

# Progress on Finite Feynman Integrals with Uniform Transcendentality



Yang Zhang 张扬

University of Science and Technology of China

Peng Huanwu Center for Fundamental Theory

Nov. 29, 2020



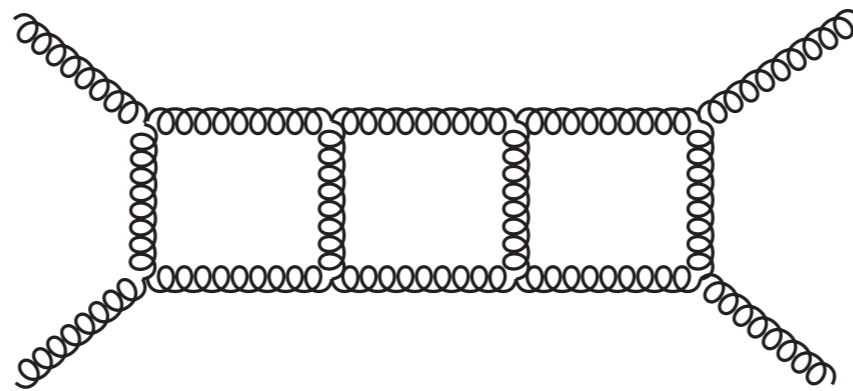
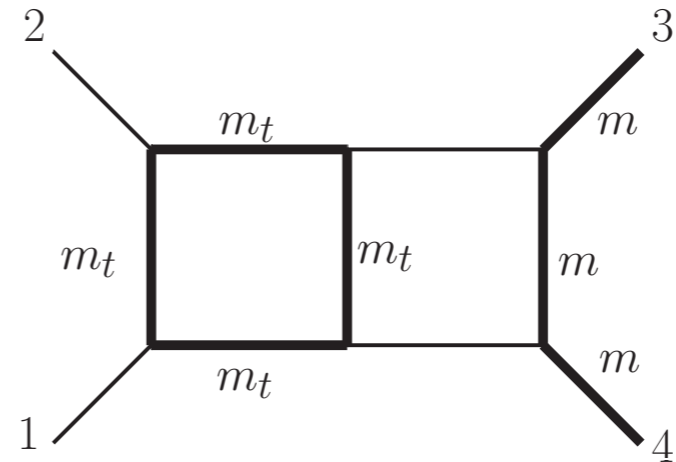
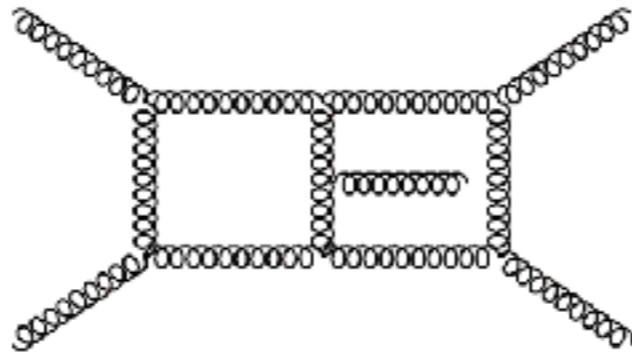
This talk is based on the ongoing project

with

Johannes Henn and Kai Yan (颜开)  
Max-Planck Institute for Physics, Munich



# Feynman integral evaluation



Multi-loop Feynman integrals are the **hardcore** objects for a perturbative QFT computation

- Important for high-energy phenomenology
- Theoretically important

# Feynman integral evaluation: the State of Art

## recent analytic results

- 2-loop qqbar  $\rightarrow$  ttbar nonplanar integrals  
(Di Vita, Gehrmann, Laporta, Mastrolia, Pierpaolo; Primo, Schubert, Ulrich, 2019 )
- 4-loop form factor planar integrals  
(von Manteuffel, Schabinger, 2019)
- 2-loop Higgs + one jet production (with t quark mass dep.) nonplanar integrals (almost all)  
(Bonciani, Del Duca, Frellesvig, Henn, Hidding, Maestri, Moriello, Salvatori, Smirnov, 2019)
- 2-loop five-point massless planar and nonplanar integrals  
(Chicherin, Gehrmann, Henn, Wasser, YZ, Zoia, 2019)

....

and there would be many more in the near future

## uncharted territory

Feynman integral with elliptic polylogarithms,  
usually in multi-loop nonplanar diagrams with “large” massive internal loop  
new transcendental functions (new, even for mathematicians)

see, eg, Broedel, Duhr, Dulat, Tancredi, 2018

# Mainstream Analytic Methods

## Differential equation for Feynman integrals

$d = 4 - 2\epsilon$   $I(\bar{x}, \epsilon)$  master integrals as a column vector

$$\frac{\partial}{\partial x_i} I(\bar{x}, \epsilon) = A_i(\bar{x}, \epsilon) I(\bar{x}, \epsilon) \quad (\text{Kotikov 1991})$$

Can be used both numerically and analytically

eg. Gehrmann, Remiddi 1999, Papadopoulos 2014, Liu, Ma, Wang, 2018 ...

## Canonical Differential equation

$$\frac{\partial}{\partial x_i} \tilde{I}(\bar{x}, \epsilon) = \epsilon \tilde{A}_i(\bar{x}) \tilde{I}(\bar{x}, \epsilon)$$

Proportional to  $\epsilon$

$$= \epsilon \left( \sum_{l=1} \frac{\partial \log(W_l)}{\partial x_i} m_l \right) \tilde{I}(\bar{x}, \epsilon)$$

Symbol letters


rational number matrix

(Henn 2013)

# Canonical Differential Equation

Can be solved order-by-order analytically

$$\tilde{I}(x) = P \exp \left( \epsilon \int_{\mathcal{C}} dA \right) \tilde{I}(x_0)$$



path-ordered

Chen's (陈国才) iterated integrals,  
homotopically invariant

In term of poly-logarithm functions

$$G(\underbrace{0, \dots, 0}_k; z) = \frac{1}{k!} (\log z)^k, \quad G(a_1, \dots, a_k; z) = \int_0^z \frac{dt}{t - a_1} G(a_2, \dots, a_k; t)$$

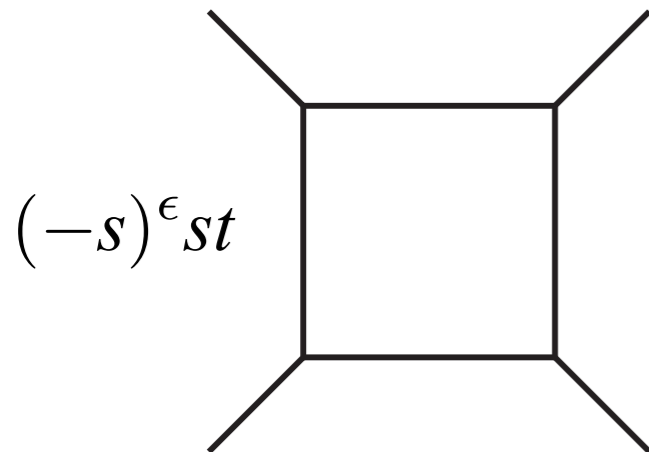
implemented in the C++ library  
GiNaC

# Uniformly transcendental (UT) integrals $\Rightarrow$ Canonical DE

$$I(x) = \epsilon^{-2L} \sum_{i=0}^{\infty} f_i(x) \epsilon^i$$

$$\mathcal{T}(\log) = 1, \mathcal{T}(\pi) = 1, \mathcal{T}(\zeta_n) = n, \mathcal{T}(\text{Li}_n) = n, \dots, \mathcal{T}(f_1 f_2) = \mathcal{T}(f_1) + \mathcal{T}(f_2)$$

$$\mathcal{T}(f_i) = i, \quad \mathcal{T}(\partial_x f_i) = i - 1$$

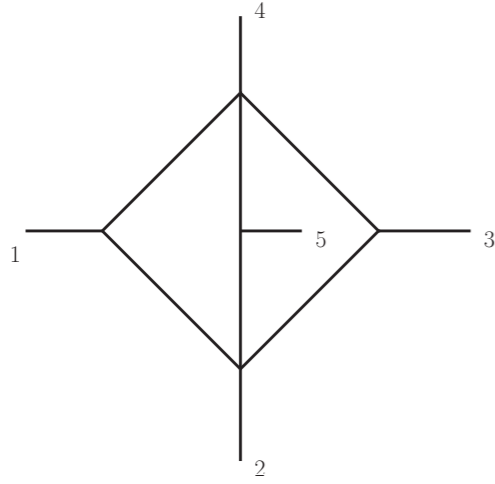


$$\begin{array}{cccc} 0 & 1 & 2 & 3 \\ \frac{4}{\epsilon^2} - \frac{2 \log(x)}{\epsilon} - \frac{4\pi^2}{3} - \frac{1}{6}\epsilon \left( -12\text{Li}_3(-x) + 12\text{Li}_2(-x) \log(x) \right. \\ \left. -2 \log^3(x) + 6 \log(x+1) \log^2(x) - 7\pi^2 \log(x) + 6\pi^2 \log(x+1) + 68\zeta(3) \right) \end{array}$$

UT

# Uniformly transcendental (UT) integrals: more complicated example

Chicherin, Gehrmann, Henn, Wasser, YZ, Zoia (2019)



$$I = \frac{1}{\epsilon^2} f^{(2)} + \frac{1}{\epsilon} f^{(3)} + f^{(4)} + \mathcal{O}(\epsilon).$$

$$f^{(2)} = -3 \left[ \text{Li}_2\left(\frac{1}{W_{27}}\right) - \text{Li}_2(W_{27}) + \text{Li}_2\left(\frac{1}{W_{28}}\right) - \text{Li}_2\left(\frac{1}{W_{27}W_{28}}\right) - \text{Li}_2(W_{28}) + \text{Li}_2(W_{27}W_{28}) \right].$$

$W_1 = v_1,$	$W_6 = v_3 + v_4,$	$W_{11} = v_1 - v_4,$	$W_{16} = v_1 + v_2 - v_4,$	$W_{21} = v_3 + v_4 - v_1 - v_2,$	$W_{26} = \frac{v_1 v_2 - v_2 v_3 + v_3 v_4 - v_1 v_5 - v_4 v_5 - \sqrt{\Delta}}{v_1 v_2 - v_2 v_3 + v_3 v_4 - v_1 v_5 - v_4 v_5 + \sqrt{\Delta}},$
$W_2 = v_2,$	$W_7 = v_4 + v_5,$	$W_{12} = v_2 - v_5,$	$W_{17} = v_2 + v_3 - v_5,$	$W_{22} = v_4 + v_5 - v_2 - v_3,$	$W_{27} = \frac{-v_1 v_2 + v_2 v_3 - v_3 v_4 - v_1 v_5 + v_4 v_5 - \sqrt{\Delta}}{-v_1 v_2 + v_2 v_3 - v_3 v_4 - v_1 v_5 + v_4 v_5 + \sqrt{\Delta}},$
$W_3 = v_3,$	$W_8 = v_5 + v_1,$	$W_{13} = v_3 - v_1,$	$W_{18} = v_3 + v_4 - v_1,$	$W_{23} = v_5 + v_1 - v_3 - v_4,$	$W_{28} = \frac{-v_1 v_2 - v_2 v_3 + v_3 v_4 + v_1 v_5 - v_4 v_5 - \sqrt{\Delta}}{-v_1 v_2 - v_2 v_3 + v_3 v_4 + v_1 v_5 - v_4 v_5 + \sqrt{\Delta}},$
$W_4 = v_4,$	$W_9 = v_1 + v_2,$	$W_{14} = v_4 - v_2,$	$W_{19} = v_4 + v_5 - v_2,$	$W_{24} = v_1 + v_2 - v_4 - v_5,$	$W_{29} = \frac{v_1 v_2 - v_2 v_3 - v_3 v_4 - v_1 v_5 + v_4 v_5 - \sqrt{\Delta}}{v_1 v_2 - v_2 v_3 - v_3 v_4 - v_1 v_5 + v_4 v_5 + \sqrt{\Delta}},$
$W_5 = v_5,$	$W_{10} = v_2 + v_3,$	$W_{15} = v_5 - v_3,$	$W_{20} = v_5 + v_1 - v_3,$	$W_{25} = v_2 + v_3 - v_5 - v_1,$	$W_{30} = \frac{-v_1 v_2 + v_2 v_3 - v_3 v_4 + v_1 v_5 - v_4 v_5 - \sqrt{\Delta}}{-v_1 v_2 + v_2 v_3 - v_3 v_4 + v_1 v_5 - v_4 v_5 + \sqrt{\Delta}},$



However, for multi-loop Feynman integrals,  
it is not easy to find a complete UT basis

(4D Leading Singularity/dlog) Arkani-Hamed, Bourjaily, Cachazo, Trnka (2010)

(Matrix transformation) Meyer (2015)

(Morse theory for DE) Lee (2015)

(4D dlog algorithm) Wasser (2017)

(D-dim Leading Singularity) Chicherin, Gehrmann, Henn, Wasser, YZ, Zoia, (2019)

(D-dim Intersection) Chen, Xu, Yang, (2020) ... ..

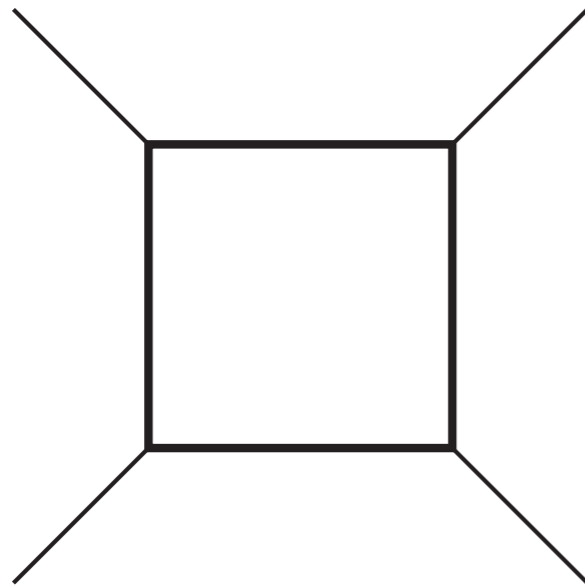
In practice, sometimes

it is **not necessary** to use a complete UT basis

# Finite UT integrals

## Example

Caron-Huot and Henn 2014



internal lines massive

$$g_1 = 2m^2 G[0, 0, 0, 3] / \epsilon^2$$

$$g_2 = -\sqrt{s(s - 4m^2)} G[1, 0, 2, 0] / \epsilon$$

$$g_3 = -\sqrt{t(t - 4m^2)} G[0, 1, 0, 2] / \epsilon$$

$$g_4 = s G[1, 1, 1, 0]$$

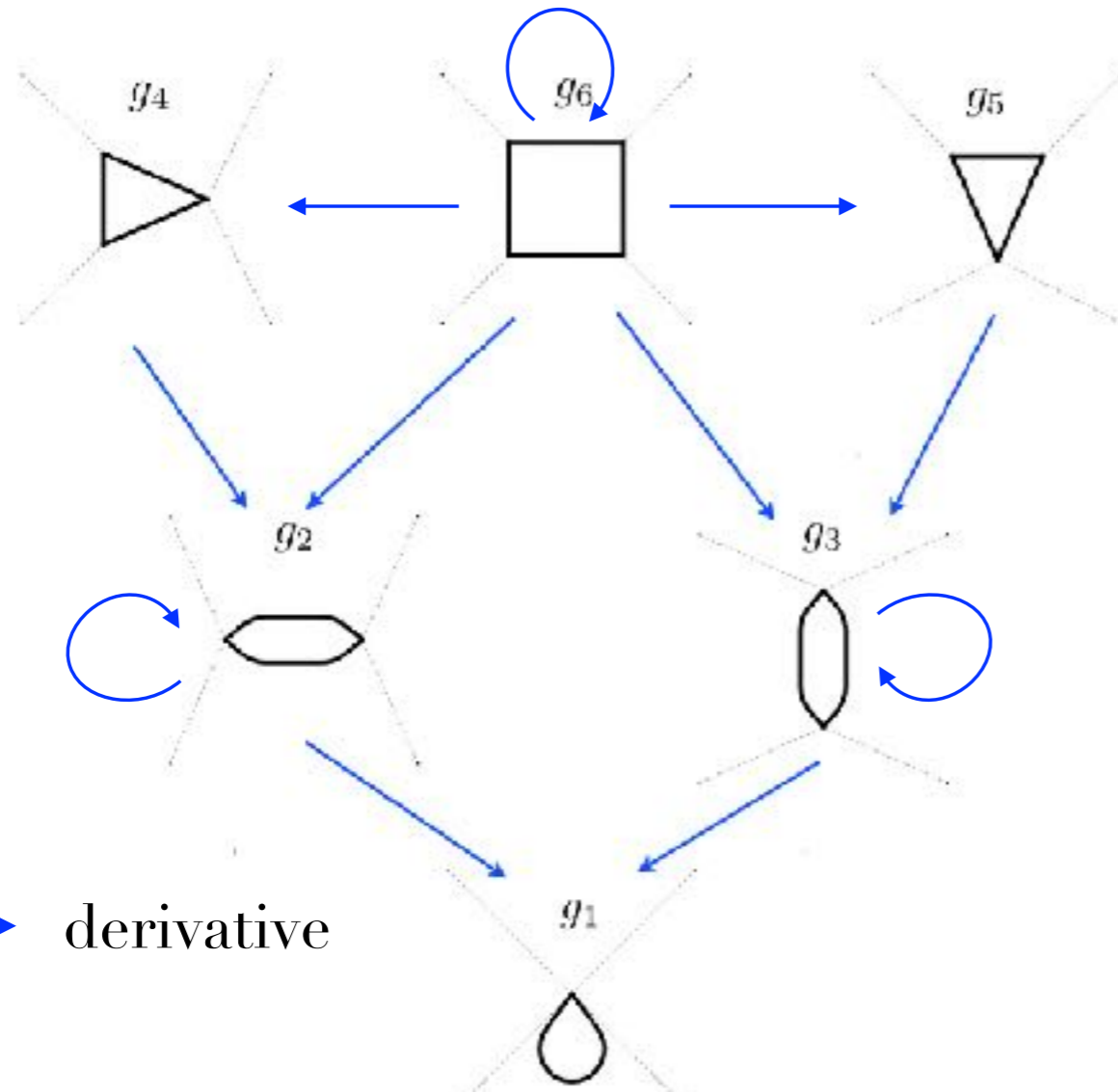
$$g_5 = t G[1, 1, 0, 1]$$

$$g_6 = \frac{1}{2} \sqrt{st(st - 4m^2(s + t))} G[1, 1, 1, 1]$$

UT basis (6 integrals)

$$d\vec{g} = \epsilon(dA)\vec{g}$$

# Finite UT integrals



$$g_1 = 2m^2 G[0, 0, 0, 3] / \epsilon^2$$

$$g_2 = -\sqrt{s(s - 4m^2)} G[1, 0, 2, 0] / \epsilon$$

$$g_3 = -\sqrt{t(t - 4m^2)} G[0, 1, 0, 2] / \epsilon$$

$$g_4 = s G[1, 1, 1, 0]$$

$$g_5 = t G[1, 1, 0, 1]$$

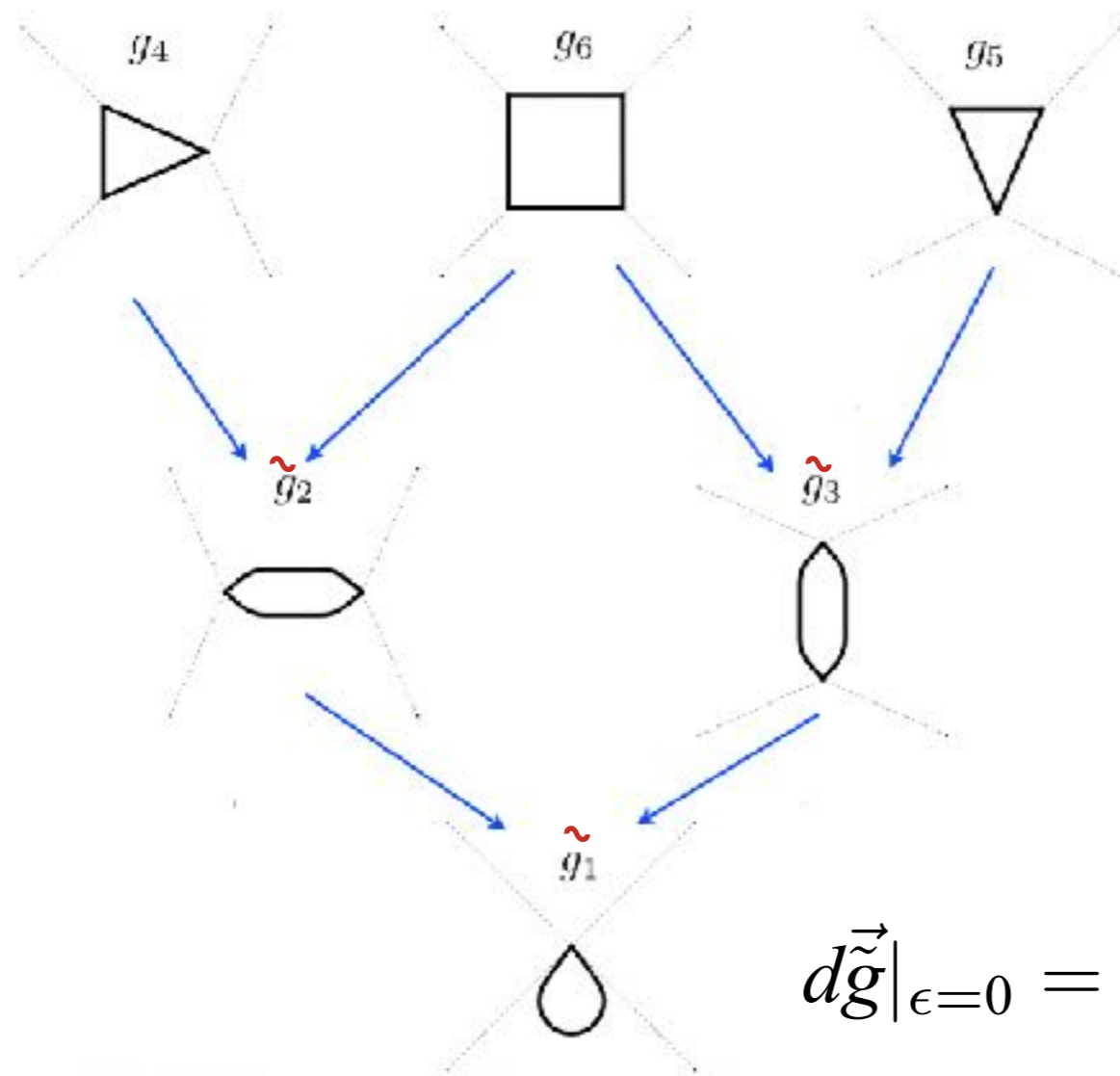
$$g_6 = \frac{1}{2} \sqrt{st(st - 4m^2(s + t))} G[1, 1, 1, 1]$$

$$d\vec{g} = \epsilon(dA)\vec{g}$$

→ derivative

$$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 2 \operatorname{ArcTanh}\left[\sqrt{1 - \frac{4m^2}{s}}\right] & \operatorname{Log}[m^2] - \operatorname{Log}[4m^2 - s] & 0 & 0 & 0 & 0 \\ 2 \operatorname{ArcTanh}\left[\sqrt{1 - \frac{4m^2}{t}}\right] & 0 & \operatorname{Log}[m^2] - \operatorname{Log}[4m^2 - t] & 0 & 0 & 0 \\ 0 & -2 \operatorname{ArcTanh}\left[\sqrt{1 - \frac{4m^2}{s}}\right] & 0 & 0 & 0 & 0 \\ 0 & 0 & -2 \operatorname{ArcTanh}\left[\sqrt{1 - \frac{4m^2}{t}}\right] & 0 & 0 & 0 \\ 0 & 4 \operatorname{ArcTanh}\left[\frac{\sqrt{(-4m^2 - s)t}}{\sqrt{s t - 4m^2(s+t)}}\right] & 4 \operatorname{ArcTanh}\left[\frac{\sqrt{s(-4m^2 - t)}}{\sqrt{s t - 4m^2(s+t)}}\right] & -4 \operatorname{ArcTanh}\left[\sqrt{1 - \frac{4m^2(s-t)}{st}}\right] & -4 \operatorname{ArcTanh}\left[\sqrt{1 - \frac{4m^2(s-t)}{st}}\right] & \operatorname{Log}[m^2] + \operatorname{Log}[s+t] - \operatorname{Log}[-st + 4m^2(s-t)] \end{pmatrix}$$

# Finite UT integrals with **truncation**



$$\tilde{g}_1 = 2m^2 G[0, 0, 0, 3] / \epsilon^2$$

$$\tilde{g}_2 = -\sqrt{s(s-4m^2)} G[1, 0, 2, 0] / \epsilon$$

$$\tilde{g}_3 = -\sqrt{t(t-4m^2)} G[0, 1, 0, 2] / \epsilon$$

$$g_4 = s G[1, 1, 1, 0]$$

$$g_5 = t G[1, 1, 0, 1]$$

$$g_6 = \frac{1}{2} \sqrt{st(st-4m^2(s+t))} G[1, 1, 1, 1]$$

$$d\vec{\tilde{g}}|_{\epsilon=0} = (d\tilde{A})\vec{\tilde{g}}|_{\epsilon=0}$$

$$\tilde{A} = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 2 \operatorname{ArcTanh}\left[\sqrt{1-\frac{4m^2}{s}}\right] & 0 & 0 & 0 & 0 \\ 2 \operatorname{ArcTanh}\left[\sqrt{1-\frac{4m^2}{t}}\right] & 0 & 0 & 0 & 0 \\ 0 & -2 \operatorname{ArcTanh}\left[\sqrt{1-\frac{4m^2}{s}}\right] & 0 & 0 & 0 \\ 0 & 0 & -2 \operatorname{ArcTanh}\left[\sqrt{1-\frac{4m^2}{t}}\right] & 0 & 0 \\ 0 & 4 \operatorname{ArcTanh}\left[\frac{\sqrt{(s-4m^2)t}}{\sqrt{(s-4m^2)(s+t)}}\right] & 4 \operatorname{ArcTanh}\left[\frac{\sqrt{(t-4m^2)s}}{\sqrt{(t-4m^2)(s+t)}}\right] & 0 & 0 \end{pmatrix}$$

All diagonal elements vanish, and the DE can be solved recursively.

# Summary

- (quasi)-finite integrals provides a simple truncated DE system
- Towards a systematical way of finding (quasi)-finite UT integrals
- graded syzygy to get a small IBP system for (quasi)-finite UT integrals
- working in progress ...

Vielen Dank