

Taub-NUT Blackholes with Additional Hair

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This talk is based on work:

Yu-Qi Chen and Hai-Shan Liu, arxiv: 2406.03692, to appear in SCPMA Hai-Shan Liu and Lei Zhang, arxiv: 2407.08208, JHEP 10 (2024) 067

Taub-NUT spacetimes

• New Taub-NUT black holes in higher derivative gravity theories

Taub-NUT spacetime

 Taub-NUT spacetime is a solution to Einstein's general relativity, first discovered by Taub in 1951, subsequently rediscovered by Newman, Tamburino and Unti in 1963.

• The solution is a simple generalization of the Schwarzschild spacetime. In addition to a Schwarzschild-like mass parameter m, it contains one more parameter which is called NUT parameter n.

• The nature of this NUT parameter n was viewed as a "gravitational magnetic charge" by some people, thus the solution is a sort of gravitational dyon and its Euclidean continuation is the solution known as Kaluza-Klein monopole.

Taub-NUT spacetime

- People pointed out that there may exist NUT charge in the real world.
- Recent works shows that CRO J165540 and M87 might contain NUT charge.

C.Chakraborty and S.Bhattacharyya, Phys.Rev.D98,043021(2018); C.Chakraborty and S.Bhattacharyya, JCAP 05,034(2019); M.Ghasemi-Nodehi and C.Chakraborty, EPJC 81,939(2021).

Primordial black holes could contain NUT charge, too.

C.Chakraborty and S.Bhattacharyya, Phys.Rev.D106,103028(2022); C.Chakraborty and B.Mukhopadhyay, EPJC 83,937(2023).

Taub-NUT spacetime

• The theory is just Einstein's general relativity

$$
\mathcal{L}=\sqrt{-g}\mathbf{R}
$$

- Taub-NUT metric is
- $ds^2 = -f(dt + 2n \cos\theta d\varphi)^2 + \frac{dr^2}{f} + (r^2 + n^2)(d\theta^2 + \sin^2\theta d\phi^2)$ $f = \frac{r^2 - 2mr - n^2}{r^2 + n^2}$

with two parameters m and n.

• This metric does not have curvature singularities, but has the so-called Minser string singularities at north and south pole ($\theta = 0, \pi$), where the metric fails to be invertible.

 $T_{\rm s}$

I: Black holes with spin-2 hair

Higher order curvature gravity

$$
L_4 = \sqrt{-g} \left(R + \beta R^2 + \alpha C^2 + \gamma GB \right)
$$

 Linear spectrum: one massive scalar mode, one massless gravity mode and one massive gravity mode.

$$
g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}; \quad h = \eta^{\mu\nu} h_{\mu\nu}, \quad h_{\mu\nu}^T = h_{\mu\nu} - \frac{1}{4} h \eta_{\mu\nu}
$$

$$
(\Box - \mu_0^2) h = 0, \qquad \Box(\Box - \mu_2^2) h_{\mu\nu}^T = 0 \quad , \qquad \mu_0^2 = \frac{1}{6\beta} \quad , \quad \mu_2^2 = -\frac{1}{2\alpha}
$$

 Asymptotic to Minkowski spacetime, the expansion of static spherical metric to linear order in large r is

$$
ds_4^2 = -hdt^2 + \frac{dr^2}{f} + r^2(d\theta^2 + \sin^2\theta \, d\phi^2),
$$

$$
h = 1 - \frac{m}{r} - \frac{c_0}{r}e^{-\mu_0 r} - \frac{c_2}{r}e^{-\mu_2 r} + \cdots, \qquad f = \dots
$$

The theory contains several modes. It doesn't guarantee the theory admits black holes that carry the corresponding hair.

• In four dimensional spacetimes, any Ricci-flat solution (e.g. Schwarzschild black hole) is still a solution of quadratic curvature gravity.

This raises a question: does there exist new black hole solution beyond Schwarzschild black hole in the high curvature gravity theory ?

Black hole solution with ghost massive spin-2 hair was constructed in Einstein-Weyl gravity($\beta = 0$).

Lu, Perkins, Pope and Stelle, Phys.Rev.Lett.114, 171601, 2015

Later, a new branch solution was constructed in Einstein-Weyl gravity($\beta = 0$) by Huang, Liu and Zhang. I: Black holes with spin-2 hair in high order curvature gravity

Huang, Liu and Zhang, JHEP 02 (2023) 057

• No new branch was discovered after these.

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• Are these all ?

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• Are these all?

 As we mentioned earlier, Taub-NUT black hole is also Ricci-flat.

No scalar hair theorem

• Theory

$$
L_4 = \sqrt{-g} (R + \beta R^2 + \alpha C^2)
$$

Taub-NUT ansatz

$$
ds_4^2 = -h(r)(dt + 2n\cos\theta d\phi)^2 + \frac{dr^2}{f(r)} + (r^2 + n^2)(d\theta^2 + \sin\theta^2 d\phi^2)
$$

• The trace of the equations of motion is

$$
\beta(\Box - m_0^2)R = 0
$$

 $m_0 = 1/\sqrt{6\beta}$ is the mass of scalar mode. The equation can be turned to $\int dx^4 \sqrt{g} R (\Box - m_0^2) R = \int dx^4 \Big[\partial_r (\sqrt{g} g^{rr} R \partial_r R) - \sqrt{g} (g^{rr} \partial_r R \partial_r R + m_0^2 R^2) \Big] = 0$ It can be seen that R must vanish $R = 0$.

• Does new Taub-NUT black hole exist? And if yes, where are they?

• The Lichnerowicz Mode analysis of the theory can give us an answer.

Making a tensor perturbation under Taub-NUT background. $\mathbf{1}$

$$
(\Delta_L + \frac{1}{2\alpha})\delta R_{\mu\nu} = 0
$$

$$
\Delta_L \delta R_{\mu\nu} = -\Box \delta R_{\mu\nu} - 2R_{\mu\nu\rho\sigma} \delta R^{\rho\sigma}
$$

• And $\delta R_{\mu\nu}$ is a transverse and traceless tensor.

$$
g^{\mu\nu}\delta R_{\mu\nu}=0~~\nabla^{\mu}\delta R_{\mu\nu}=0
$$

• It implies that once the Lichnerowicz operator has a transverse traceless eigenfunction with eigenvalue $\lambda = -$ 1 2α , then there exists a linearized perturbation away from the Taub-NUT solution.

The Lichnerowicz equation can not be solved analytically, but it can be solve in the large r, which gives the Yukawa like mode

• The equation can be solved numerically, the relation between r_0 and NUT parameter n is shown for $\alpha = 1/2$.

 $e^{\pm\sqrt{-\lambda}r}$

It can be seen that when $n=0, r_0 = 0.876$, which covers the Schwarzschild case

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The Lichnerowicz is very important for us to find new black hole numerically.

First, it can tell us if a black hole exists.

Second, it can tell us where the new black hole will emerge.

• Theory

$$
L_4 = \sqrt{-g} \ (R + \alpha \ C^2)
$$

Taub-NUT ansatz

$$
ds_4^2 = -h(r)(dt + 2n\cos\theta d\phi)^2 + \frac{dr^2}{f(r)} + (r^2 + n^2)(d\theta^2 + \sin\theta^2 d\phi^2)
$$

• Near horizon r_0 , the expansions are

$$
f(r) = f_1(r - r_0) + f_2(r - r_0)^2 + f_3(r - r_0)^3 + \dots
$$

$$
h(r) = h_1(r - r_0) + h_2(r - r_0)^2 + h_3(r - r_0)^3 + \dots
$$

$$
h(r) = h_1(r - r_0) + h_2(r - r_0)^2 + h_3(r - r_0)^3 + \dots
$$

with

$$
f_2 = \frac{-3n^2[1 + 4\alpha f_1(f_1 - h_1) - f_1r_0]}{8\alpha f_1r_0(r_0^2 + n^2)} + \frac{8\alpha f_1(1 - 2f_1r_0) + 3r_0(f_1r_0 - 1)}{8\alpha f_1(r_0^2 + n^2)}
$$

$$
h_2 = \frac{h_1n^2[1 + 4\alpha f_1(f_1 - h_1) - f_1r_0]}{8\alpha f_1^2r_0(r_0^2 + n^2)} + \frac{8\alpha f_1(-1 + 2f_1r_0) + r_0(f_1r_0 - 1)}{8\alpha f_1^2(r_0^2 + n^2)}
$$

• Numerically, we shall integrate from horizon to infinity, using shooting method to get rid of the diverging mode $e^{+\sqrt{-\lambda}r}$ and keep the converging mode $e^{-\sqrt{-\lambda}r}$.

Metric profile with different NUT parameter n

• $\alpha = 0.5, r_0 = 1$

Critical line of m=0

• Temperature and entropy

$$
T = \frac{\sqrt{f'(r_0)h'(r_0)}}{4\pi} = \frac{\sqrt{1+r_0\delta}}{4\pi r_0}
$$

$$
S = -\frac{1}{8} \int \sqrt{h} d\Omega \epsilon_{ab} \epsilon_{cd} \frac{\partial L}{\partial R_{abcd}} = \pi (r_0^2 + n^2) - 4\pi \alpha r_0 \delta
$$

 1.0

Radius (ro)

 1.2

 $n = 0.5$

 1.4

II scalarization: Black holes with scalar hair

• No hair theorem: it is well-know that Einstein gravity minimally coupled to massless scalar can not supports a black hole with non trivial scalar hair.

$$
L_4 = \sqrt{g} \left(R - \frac{1}{2} (\partial \phi)^2 \right)
$$

• The integration $\int dx^4 \sqrt{g} \phi \nabla_{\mu} \nabla^{\mu} \phi = - \int dx^4 \sqrt{g} (\nabla^{\mu} \phi)^2$ EOM for scalar is

$$
\nabla_{\!\mu}\nabla^{\mu}\phi=0
$$

so

$$
\phi = 0
$$

Scalarization in EsGB

 Recently, a way of scalariztion of black holes by using Gauss-Bonnet terms was carried out.

> Antoniou,Bakopoulos, Kanti, Phys.Rev.Lett.120,131102,2018; Doneva and Yazadjiev, Phys.Rev.Lett.120,131103,2018; Silva, Sakstein, Gualtieri,Sotiriou, Berti, Phys.Rev.Lett.120,131104,2018.

• It was shown that Einstein-scalar theory can excite new scalar hairy black hole by adding non-minimally coupled term

$$
L_4 = \sqrt{g} \left(R - \frac{1}{2} (\partial \phi)^2 \right) + \frac{\lambda^2}{12} (1 - e^{-6\phi^2})GB
$$
\n
$$
\begin{array}{c}\n \stackrel{\text{a.s.}}{=} \overbrace{\left(\begin{array}{c}\n \stackrel{\text{a.s.}}{=} \text{Bessel branch} \\
\stackrel{\text{b.s.}}{=} \text{Bessel branch} \\
\stackrel{\text{c.s.}}{=} \text{Bessel branch}\n \end{array}\right)}\n \stackrel{\text{a.s.}}{=} \overbrace{\left(\begin{array}{c}\n \stackrel{\text{b.s.}}{=} \text{Bessel branch} \\
\stackrel{\text{c.s.}}{=} \text{Bessel branch}\n \end{array}\right)}\n \stackrel{\text{b.s.}}{=} \overbrace{\left(\begin{array}{c}\n \stackrel{\text{d.s.}}{=} \text{Bessel branch} \\
\stackrel{\text{d.s.}}{=} \text{Bessel branch}\n \end{array}\right)}\n \stackrel{\text{d.s.}}{=} \overbrace{\left(\begin{array}{c}\n \stackrel{\text{d.s.}}{=} \text{Bessel graph} \\
\stackrel{\text{d.s.}}{=} \text{Bessel graph}\n \end{array}\right)}\n \stackrel{\text{d.s.}}{=} \overbrace{\left(\begin{array}{c}\n \stackrel{\text{d.s.}}{=} \text{Bessel graph} \\
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\stackrel{\text{d.s.}}{=} \text{Bessel graph}\n \end{array}\right)}\n \stackrel{\text{d.s.}}{=} \overbrace{\left(\begin{array}{c}\n \stackrel{\text{d.s.}}{=} \text{Bessel graph} \\
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$$

Scalarization of Taub-NUT

Theory

$$
L_4 = \sqrt{g} \left(R - \frac{1}{2} (\partial \phi)^2 \right) + \frac{\lambda^2}{12} (1 - e^{-6\phi^2}) GB
$$

Test scalar under Taub-NUT background

$$
\nabla_\alpha\nabla^\alpha\delta\varphi+\frac{\lambda^2}{4}R_{GB}^2\delta\varphi=0
$$

Performing a spherical harmonics of the scalar, the radial equation is

$$
\frac{1}{(r^2+n^2)}\frac{d}{dr}\left[(r^2-2mr-n^2)\frac{dU_l}{dr}\right] - \left[\frac{l(l+1)}{n^2+r^2} - \frac{\lambda^2}{4}R_{GB}^2\right]U_l = 0.
$$

It becomes an eigenvalue problem, fixing (n, λ), the equation can admit solutions for a discrete parameter m.

Scalarization of Taub-NUT

The equation can be solved numerically, which gives

 This analysis shows the Taub-NUT black hole is not stable in these points, where new black holes with scalar hair may emerge.

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For Taub-NUT ansatz, the equations of motion are very complicated.

$$
4fh^2r\left[2n^2\Psi_r + r(n^2+r^2)\Psi'_r\right] - 3n^2f^2(n^2+r^2)\left(-1+2\Psi_r h' + 4h\Psi'_r\right) -\nfh(n^2+r^2)^3\varphi'^2 - fh(n^2+r^2)\left[2n^2+r^2-6r^2\Psi_r h' + 4(n^2+r^2)\Psi'_r\right] +\nf(n^2+r^2)^2\left[1-(r+2\Psi_r)h'\right] = 0,
$$

$$
4f^{2}h\left[n^{2}+(n^{2}+r^{2})^{2}\varphi'^{2}\right]-h(n^{2}+r^{2})(n^{2}+r^{2}-4hr\Psi_{r})f'^{2}+2f^{2}r(n^{2}+r^{2})h'-4n^{2}f^{3}\left(2\Psi_{r}h^{'}+4h\Psi_{r}^{'}-1\right)+f(n^{2}+r^{2})^{2}f^{'}h^{'}+2fhr(n^{2}+r^{2})\left(1-6\Psi_{r}^{'}h^{'}\right)f^{'}+2fh(n^{2}+r^{2})^{2}f''-8fh^{2}r(n^{2}+r^{2})\Psi_{r}f''-8fh^{2}\left[n^{2}\Psi_{r}+r(n^{2}+r^{2})\Psi_{r}^{'}\right]f^{'}=0,
$$

$$
fh(n^{2} + r^{2})^{4}\varphi' f' + f^{2} (n^{2} + r^{2})^{4} \varphi' h' + 2f^{2}h (n^{2} + r^{2})^{2} [2r(n^{2} + r^{2})\varphi' + (n^{2} + r^{2})^{2}\varphi''] +
$$

$$
\left\{ fh(n^{2} + r^{2})^{2} [3r^{2}\lambda^{2}h' f' - 2\lambda^{2}(n^{2} + r^{2})f''] - 3n^{2}\lambda^{2}ff' (n^{2} + r^{2})^{2} (f'h + fh') +
$$

$$
4n^{2}\lambda^{2}f^{3} [2h (n^{2} - 3r^{2}) + r (n^{2} + r^{2})h'] + \lambda^{2}h(n^{2} + r^{2})(n^{2} + r^{2} - r^{2}h)f' +
$$

$$
2\lambda^{2}fh^{2}r(n^{2} + r^{2}) [2n^{2}f' + r(n^{2} + r^{2})f''] - \lambda^{2}f (n^{2} + r^{2})^{3}f'h' +
$$

$$
24n^{2}rf^{2}h(n^{2} + r^{2})^{2}f' - 6n^{2}\lambda^{2}f^{2}h (n^{2} + r^{2})^{2}f' \right\} \frac{dF(\varphi)}{d\varphi} = 0,
$$

$$
(n^{2} + r^{2})h[r(n^{2} + r^{2}) + 2(n^{2} + r^{2} - 3r^{2}h)\Psi_{r}]f' - n^{2}f^{2}(n^{2} + r^{2} + 8rh\Psi_{r}) - (n^{2} + r^{2})^{2}f - (n^{2} + r^{2})fh[-r^{2} - 6n^{2}f'\Psi_{r} + (n^{2} + r^{2})^{2}\varphi'^{2}] = 0,
$$

Asymptotic form in the horizon

$$
h(r) = \frac{1+\delta}{r_h} (r - r_h) + \mathcal{O}(r - r_h)^2,
$$

\n
$$
f(r) = \frac{r_h \left[e^{\left(-12\varphi_h^2\right)} \left(n^2 + r_h^2\right) \delta + 6\varphi_h^2 (1+\delta)^2 \lambda^4 \right]}{6n^2 \varphi_h^2 (1+\delta) \lambda^4} (r - r_h) + \mathcal{O}(r - r_h)^2,
$$

\n
$$
\varphi(r) = \varphi_h - \frac{e^{6\varphi_h^2} r_h \delta}{2\varphi_h (1+\delta) \lambda^2} (r - r_h) + \mathcal{O}(r - r_h)^2.
$$

• Asymptotic form in the infinity

$$
h(r) = 1 - \frac{2M}{r} + \frac{D^2 - 2n^2}{r^2} + \frac{M(D^2 + 2n^2)}{r^3} + \frac{(4D^2M^2 + 6n^4 + 2n^2D^2)}{3r^4} + \frac{M}{12r^5} \{63D^4 - 4D^2 (142M^2 + 65n^2 - 4\lambda^2) + 8 (80M^4 + 12n^2 - 51n^2\lambda^2) + 8M^2 (80n^2 - 21\lambda^2) \} + \dots, f(r) = 1 - \frac{2M}{r} - \frac{2n^2}{r^2} + \frac{M(D^2 + 6n^2)}{3r^3} + \frac{2D^2M^2 + 2n^2(D^2 + 3n^2)}{3r^4} + \frac{M}{60r^5} \{55D^4 - 4D^2 (130M^2 + 50n^2 - 28\lambda^2) + 8 (80M^4 - 51n^2\lambda^2) + 8M^2 (80n^2 - 2\lambda^2) \} + \dots, \n\varphi(r) = \frac{D}{r} + \frac{DM}{r^2} - \frac{D[D^2 - 2(4M^2 + n^2)]}{6r^3} + \frac{DM (6M^2 - 2D^2 + 3n^2)}{3r^4} + \frac{D}{120r^5} \{9D^4 + 24 [16M^4 + n^4 + 3n^2\lambda^2 + 3M^2 (4n^2 - \lambda^2)] - 4D^2 (58M^2 + 13n^2) \} + \dots
$$

The equations can be solved numerically

The entropy of the black hole

 The entropy of scalarized black hole is universal at the bifurcation point for different NUT parameter, and it is the highest value of the scalarized black hole.

 $S_{\text{scalar-free}}(m, n) < S_{\text{scalarized}}(m, n) \leq S_{\text{scalarized}}(m_{\text{max}}(n)) = S_{\text{scalar-free}}(m_{\text{max}}(n))$

 The universal value at bifurcation points shows a relation between m and n

The universal value only holds for positive mass parameter.

Conclusion and outlook

 We constructed new Taub-NUT black holes with scalar hair or spin-2 hair in higher derivative gravity theories.

 The thermodynamics, especially the definition of mass and NUT charge for new Taub-NUT black hole in higher derivative gravity is worth further investigation.

Thank you!