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 In SR and for classical field theory that play a role in physics, the positive total energy is the integral of a positive energy density ρ;

Responsible for ensuring the stability of the ground state

• In gravity, the situation is very different.

There is no satisfactory way to define the local energy density including gravity

 It is an old conjecture that this <u>total energy</u> is in fact always strictly positive, except for flat Minkowski space, which has zero energy



Positive energy theorem:

Under <u>dominant energy condition</u>, the <u>ADM</u> energy of any <u>asymptotically flat spacetime</u> solving the <u>Einstein equations</u> is non-negative, and zero precisely only for <u>Minkowski spacetime</u>.

• The first proof is given by Richard Schoen and Shing-Tung Yau:

Commun. Math. Phys. 65, 45 (1979); Commun. Math. Phys. 79, 231 (1981)

 An alternative proof for Einstein gravity, making crucial use of <u>spinor</u> algebra Edward Witten, Commun. Math. Phys. 80, 381 (1981)

 Gibbons, Hawking, Horowitz and Perry proved extensions of the theorem to asymptotically antide Sitter spacetimes and to Einstein–Maxwell theory.

 $M \ge \sqrt{Q^2 + P^2}$ Here Q is charge, P is magnetic charge

Commun. Math. Phys. 88 (3): 295–308



- An initial data site has an apparent horizon with area A_H and ADM mass E;
- The system then will evolve and settle down to a Kerr-Newman black hole;
- Assume that the dynamics is dominated by classical physics and null energy condition (NEC) is satisfied, then $A_H \leq A_{H,f}$
- For the final Kerr-Newman black hole,

$$A_{H,f} = 4\pi (r_h^2 + a^2), \qquad M_f = \frac{r_h^2 + a^2}{2r_h}$$

• We can see that $A_{H,f} \le 16\pi M_f^2$. Then we see $A_H \le 16\pi M_f^2 \le 16\pi M^2$

A gap: Only NEC is involved in above physical heuristic argument





• We have four different variables to describe the size of a black hole

Radius of horizon r₊

Radius of photon sphere r_{ph} Radius of shadow r_{sh}

• In Schwarzschild black hole, we find $\frac{3r_{+}}{2} = r_{ph} = \frac{r_{sh}}{\sqrt{3}} = 3M$



$$\frac{3r_+}{2} \le r_{ph} \le \frac{r_{sh}}{\sqrt{3}} \le 3M$$

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Particularly, we proved that if NEC is true, then

 $r_h < r_{ph} \le 3M$

This suggests NEC is enough to insure positive energy!



New theorem of "positive mass" on black hole

For asymptotically flat stationary black hole, if

- (1) Einstein equation and null energy condition are satisfied, and
- (2) the cross-section of event horizon has S^2 topology, then

$$3M \ge \frac{\kappa A}{4\pi} + \sqrt{\frac{A}{4\pi}}$$

• Particularly, a regular static spacetime has nonnegative ADM energy.



In static spherically symmetric case

• We consider the spherically symmetric spacetime.

$$ds^{2} = -f(r)e^{-\chi(r)}dt^{2} + \frac{dr^{2}}{f(r)} + r^{2}(d\theta^{2} + \sin^{2}\theta d\varphi^{2}) \qquad f(r) = 1 - \frac{2M}{r} + \cdots, \quad |\chi(r)| \to \mathcal{O}(1/r^{1+\alpha}), \quad r \to \infty$$

The energy momentum tensor has a form

$$f(r) = 1 - \frac{2M}{r} + \cdots, \quad |\chi(r)| \to \mathcal{O}(1/r^{1+\alpha}), \quad r \to \infty$$

- Surface gravity and area of horizon are given by $\kappa=rac{f'(r_h)e^{-\chi(r_h)/2}}{2}, \ \ \mathcal{A}=\pi r_h^2$

The Einsteins equation gives the following independent equations

$$\begin{aligned} f' &= (1 - 8\pi r^2 \rho - f)/r ,\\ \chi' &= -\frac{8\pi r}{f} (\rho + p_r) , & \text{NEC insures } \chi' \leq 0 & \longrightarrow \text{NEC insures } \chi \geq 0 \\ p'_r &= \frac{\rho + 4p_T - 3p_r}{2r} - \frac{(p_r + \rho)(8\pi p_r r^2 + 1)}{2rf} \end{aligned}$$



In static spherically symmetric case

Let us define a new quasi-local masses

$$|\chi(r)| \rightarrow \mathcal{O}(1/r^{1+\alpha}), r \rightarrow \infty$$

At infinity, they both give
us black hole mass

 $m_{\rm s}(\infty) = m_n(\infty) = M$

• At horizon, they are all positive

$$m_n(r_h) = \frac{\kappa \mathcal{A}}{12\pi} + \frac{r_h}{3}$$

 $m'_{n}(r) = \frac{8\pi}{3}e^{-\chi/2}r^{2}(\rho + p_{T}) + \frac{1}{3}[1 - e^{-\chi(r)/2}]$ Using Einstein equation, one can verify

NEC gives us theorem on positivity of black hole mass!



Beyond spherically symmetric case

For a static black hole, we can make ADM decomposition

$$\mathrm{d}s^2 = -N^2 \mathrm{d}t^2 + h_{ab} \mathrm{d}x^a \mathrm{d}x^b$$

We can foliate the equal t surface by $\{S_y\}$ and find a quasi-local mass m(y) on every surface S_y

- (1) $m(0) \leq m(\infty);$
- (2) At the initial surface we have $3m(0) = \sqrt{A/(4\pi)} + \kappa A/(4\pi);$
- (3) At the infinity $m(\infty) = E$;





Basic idea of proof

- Geroch, Pong Soo Jang and Wald once proposed a very important tool called "inverse mean curvature flow" to solve Penrose inequality;
- For a given initial surface, the flow is generated by the vector field $v^{\mu} = k^{-1}r^{\mu}$ and r^{μ} is the outward unit normal vector of the surface.
- The Hawking-Geroch mass of surface S_v is defined as

$$m_H(S_y) = \frac{\sqrt{A(S_y)}}{64\pi^{3/2}} \int_{S_y} (2\mathcal{R} - k^2) dS$$

Hawking-Geroch mass satisfies

$$m_H(0) = \frac{\sqrt{A(S_0)}}{64\pi^{3/2}} \int_{S_0} (2\Re - k^2) dS = \sqrt{\frac{A}{16\pi}} \qquad m_H(\infty) = E$$





Geroch, Wald and Jang found that

$$m_H(y) - m_H(0) \ge \int_{\Sigma_t} \rho Q \mathrm{d}V$$



Here Q is a positive function and will approach to 1 rapidly.

 $\mathrm{d}s^2 = -N^2 \mathrm{d}t^2 + h_{ab} \mathrm{d}x^a \mathrm{d}x^b$

We now introduce a scalar field W on according to following elliptic equation

$$D^{2}(NW) = 8\pi Q(\rho + T/2) \qquad W|_{\partial \Sigma_{t}} = 1$$

$$\partial_a (WN) \mathrm{d}S^a |_{\partial \Sigma_t} = \partial_a N \mathrm{d}S^a |_{\partial \Sigma_t}$$

We consider a new quasi-local mass

$$m_{K}(y) = \frac{1}{4\pi} \int_{S_{y}} \partial_{a} (NW) dS^{a} \longrightarrow m_{K}(\infty) - m_{K}(0) = 2 \int_{\Sigma_{t}} Q(\rho + T/2) dV$$
$$\lim_{y \to 0} \frac{1}{4\pi} \int_{S_{y}} \partial_{a} N dS^{a} = \frac{\kappa \mathcal{A}}{4\pi}, \quad \lim_{y \to \infty} \frac{1}{4\pi} \int_{S_{y}} \partial_{a} N dS^{a} = E \longrightarrow m_{K}(0) = \frac{\kappa \mathcal{A}}{4\pi}, \quad m_{K}(\infty) = E$$



Define the combination
$$m(r) := \frac{2}{3}m_H(r) + \frac{1}{3}m_K(r)$$

 $m_H(y) - m_H(0) \ge \int_{\Sigma_t} \rho Q dV$
 $m_K(\infty) - m_K(0) = 2 \int_{\Sigma_t} Q(\rho + T/2) dV$
 $m(\infty) - m(0) \ge \frac{1}{3} \int_{\Sigma_t} (4\rho + T)Q dV$
 MEC insures it nonnegative

We then conclude that NEC can insures the nonnegativity of ADM energy for static black hole!



Summary

• I conjecture there may be a universal inequality by only requiring NEC;

$$3E \ge \frac{\kappa A}{4\pi} + \sqrt{\frac{A}{4\pi}}, \qquad 2E \ge \sqrt{\frac{A}{4\pi}}$$

- A proof for static black hole that contains one connect compact horizon is given;
- This begs a new challenge for a famous settled question of general relativity:

In what general case can NEC replace DEC to insure nonnegative ADM energy?

