

Refined Topological Strings on Compact Elliptic-fibered Calabi-Yau 3-folds

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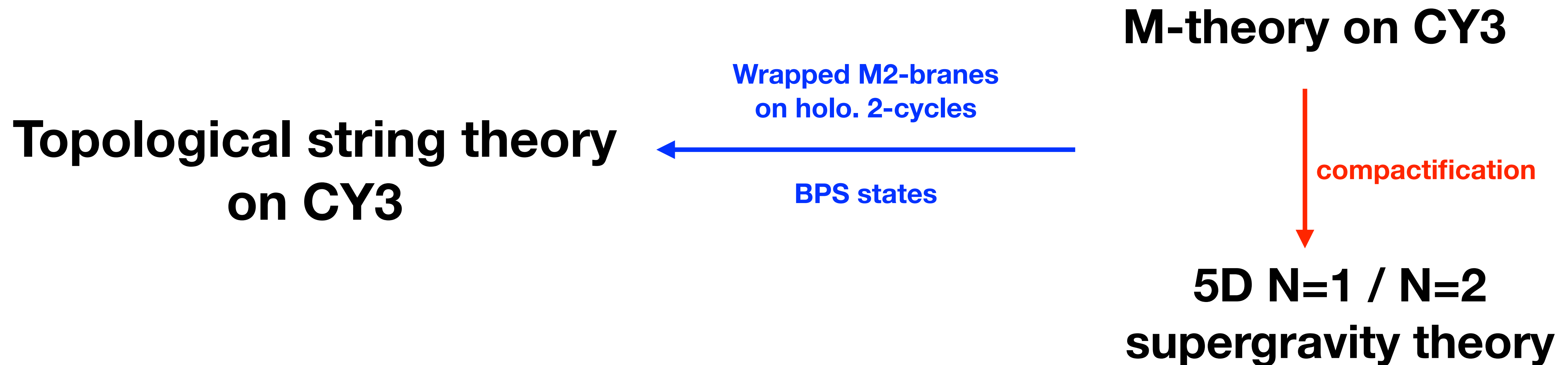
Korea Institute for Advanced Study

第五届全国场论与弦论学术研讨会

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Motivation

For M-theory compactification on a generic Calabi-Yau 3-fold (CY3), the microscopic degeneracies N_{j_L, j_R} of spinning BPS states are captured *exactly* by topological string theory on CY3



Motivation

The BPS states in 5D are characterized by the spin (j_L, j_R) under the spatial rotation group

$$SO(4) \simeq SU(2)_L \times SU(2)_R$$

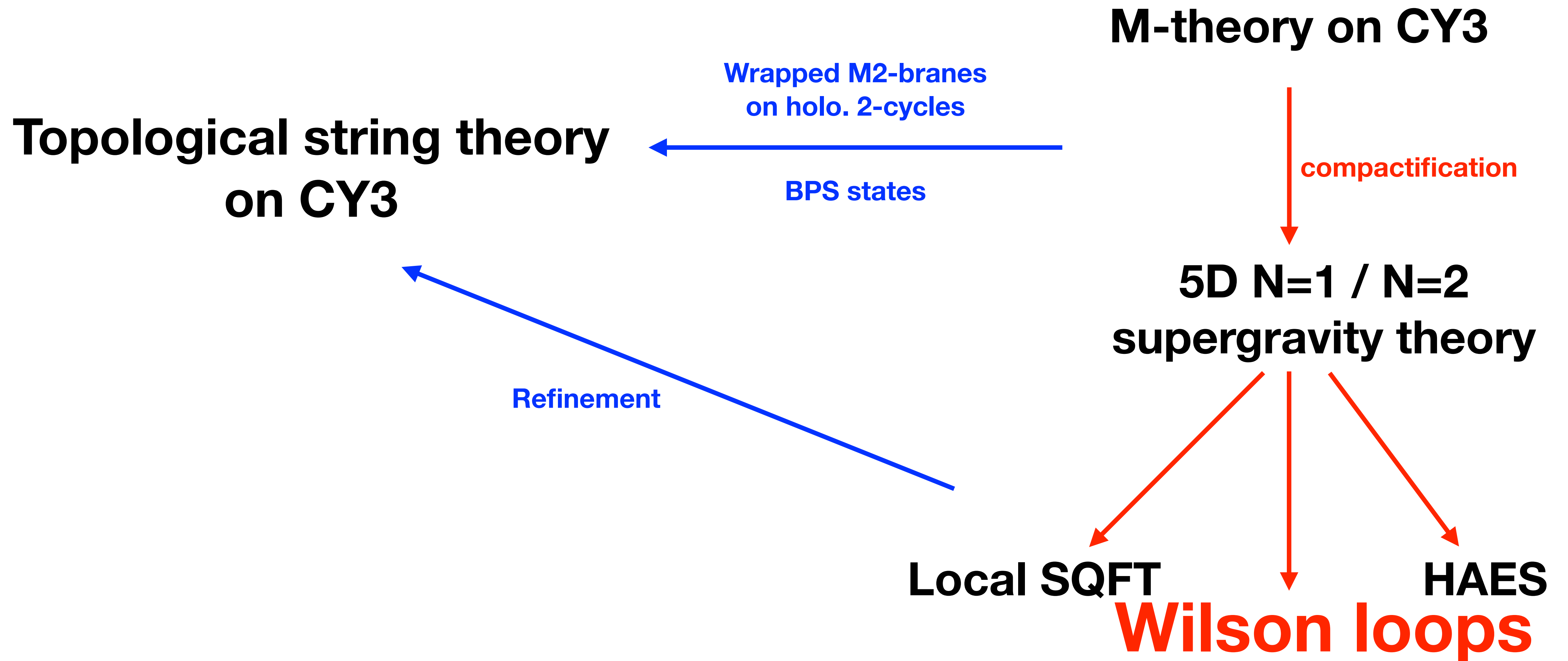
which can be captured by the **topological A-model string amplitudes**

However, the $SU(2)_R$ contents are irrelevant in the conventional topological string calculations.

Turning on $SU(2)_R$ gives a **refined** version of topological strings, known as **refined topological strings** — Hard to define / calculate

The purpose of today's talk :

A proposal for REFINED topological strings on COMPACT Calabi-Yau 3-folds

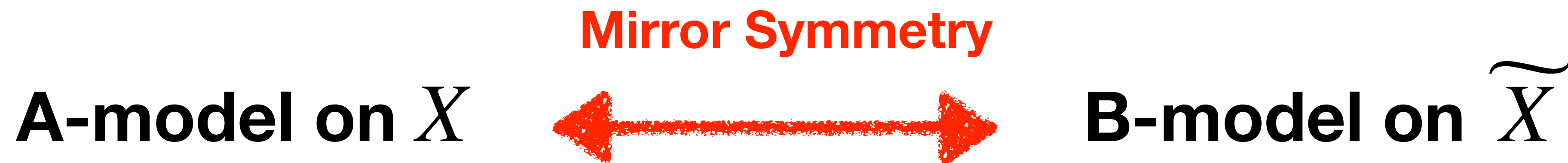


Topological String Theory

- Topological String Theory is a simplified version of string theory focusing on the topological aspects of the string's worldsheet [Witten '89]
- It defines a map from a 2D $\mathcal{N} = (2, 2)$ worldsheet theory to the target space

$$\phi_i : \Sigma \rightarrow X = \text{CY3} \quad \text{vanishing first Chern class}$$

- We can define the **A-model** and **B-model** of the topological string theory by different types of topological twists.



Topological String Theory

- The A-model is defined as Gromov-Witten theory in mathematics.

$$\mathcal{F}^{(g)}(t_i) = \sum_{\beta \in H_2(X; \mathbb{Z})} \text{GW}_g^\beta e^{-\beta \cdot t}$$

- The generating function of these worldsheet instanton numbers is called the free energy of the topological strings. It **only** depends on the Kähler deformation parameters t_i — the masses of the BPS particles

$$\mathcal{F} = \sum_{g=0}^{\infty} \lambda^{2g-2} \mathcal{F}^{(g)}(t_i)$$

Topological String Theory

- The generating function can be expanded in terms of integral **Gopakumar-Vafa (GV) invariants**:

$$\mathcal{F} = \sum_{g=0}^{\infty} \lambda^{2g-2} \mathcal{F}^{(g)}(t_i) = \sum_{\beta, g} \sum_{k=1}^{\infty} \frac{n_g^\beta}{k} \left(2 \sinh\left(\frac{k\lambda}{2}\right) \right)^{2g-2} e^{-k\beta \cdot t}$$

which is derived by [Gopakumar-Vafa '98], from the Schwinger one-loop calculation for an electric particle in a background constant electromagnetic field

- The **GV invariant** is a topological invariant, it is related to the degeneracies

$$N_{j_L, j_R}$$

$$n_g^\beta \sim \sum_{j_R} (2j_R + 1) N_{j_L, j_R}^\beta$$

Topological String Theory

- The A-model GW or GV invariants are usually very hard to compute, but they can be relatively easy to compute from the B-model geometry via mirror symmetry.
- At genus 0, for a hypersurface CY, we can solve the genus 0 GV from the solutions of GKZ systems. E.g. Quintic CY3

d	n_0^d
1	2875
2	609250
3	317206375
4	242467530000
5	229305888887625

Topological String Theory

- Higher genus invariants can be computed from **Holomorphic anomaly equations** [Bershadsky, Cecotti, Ooguri, Vafa '93]

➡ Quintic CY3

– **gap conditions: genus 51(+2)** [Huang, Klemm, Quackenbush '06]

– **Modularities on the DT invariants: genus 80**

[Alexandrov, Feyzbakhsh, Klemm, Pioline, Schimannek '23]

[Alexandrov, Feyzbakhsh, Klemm, Pioline '23]

- Best playground of the resurgence in string theory
- Is there an all-genus expression? Black hole entropy [Vafa's talk at String Math 2024]

Topological String Theory

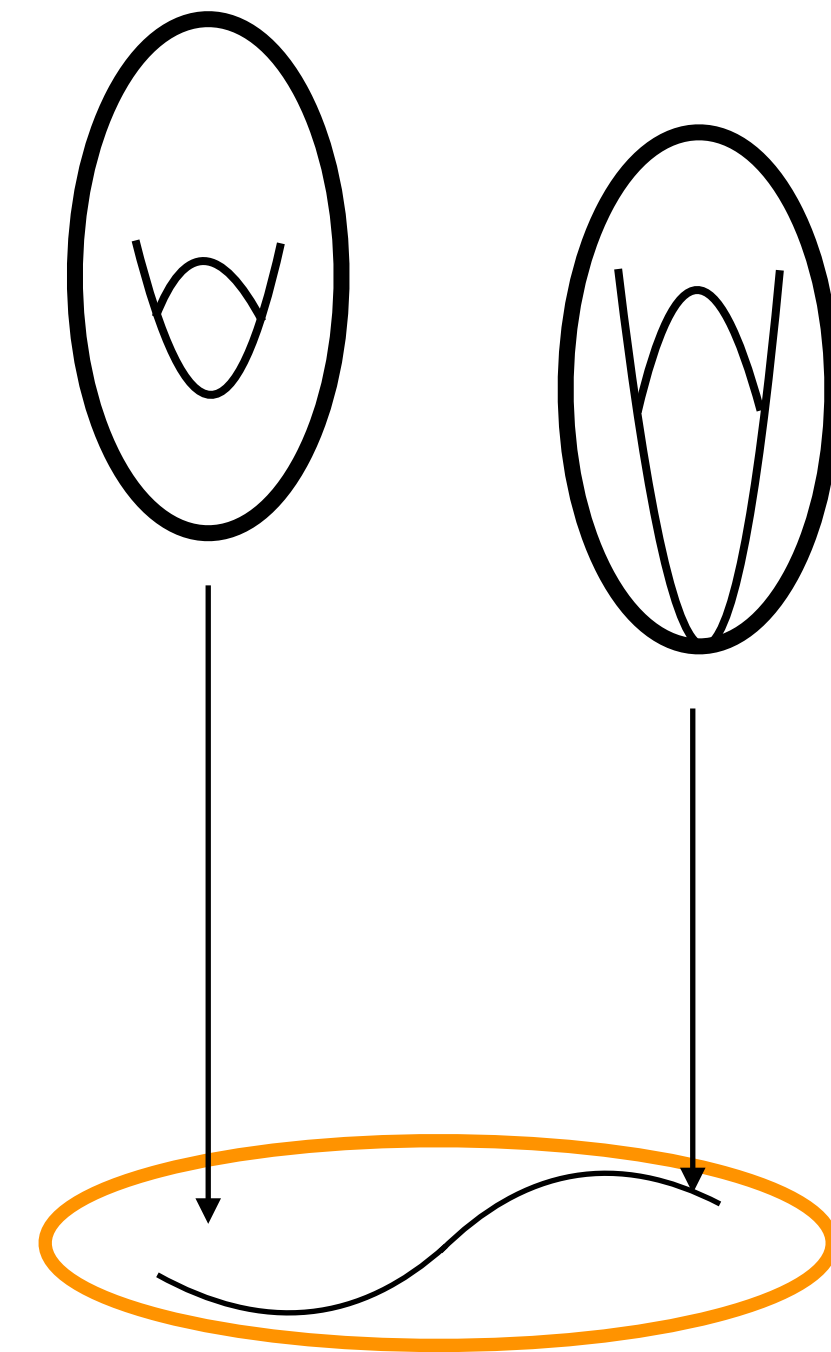
Elliptic fibered Calabi-Yau 3-folds:

E.g. elliptic fibration over \mathbb{P}^2

$$\begin{aligned} l^{(1)} &= (-6 \ ; \ 2 \ 3 \ \boxed{1} \ 0 \ 0 \ 0 \) && \text{Fiber curve} \\ l^{(2)} &= (0 \ ; \ 0 \ 0 \ \boxed{-3} \ 1 \ 1 \ 1 \) && \text{Base curve} \end{aligned}$$

\downarrow
 D

Compact divisor in local \mathbb{P}^2



Topological String Theory

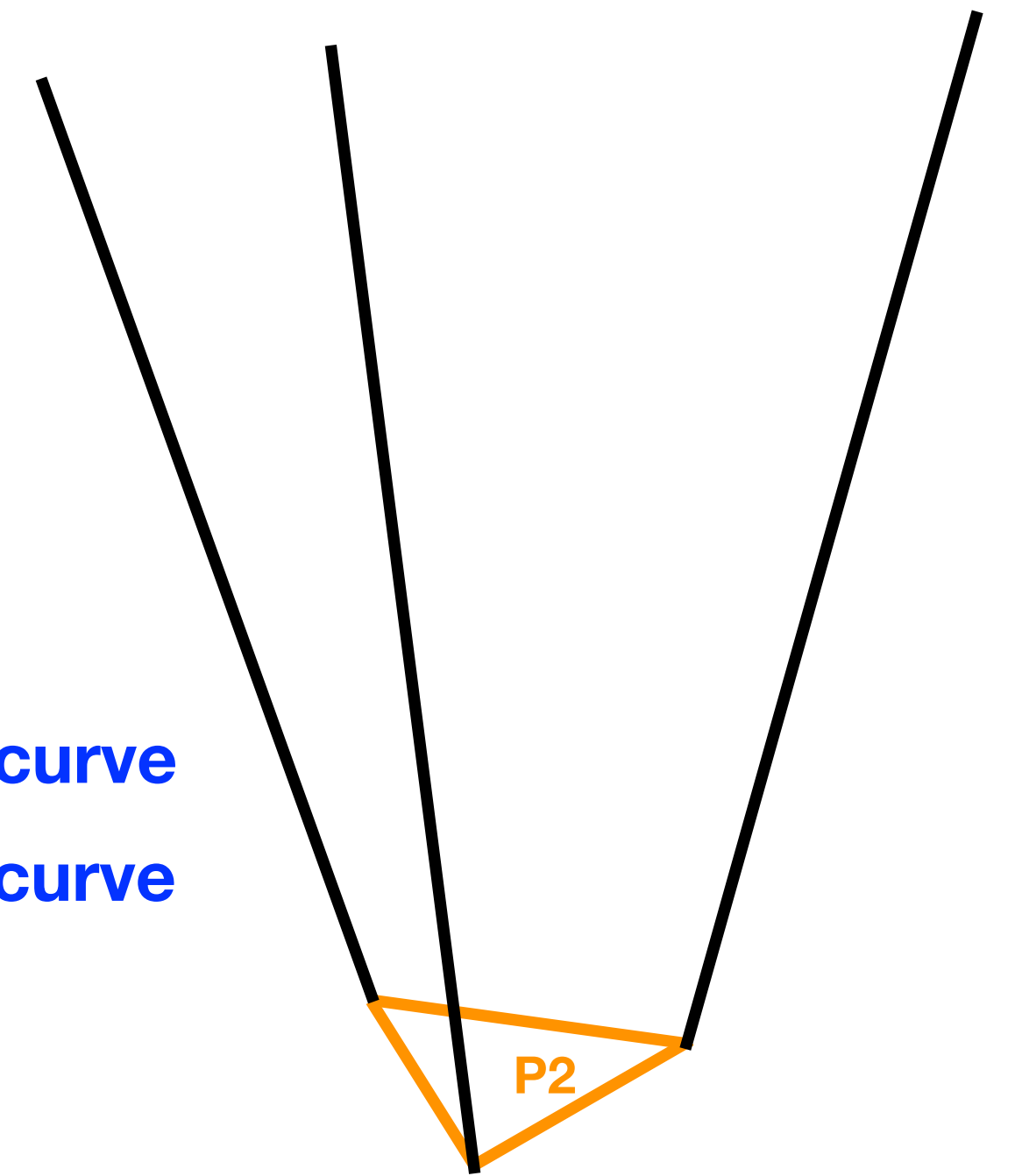
Local Calabi-Yau 3-folds:

E.g. local \mathbb{P}^2 , large fiber limit of elliptic \mathbb{P}^2

$$\begin{aligned} l^{(1)} &= (-6 \ ; \ 2 \ 3 \ \boxed{1} \ 0 \ 0 \ 0) && \text{Fiber curve} \\ l^{(2)} &= (0 \ ; \ 0 \ 0 \ \boxed{-3} \ 1 \ 1 \ 1) && \text{Base curve} \end{aligned}$$

\downarrow
 D

Compact divisor in local \mathbb{P}^2



Local \mathbb{P}^2

- Holomorphic anomaly equation
- Topological vertex
- Topological recursion / Remodeling conjecture

Refined Topological String Theory

Refinement of Local Calabi-Yau 3-folds:

- 4D or 5D supersymmetric gauge theories (on \mathbb{R}^4 or $\mathbb{R}^4 \times S^1$) can be geometric engineered from local Calabi-Yau 3-folds X , from IIA/M-theory compactification on X
- The instanton partition function of the gauge theory in the **(extended) Coulomb branch** can be computed via supersymmetric localization, it is equal to the topological string partition function

$$Z_{\text{gauge}}(t, \lambda) = Z_{\text{top}}(t, \lambda)$$

	Gauge theory	Topological string theory
t	VEV of scalars in the vector multiplets, mass parameters,...	Kähler parameters
λ	Chemical potential for the rotation symmetry	Topological string coupling

Refined Topological String Theory

Refinement of Local Calabi-Yau 3-folds:

- Nekrasov proposed the gauge theory on the Omega-deformed background, by turning on the chemical potentials $\epsilon_{\pm} = \frac{1}{2}(\epsilon_1 \pm \epsilon_2)$ for the 4D rotation symmetry

$$SO(4) \simeq SU(2)_L \times SU(2)_R$$

$$Z_{\text{gauge}}(t, \lambda) \xrightarrow{\text{refinement}} Z_{\text{gauge}}(t, \epsilon_1, \epsilon_2)$$

- The BPS particles are characterized by the spins (j_L, j_R)

Refined Topological String Theory

Refinement of Local Calabi-Yau 3-folds:

- Refined topological strings are proposed from the refinement of gauge theory

$$\begin{aligned} Z_{\text{ref}}(t, \epsilon_1, \epsilon_2) &= Z_{\text{gauge}}(t, \epsilon_1, \epsilon_2) \\ &= \exp \left[\sum_{\beta \in H_2(X, \mathbb{Z})} \sum_{k=1}^{\infty} \sum_{j_L, j_R} (-1)^{2j_L+2j_R} \frac{1}{k} N_{j_L, j_R}^{\beta} \frac{\chi_{j_L}(k\epsilon_-) \chi_{j_R}(k\epsilon_+)}{2 \sinh(\frac{k\epsilon_1}{2}) 2 \sinh(\frac{k\epsilon_2}{2})} e^{-k\beta \cdot t} \right] \end{aligned}$$

- The non-negative number N_{j_L, j_R}^{β} counts the degeneracy of the BPS particle with spin (j_L, j_R) and mass $\beta \cdot t$, it comes from M2-branes wrapping over the curve class β
- It is called refined BPS invariants (non-negative integers)

Refined Topological String Theory

Refined topological strings on non-compact CY3

- Refined holomorphic anomaly equations
- Refined topological vertex
- Blowup equations
- Elliptic fibration over non-compact toric surface (6d SCFTs, 6d LSTs)
 - Modular bootstrap: can be refined

$$F = \sum_{g=0}^{\infty} (\epsilon_1 + \epsilon_2)^{2n} (\epsilon_1 \epsilon_2)^{g-1} F^{(n,g)}(t_i)$$

genus $g \rightarrow$ genus (n, g)

Refined Topological String Theory

Known results for refined topological strings on compact CY3's

- $K3 \times T2$ [Katz, Klemm, Pandharipande, '14]
- Elliptic CY3, $(n \leq 1, g)$ done by [Huang Katz, Klemm, '20]

Our result: refined topological strings on compact (elliptic-fibered) CY3 for any (n, g)

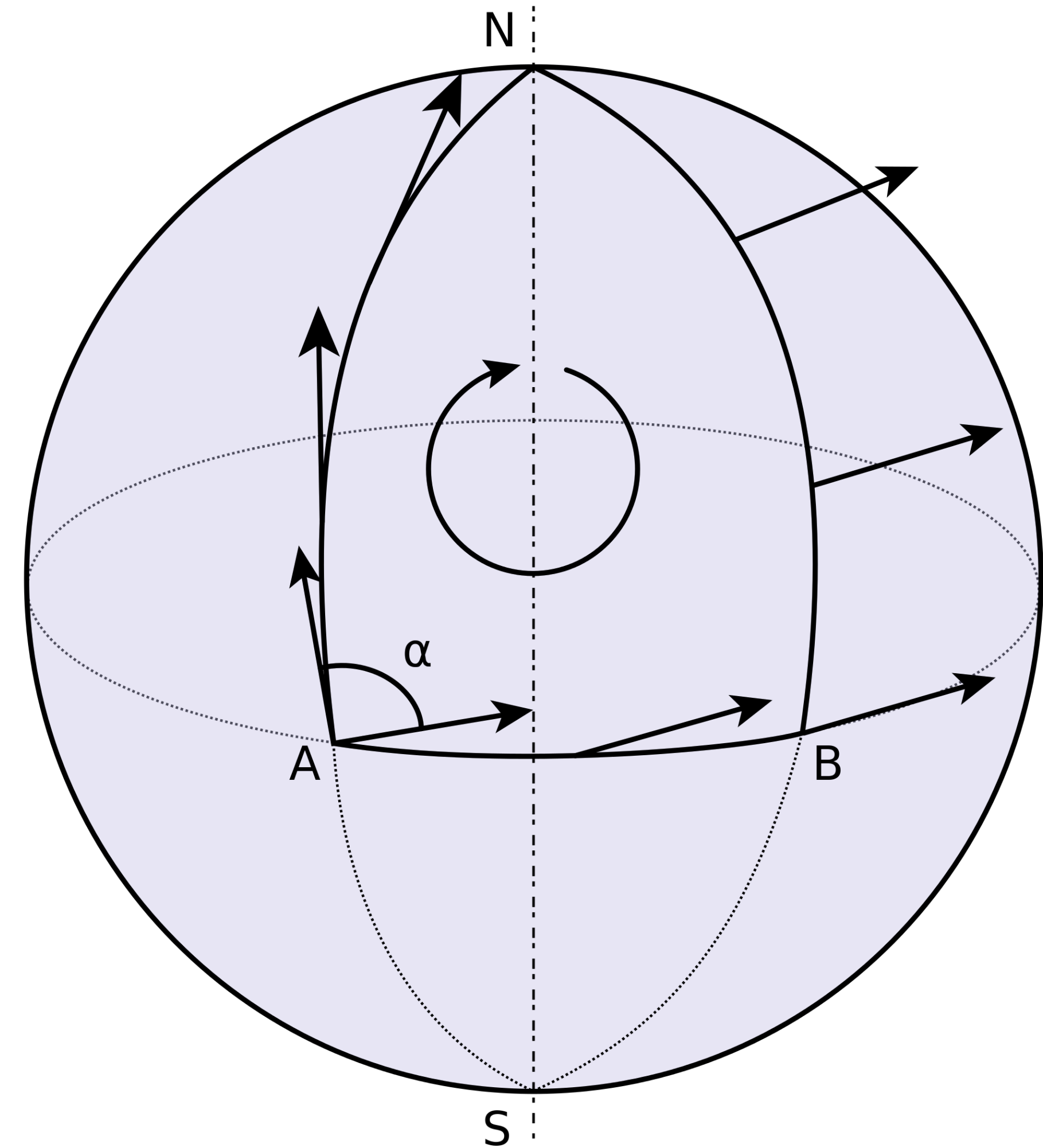
$$F = \sum_{g=0}^{\infty} (\epsilon_1 + \epsilon_2)^{2n} (\epsilon_1 \epsilon_2)^{g-1} F^{(n,g)}(t_i)$$
$$= \exp \left[\sum_{\beta \in H_2(X, \mathbb{Z})} \sum_{k=1}^{\infty} \sum_{j_L, j_R} (-1)^{2j_L + 2j_R} \frac{1}{k} N_{j_L, j_R}^{\beta} \frac{\chi_{j_L}(k\epsilon_-) \chi_{j_R}(k\epsilon_+)}{2 \sinh(\frac{k\epsilon_1}{2}) 2 \sinh(\frac{k\epsilon_2}{2})} e^{-k\beta \cdot t} \right]$$

Remark: The number N_{j_L, j_R}^{β} may not be an invariant but depends on the complex deformation of the CY3

Wilson loop

- In gauge theory, Wilson lines/loops are gauge invariant operators
- They arise from the parallel transport of gauge variables around closed loops
- They can be generated from the worldline C of static infinitely massive quarks in rep. \mathbf{r} —
Polyakov loop

$$W_{\mathbf{r}}(C) = \text{Tr}_{\mathbf{r}} \left[\mathcal{P} \exp \left(i \oint_C A_{\mu} dx^{\mu} \right) \right]$$



Wilson loop

- Supersymmetric version

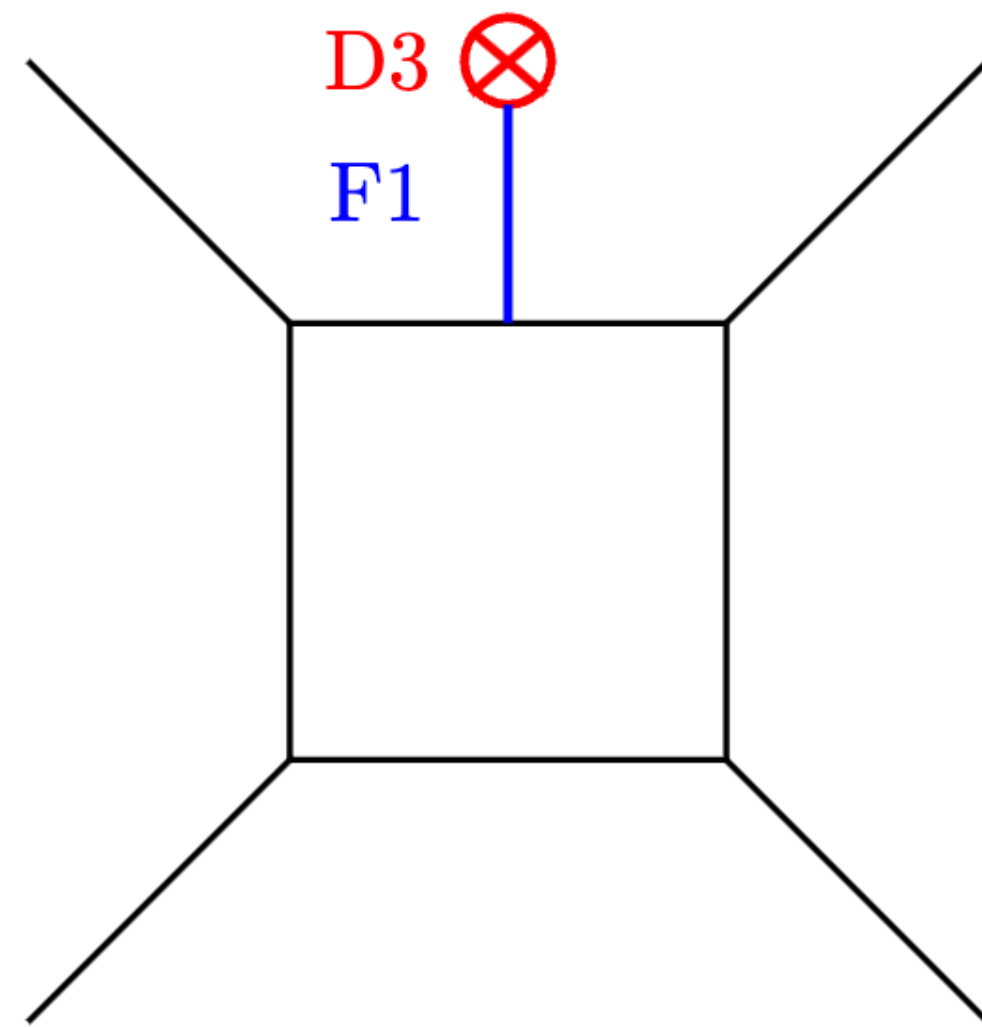
$$W_{\mathbf{r}}[C] = \text{Tr}_{\mathbf{r}} \mathcal{P} \exp \int_C (iA_{\mu} \dot{x}^{\mu} + |\dot{x}| \phi) ds$$

- a scalar is added to preserve some of the supersymmetries.
- Half-BPS operators in 5D $\mathcal{N} = 1$ gauge theories on S^1
 - The rotation symmetry $SO(4)$ is preserved
 - GV-like expansion [Huang, Lee, XW, '22][Kim, Kim, Kim, '21]

Wilson loop

- **Coulomb branch:** the scalar field ϕ gets the expectation value in the Cartan subalgebra of the gauge group, which breaks the gauge group to $U(1)^r$.
- The scalar expectation values $\phi_i, i = 1, \dots, r$ parametrize the moduli space on the Coulomb branch.
- The **representation** of the Wilson loop becomes the **electric charge** of the **Wilson loop particle** — — a **heavy, stationary electric particle** located at the origin of the space \mathbb{R}^4

Half-BPS Wilson loop operators – IIB realization [David Tong, '14] ...

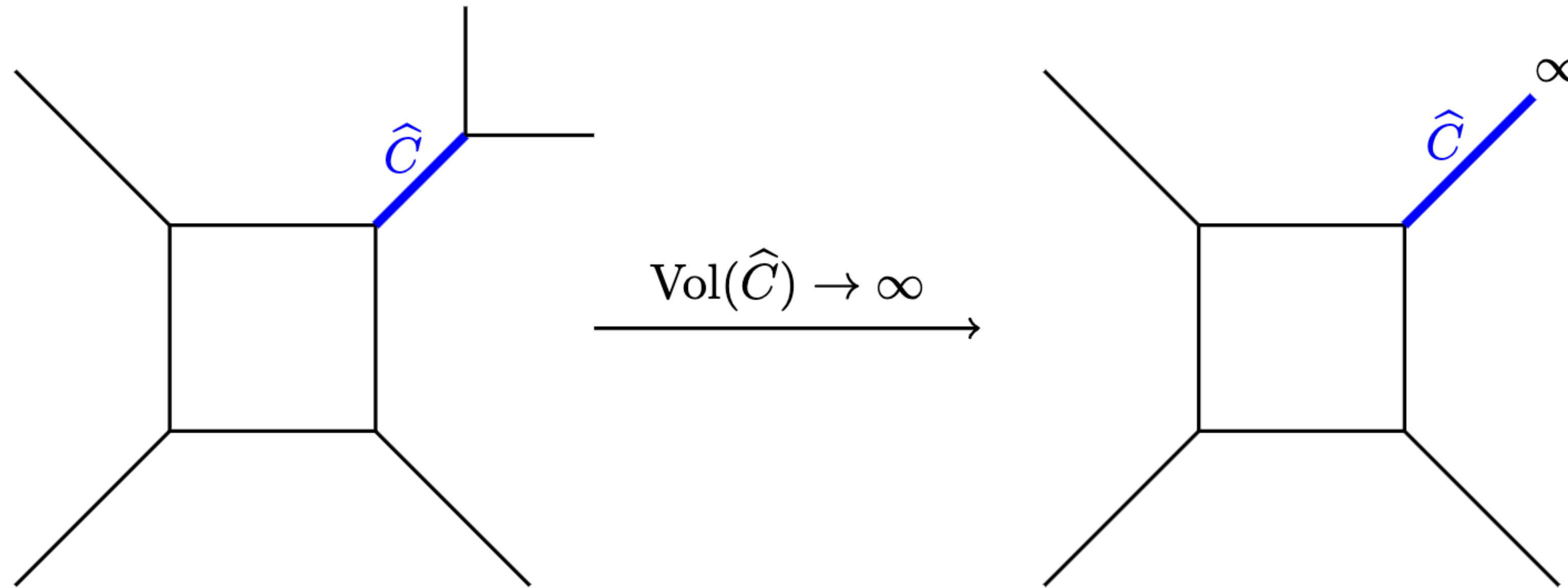


	0	1	2	3	4	5	6	7	8	9
D5	•	•	•	•	•	•				
NS5	•	•	•	•	•		•			
$5_{(p,q)}$	•	•	•	•	•	θ	θ			
F1	•						•			
D3	•							•	•	•

The half-BPS Wilson loop operators are realized by adding **semi-infinite F1 strings** with charge 1, stretched between D3 branes and D5 branes.

The lowest energy modes on such F1 strings are **fermionic**, so there can be at most one F1 string stretched between a D3 brane and a D5 brane.

Half-BPS Wilson loop operators – IIB realization



To have a geometric realization, we can also add a flavor matter by blowing up the geometry at one point.

In the large volume limit of the exceptional curve, only single F1 states contribute

Half-BPS Wilson loop operators – IIB realization

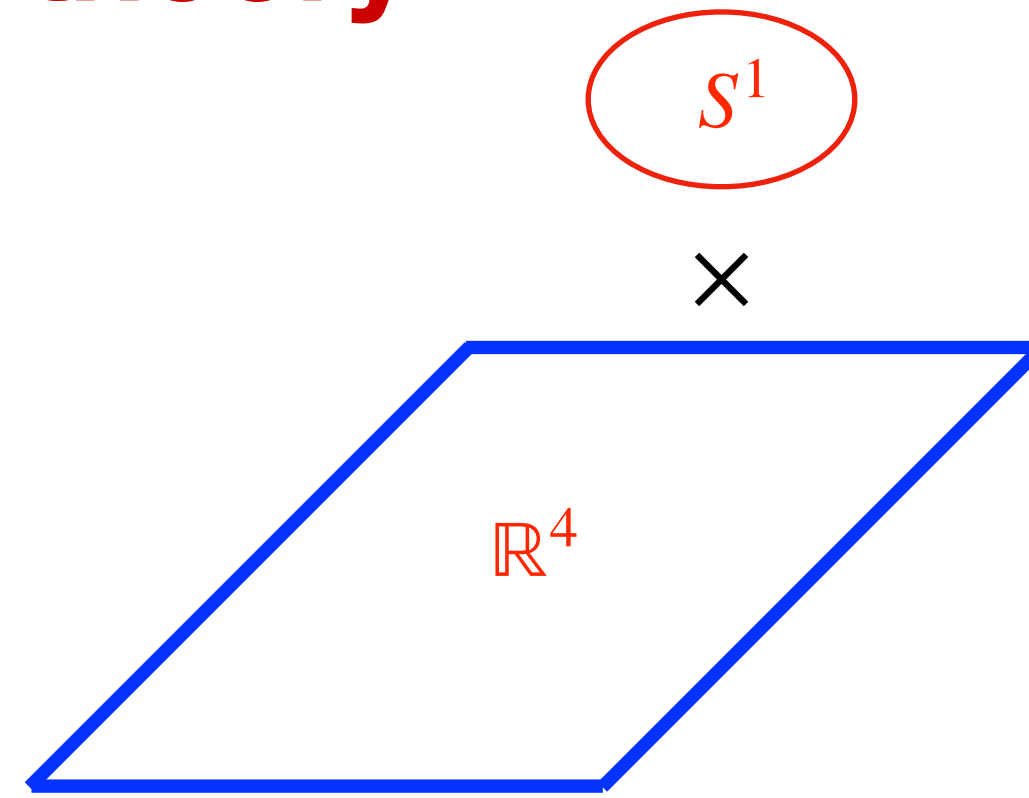
- The geometric definition can be extended to non-gauge theory, e.g. local P2

d_h/d_C	0	1	2
0	0	1	0
1	3	-2	1
2	-6	5	-4
3	27	-32	35
4	-192	286	-400
5	1695	-3038	5187

Genus 0 invariants for Wilson loops of local P2

- And theories with broken one-form symmetry

Realization in M theory

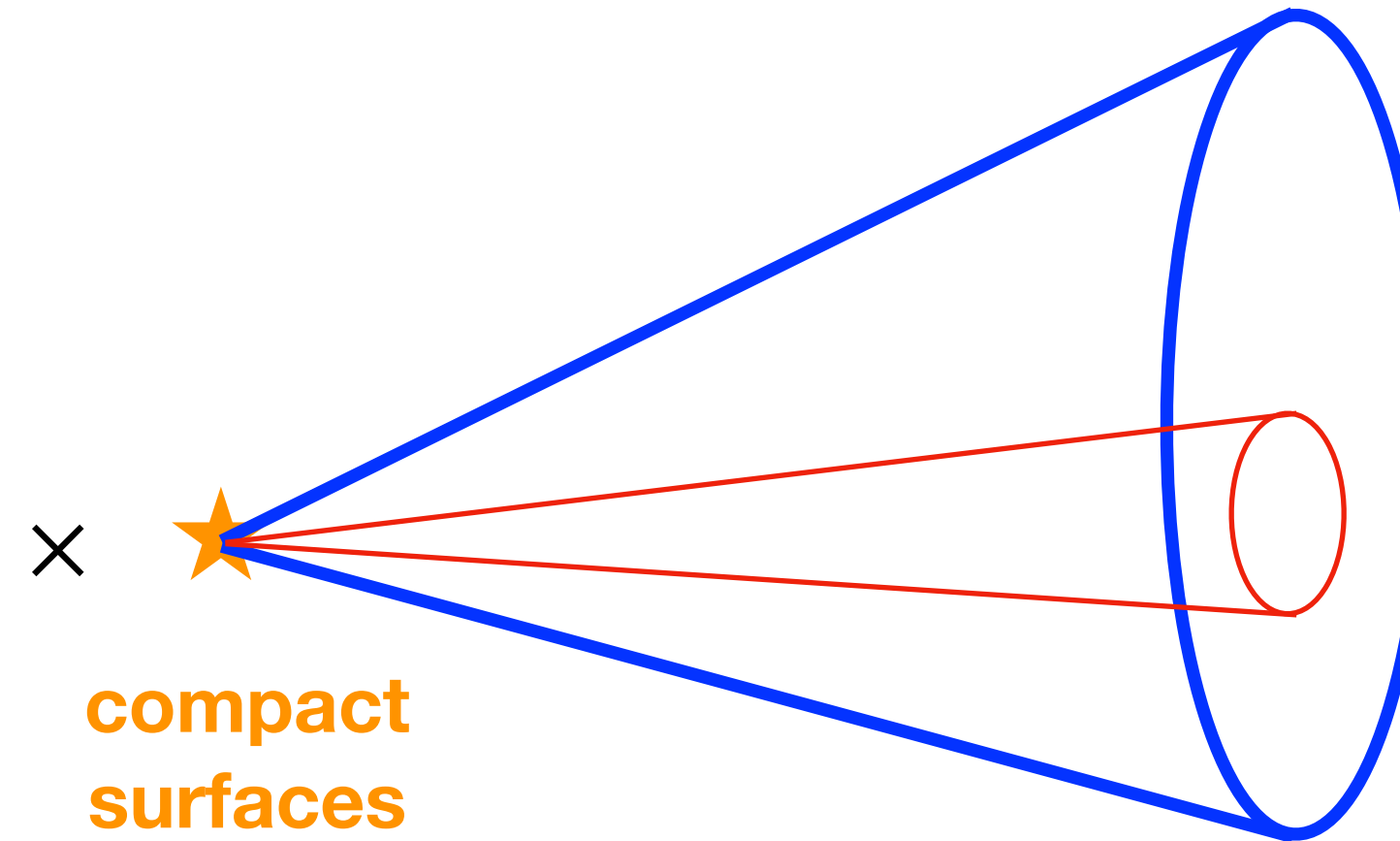


Wilson loop

heavy, stationery electric particle

BPS particles

Representation



Geometry X

a non-compact curve C in CY3 extended to infinity
Heavy

M2-branes wrapping around $C + C, C \in H_2(X; \mathbb{Z})$

Charges of C

$$q_i = D_i \cdot C$$

Compact divisor

How the Wilson loops of a local CY3 are connected to the compact CY3

Genus 0 GV invariants:

d_h/d_C	0	1	2
0	0	1	0
1	3	-2	1
2	-6	5	-4
3	27	-32	35
4	-192	286	-400
5	1695	-3038	5187

Wilson loops for local P2

d_h/d_C	0	1	2
0	0	1	0
1	3	-2	0
2	-6	5	0
3	27	-32	7
4	-192	286	-110
5	1695	-3038	1651

one point blowup of local P2

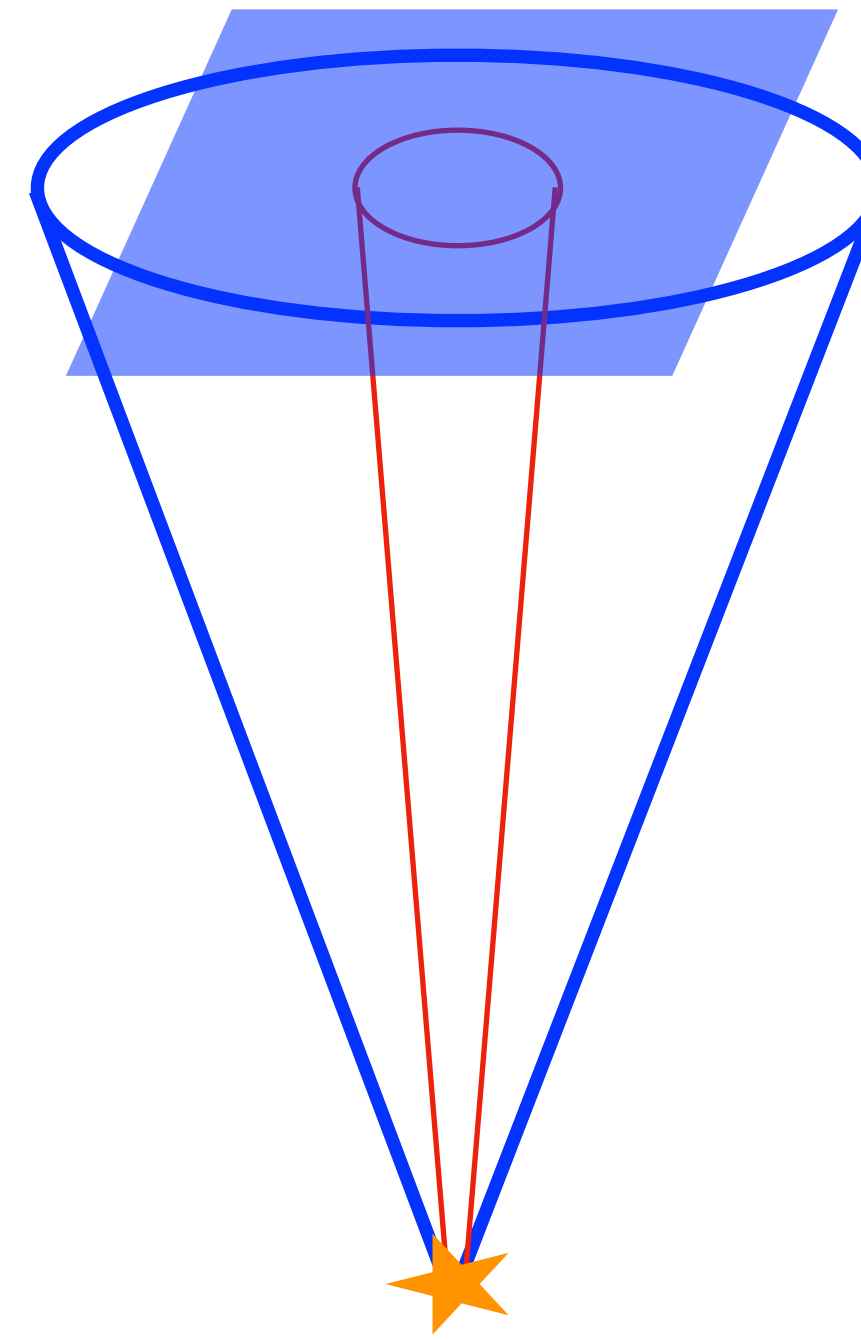
d_h/d_C	0	1	2
0	0	$540 \times (1)$	540
1	3	$540 \times (-2)$	143370
2	-6	$540 \times (5)$	-574560
3	27	$540 \times (-32)$	5051970
4	-192	$540 \times (286)$	-57879900
5	1695	$540 \times (-3038)$	751684050

elliptic P2

The Euler number for elliptic P2 is -540

This is not a coincidence, the red numbers are all related to Wilson loops of local P2

How the Wilson loops of a local CY3 are connected to the compact CY3



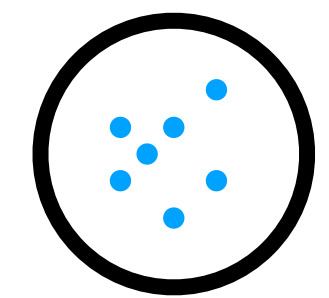
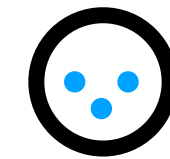
Ending of the 2-cycle at infinity

Wilson loop of local P2

dP1

elliptic P2

Wilson loop
particle



1

1

$$\frac{1}{2 \sinh(\frac{\epsilon_1}{2}) \cdot 2 \sinh(\frac{\epsilon_2}{2})}$$

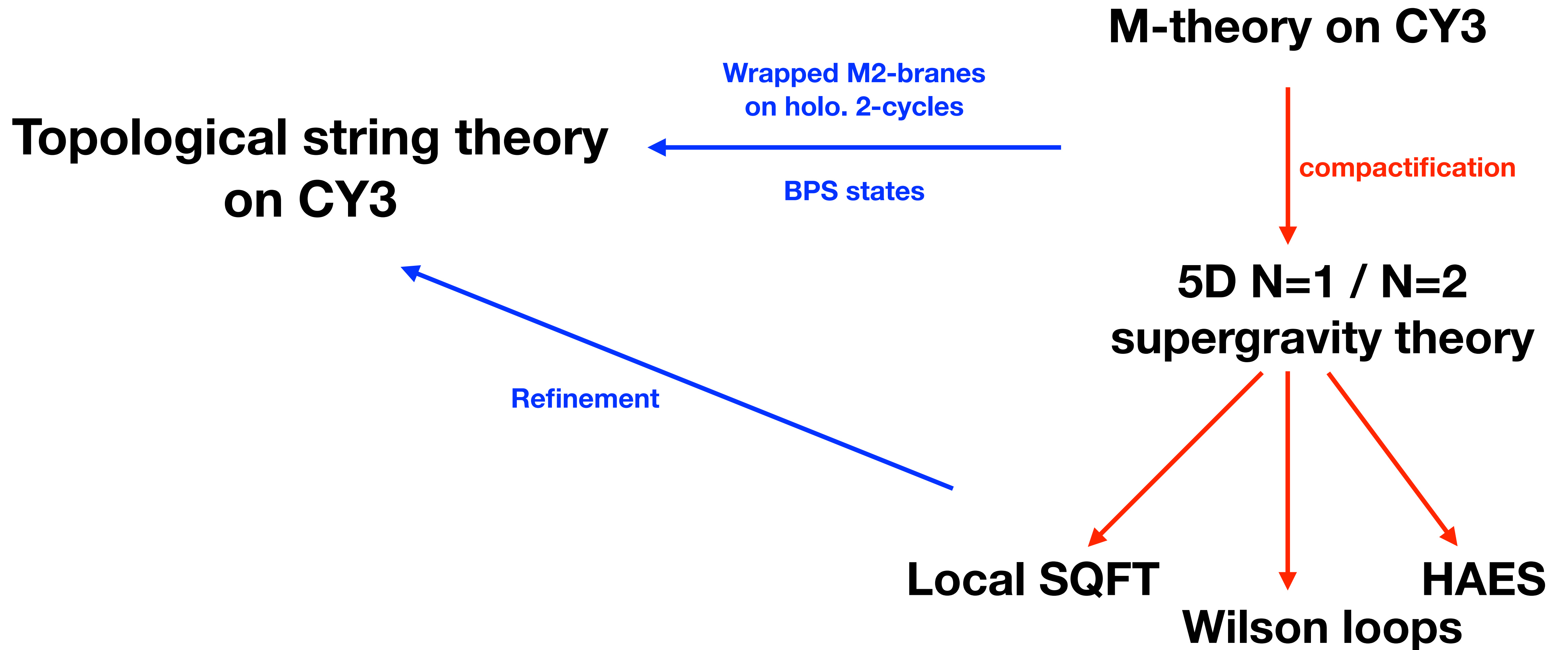
$$\frac{546[0, 0] \oplus [1, \frac{1}{2}]}{2 \sinh(\frac{\epsilon_1}{2}) \cdot 2 \sinh(\frac{\epsilon_2}{2})}$$

If a half-BPS particle in the 5D supergravity theory is heavy enough

It becomes the half-BPS Wilson loop particle in the local 5D quantum field theory

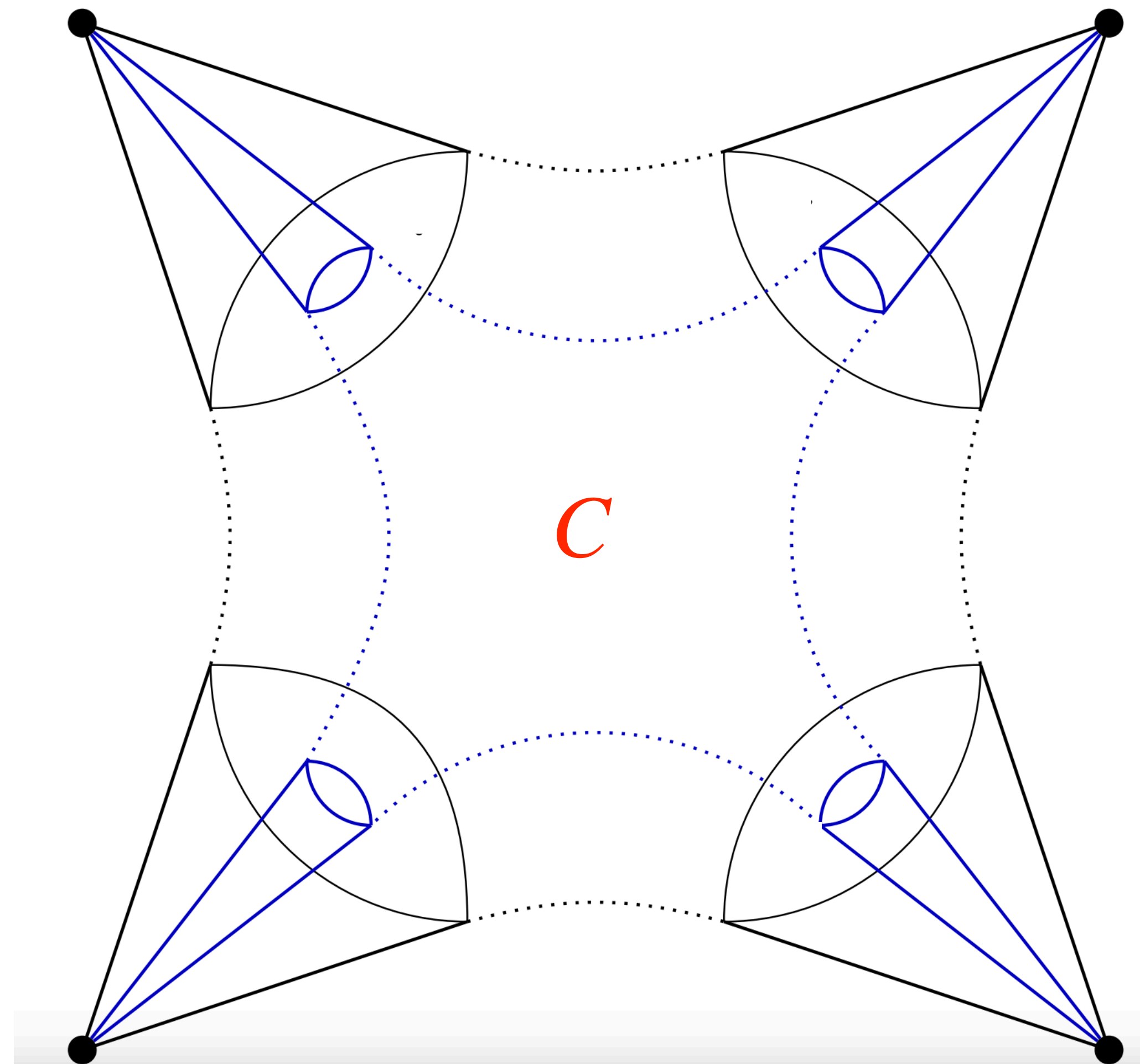
We can refine the BPS spectrum of the 5D supergravity theory via refinement of the 5D Wilson loops

Conclusion



Conclusion

We can also consider the gluing of multiple local theories



Thank you!