (\partial_{1\mu} k_{1\mu} + \partial_{2\mu} k_{2\nu}) v_1^{[\mu} v_2^{\nu]} \dots$$

$$\mathbf{P}_2 A_1 = 0 \rightarrow 2A_2 + B_3 - 2B_1 - 2(v_1 \cdot v_2) B_2 = 0$$

$$\mathbf{P}_3 A_1 = 0 \rightarrow A_3 - A_2 = 0$$

$$\mathbf{P}_1 B_3 = 0 \rightarrow B_3 - 2C_1 = 0 \dots$$



# Toy example: 2-loop HQET integral

## Cross loop order relation

$$\text{Triangle with loop on right} = \frac{1}{4} \text{Triangle with loop on bottom}$$

$$\text{Triangle with loop on left} = \frac{1}{4} \text{Triangle with loop on top}$$

$$\int \frac{d^{4-2\epsilon} k_1}{i\pi^{2-\epsilon}} \frac{1}{(-2k_1 \cdot v_1 + \delta)^{a_1} (-(k_1 - k_2)^2)} = \frac{1}{(a_1 - 1)(a_1 - 2)} \frac{1}{(-2k_2 v_1 + \delta)^{a_1 - 2 + 2\epsilon}}$$

$$\int \frac{d^{4-2\epsilon} k_1}{i\pi^{2-\epsilon}} \frac{1}{(-2k_1 \cdot v_1 + \delta)^{a_1} (-(k_1 - k_2)^2)} \rightarrow \frac{L-1}{L} \frac{1}{(a_1 - 1)(a_1 - 2)} \frac{1}{(-2k_2 v_1 + \delta)^{a_1 - 2}}$$

admissible differential equation for our IBP system by Initial

$$\frac{d}{dx} \begin{pmatrix} F_1 \\ F_2 \\ F_3 \\ F_4 \end{pmatrix} = \begin{pmatrix} 0 & \frac{2}{x} \\ \frac{0}{x(1-x)(1+x)} & 0 \\ \frac{4}{x} & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} F_1 \\ F_2 \\ F_3 \\ F_4 \end{pmatrix}$$

$$\{F_1, F_2, F_3, F_4, F_5\} = \left\{ \frac{(x^2 - 1)^2}{2x^2} A_1, \frac{x^2 - 1}{x} \left( B_1 + \frac{1 + x^2}{2x} B_2 \right), \frac{2(x^2 - 1)}{x} C_1, C_2, \frac{(x^2 - 1)^2}{2x^2} B_2 \right\}$$

$$\frac{d}{dx} \begin{pmatrix} F_5 \\ F_3 \\ F_4 \end{pmatrix} = \begin{pmatrix} 0 & \frac{1}{x} & 0 \\ \frac{4}{x} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} F_5 \\ F_3 \\ F_4 \end{pmatrix}$$

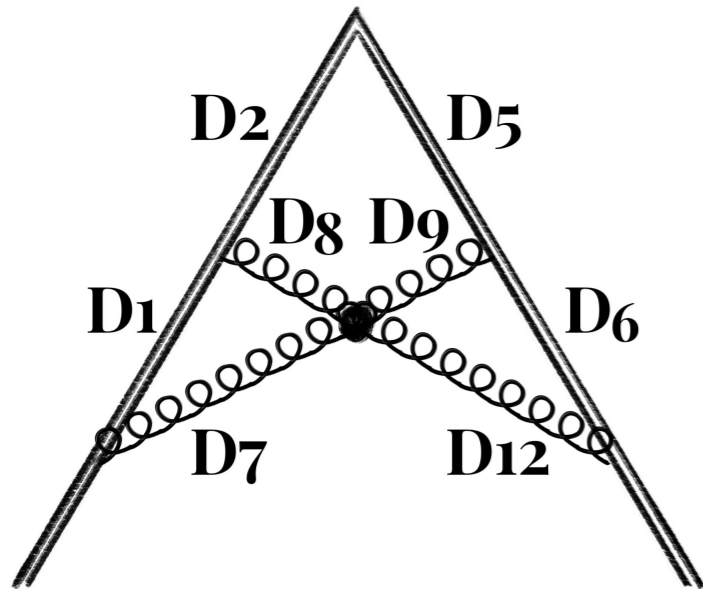
strictly upper triangular matrix

easy to solve with  
boundary condition

$$F_1 = \frac{1}{2} \left[ \frac{1}{3} \ln^3 x + \ln x \operatorname{Li}_2(x^2) - \operatorname{Li}_3(x^2) + \zeta_2 \ln x + \zeta_3 \right]$$

$$F_5 = \frac{1}{4} \ln^2 x$$

# 3-loop HQET



13 irreducible integrals

$$G[1, 1, 0, 0, 1, 1, 1, 1, 1, 0, 0, 1]$$

10503 admissible integrals

23575 4D IBPs

cross loop order relation act on  
the sector  $(1, 0, 0, 0, 1, 1, 0, 1, 1, 0, 0, 1)$   
and the sector  $(1, 0, 0, 0, 1, 1, 1, 0, 1, 1, 0, 0)$

This linear system is very sparse,  
solve with *FiniteFlow* (Peraro 2019) in 2mins with 1 core

$$\{I_1, \dots, I_{13}\} = \left\{ \begin{aligned} &g[2, 2, 0, 0, 1, 1, 1, 1, 1, 0, 0, 1, -1, 0, 0], \quad G[1, 1, -1, 0, 1, 2, 1, 1, 1, 0, 0, 1], \\ &G[1, 1, 0, 0, 1, 1, 1, 1, 1, 0, 0, 1], \quad g[3, 1, 0, 0, 2, 2, 0, 1, 1, 0, 0, 1, 0, 0, -1], \\ &G[3, 1, -1, 0, 1, 2, 0, 1, 1, 0, 0, 1], \quad g[1, 0, 0, 0, 3, 2, 1, 1, 1, 0, 0, 1, 0, 0, -1], \\ &G[1, 0, -1, 0, 2, 2, 1, 1, 1, 0, 0, 1], G[1, -1, 0, 0, 3, 1, 1, 1, 1, 0, 0, 1], G[1, 0, 0, 0, 2, 1, 1, 1, 1, 0, 0, 1], \\ &H_2[2, 0, 1, 1, 1, 0, 1], \quad H_2[1, 0, 2, 1, 1, 0, 1], \quad H_1[1, 1, 1], \quad H_1[0, 2, 1] \end{aligned} \right\}$$

# $I_3^* I_3$ DE matrix

$$\begin{pmatrix}
 -\frac{2(1+x^2)}{(-1+x)x(1+x)} & 0 & 0 & 0 & 0 & \frac{4}{(-1+x)(1+x)} & 0 & 0 & 0 & \frac{2(1+x^2)}{3(-1+x)x(1+x)} & 0 & -\frac{1}{3(-1+x)(1+x)} & 0 \\
 -\frac{2(-1+x)(1+x)}{x^2} & 0 & -\frac{2}{(-1+x)(1+x)} & 0 & 0 & 0 & 0 & 0 & \frac{1+x^2}{(-1+x)x(1+x)} & \frac{4(-1+x)(1+x)}{3x^2} & 0 & 0 & 0 \\
 0 & \frac{2}{(-1+x)(1+x)} & -\frac{2(1+x^2)}{(-1+x)x(1+x)} & 0 & 0 & 0 & 0 & 0 & \frac{2}{(-1+x)(1+x)} & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & -\frac{2(1+x^2)}{(-1+x)x(1+x)} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -\frac{x^2}{3(-1+x)^3(1+x)^3} & \frac{x(1+x^2)}{6(-1+x)^3(1+x)^3} \\
 0 & 0 & 0 & -\frac{2(-1+x)(1+x)}{x^2} & 0 & 0 & 0 & 0 & 0 & -\frac{2}{3(-1+x)(1+x)} & \frac{1+x^2}{3(-1+x)x(1+x)} & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{2(-1+x)(1+x)} & \frac{1}{2(-1+x)(1+x)} & \frac{-1-x^2}{2(-1+x)x(1+x)} & -\frac{2}{3(-1+x)(1+x)} & \frac{-1-x^2}{3(-1+x)x(1+x)} & \frac{-1-x^2}{12(-1+x)x(1+x)} & \frac{1}{2(-1+x)(1+x)} \\
 0 & 0 & 0 & 0 & 0 & -\frac{4(-1+x)(1+x)}{x^2} & 0 & \frac{2(1+x^2)}{(-1+x)x(1+x)} & -\frac{4}{(-1+x)(1+x)} & 0 & 0 & \frac{(-1+x)(1+x)}{3x^2} & 0 \\
 0 & 0 & 0 & 0 & 0 & \frac{2(-1+x)(1+x)}{x^2} & \frac{1+x^2}{(-1+x)x(1+x)} & \frac{1+x^2}{(-1+x)x(1+x)} & -\frac{(1+x^2)^2}{(-1+x)x^2(1+x)} & 0 & \frac{2(-1+x)(1+x)}{3x^2} & -\frac{(-1+x)(1+x)}{6x^2} & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & \frac{2}{(-1+x)(1+x)} & \frac{4}{(-1+x)(1+x)} & -\frac{3(1+x^2)}{(-1+x)x(1+x)} & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -\frac{2(1+x^2)}{(-1+x)x(1+x)} & 0 & \frac{1}{(-1+x)(1+x)} & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{2}{(-1+x)(1+x)} & -\frac{1+x^2}{(-1+x)x(1+x)} & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -\frac{1+x^2}{(-1+x)x(1+x)} & \frac{2}{(-1+x)(1+x)} \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0
 \end{pmatrix}$$



4D initial algorithm

$$\begin{pmatrix}
 0 & -\frac{2}{(-1+x)(1+x)} & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & \frac{1}{x} & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & \frac{1}{x} & \frac{2x}{(-1+x)(1+x)} & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & \frac{1}{x} & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{x} & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{x} \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & -\frac{1}{x} \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0
 \end{pmatrix}$$

8\*8 strictly upper triangular DE matrix

$$\begin{pmatrix} 0 & -\frac{2}{(-1+x)(1+x)} & 0 & 0 & 0 & 0 & 0 & 0 \\ & & \frac{1}{x} & 0 & 0 & 0 & 0 & 0 \\ & & & \frac{1}{x} & \frac{2x}{(-1+x)(1+x)} & 0 & 0 & 0 \\ & & & & & \frac{1}{x} & 0 & 0 \\ & & & & & & \frac{1}{x} & 0 \\ & & & & & & & 0 \\ & & & & & & & -\frac{1}{x} \\ & & & & & & & & 0 \end{pmatrix}$$

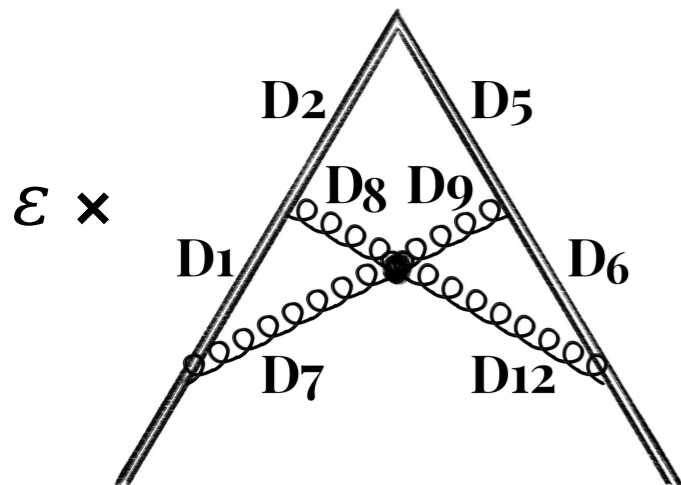
we can integrate  
the UT basis easily



$$J_7 = -J_6 = \frac{8}{3} \ln x, \quad J_5 = -J_4 = \frac{4}{3} \ln^2 x$$

reduce system dimension from 8 to 6

with the boundary condition at  $x=1$



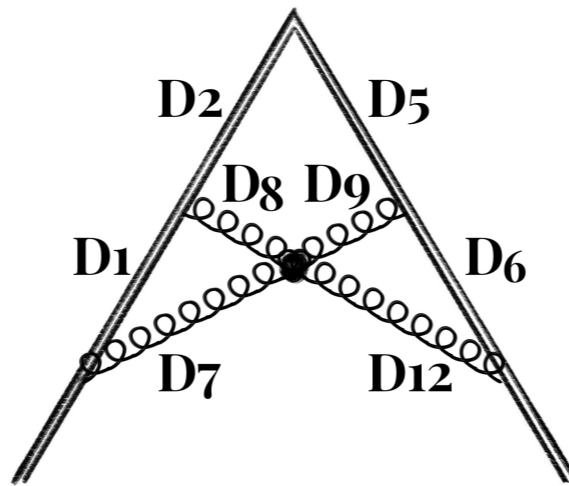
$$\begin{aligned} J_1 = & \frac{8}{3} H_{-1,-2,0,0} - \frac{8}{3} H_{-1,2,0,0} + \frac{8}{3} H_{1,-2,0,0} - \frac{8}{3} H_{1,2,0,0} - \frac{8}{3} H_{-1,0,0,0,0} \\ & - \frac{8}{3} H_{1,0,0,0,0} - \zeta_4 \ln(1-x) + \zeta_4 \ln(1+x) + \frac{4}{3} \zeta_3 \ln(1-x) \ln x \\ & - \frac{4}{3} \zeta_3 \ln x \ln(1+x) + \frac{8}{3} \zeta_3 \text{Li}_2(x) - \frac{2}{3} \zeta_3 \text{Li}_2(x^2) \end{aligned}$$

our method

arXiv[1510.07803]  
Grozin, Heen, Korchemsky and Marquard

6 master integrals to get  $J_1$

39 master integrals to get  $J_1$





# Implementation notes

syzygy generation : A simple Mathematica code

Similar to Schabinger, JHEP 01(2012)001

IBP reduction: *FiniteFlow*

find UT : *4D Initial* algorithm in Mathematica + *FiniteFlow*

# Summary

- Simplified canonical DE for divergent part of Feynman integrals
- Graded Syzygy + Integrating out bubble + 4D Initial
- First application: 3-loop HQET integrals for cusp anomalous dimension
- More generalisation + applications to come ...
- Feynman integral with subdivergences

Thanks