

Taub-NUT Blackholes with Additional Hair

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This talk is based on work:

Yu-Qi Chen and Hai-Shan Liu, arxiv: 2406.03692, to appear in SCPMA Hai-Shan Liu and Lei Zhang, arxiv: 2407.08208, JHEP 10 (2024) 067



• Taub-NUT spacetimes

• New Taub-NUT black holes in higher derivative gravity theories

Taub-NUT spacetime

 Taub-NUT spacetime is a solution to Einstein's general relativity, first discovered by Taub in 1951, subsequently rediscovered by Newman, Tamburino and Unti in 1963.

• The solution is a simple generalization of the Schwarzschild spacetime. In addition to a Schwarzschild-like mass parameter m, it contains one more parameter which is called NUT parameter n.

• The nature of this NUT parameter n was viewed as a "gravitational magnetic charge" by some people, thus the solution is a sort of gravitational dyon and its Euclidean continuation is the solution known as Kaluza-Klein monopole.

Taub-NUT spacetime

- People pointed out that there may exist NUT charge in the real world.
- Recent works shows that CRO J165540 and M87 might contain NUT charge.

C.Chakraborty and S.Bhattacharyya, Phys.Rev.D98,043021(2018); C.Chakraborty and S.Bhattacharyya, JCAP 05,034(2019); M.Ghasemi-Nodehi and C.Chakraborty, EPJC 81,939(2021).

Primordial black holes could contain NUT charge, too.

C.Chakraborty and S.Bhattacharyya, Phys.Rev.D106,103028(2022); C.Chakraborty and B.Mukhopadhyay, EPJC 83,937(2023).

Taub-NUT spacetime

• The theory is just Einstein's general relativity

$$\mathcal{L} = \sqrt{-g}\mathbf{R}$$

Taub-NUT metric is

 $ds^{2} = -f(dt + 2n\cos\theta d\varphi)^{2} + \frac{dr^{2}}{f} + (r^{2} + n^{2})(d\theta^{2} + \sin^{2}\theta d\phi^{2})$ $f = \frac{r^{2} - 2mr - n^{2}}{r^{2} + n^{2}}$

with two parameters m and n.

• This metric does not have curvature singularities, but has the so-called Minser string singularities at north and south pole ($\theta = 0, \pi$), where the metric fails to be invertible.

 T_{s}

I: Black holes with spin-2 hair

Higher order curvature gravity

$$L_4 = \sqrt{-g} (R + \beta R^2 + \alpha C^2 + \gamma GB)$$

• Linear spectrum: one massive scalar mode, one massless gravity mode and one massive gravity mode.

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}; \quad h = \eta^{\mu\nu}h_{\mu\nu}, \quad h_{\mu\nu}^T = h_{\mu\nu} - \frac{1}{4}h\eta_{\mu\nu}$$
$$(\Box - \mu_0^2)h = 0, \qquad \Box(\Box - \mu_2^2)h_{\mu\nu}^T = 0 \quad , \qquad \mu_0^2 = \frac{1}{6\beta}, \quad \mu_2^2 = -\frac{1}{2\alpha}$$

• Asymptotic to Minkowski spacetime, the expansion of static spherical metric to linear order in large r is

$$ds_4^2 = -hdt^2 + \frac{dr^2}{f} + r^2(d\theta^2 + Sin^2\theta \ d\phi^2),$$

$$h = 1 - \frac{m}{r} - \frac{c_0}{r}e^{-\mu_0 r} - \frac{c_2}{r}e^{-\mu_2 r} + \cdots, \qquad f = \dots$$

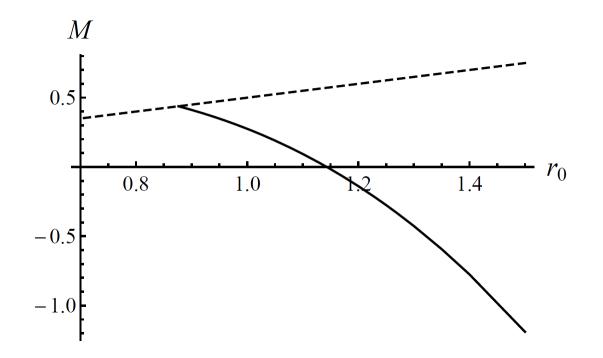
The theory contains several modes. It doesn't guarantee the theory admits black holes that carry the corresponding hair.

• In four dimensional spacetimes, any Ricci-flat solution (e.g. Schwarzschild black hole) is still a solution of quadratic curvature gravity.

• This raises a question: does there exist new black hole solution beyond Schwarzschild black hole in the high curvature gravity theory ?

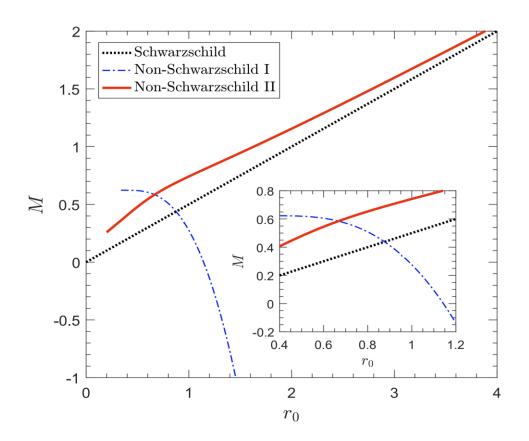
Black hole solution with ghost massive spin-2 hair was constructed in Einstein-Weyl gravity($\beta = 0$).

Lu, Perkins, Pope and Stelle, Phys.Rev.Lett.114, 171601, 2015



I: Black holes with spin-2 hair in high order curvature gravity Later, a new branch solution was constructed in Einstein-Weyl gravity($\beta = 0$) by Huang, Liu and Zhang.

Huang, Liu and Zhang, JHEP 02 (2023) 057



No new branch was discovered after these.

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• Are these all ?

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• Are these all ?

 As we mentioned earlier, Taub-NUT black hole is also Ricci-flat.

No scalar hair theorem

Theory

$$L_4 = \sqrt{-g} \left(R + \beta \ R^2 + \alpha \ C^2 \right)$$

Taub-NUT ansatz

$$ds_4^2 = -h(r)(dt + 2n\cos\theta d\phi)^2 + \frac{dr^2}{f(r)} + (r^2 + n^2)(d\theta^2 + \sin\theta^2 d\phi^2)$$

• The trace of the equations of motion is

$$\beta(\Box - m_0^2)R = 0$$

 $m_0 = 1/\sqrt{6\beta}$ is the mass of scalar mode. The equation can be turned to $\int dx^4 \sqrt{g} R(\Box - m_0^2) R = \int dx^4 \Big[\partial_r (\sqrt{g} g^{rr} R \partial_r R) - \sqrt{g} (g^{rr} \partial_r R \partial_r R + m_0^2 R^2) \Big] = 0$ It can be seen that R must vanish R = 0.

• Does new Taub-NUT black hole exist ? And if yes, where are they?

• The Lichnerowicz Mode analysis of the theory can give us an answer.

•Making a tensor perturbation under Taub-NUT background.

$$(\Delta_L + \frac{1}{2\alpha})\delta R_{\mu\nu} = 0$$
$$\Delta_L \delta R_{\mu\nu} = -\Box \delta R_{\mu\nu} - 2R_{\mu\nu\rho\sigma}\delta R^{\rho\sigma}$$

• And $\delta R_{\mu\nu}$ is a transverse and traceless tensor.

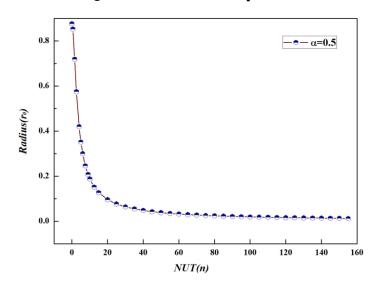
$$g^{\mu\nu}\delta R_{\mu\nu} = 0 \quad \nabla^{\mu}\delta R_{\mu\nu} = 0$$

• It implies that once the Lichnerowicz operator has a transverse traceless eigenfunction with eigenvalue $\lambda = -\frac{1}{2 \alpha}$, then there exists a linearized perturbation away from the Taub-NUT solution.

•The Lichnerowicz equation can not be solved analytically, but it can be solve in the large r, which gives the Yukawa like mode

• The equation can be solved numerically, the relation between r_0 and NUT parameter n is shown for $\alpha = 1/2$.

 $e^{\pm \sqrt{-\lambda r}}$

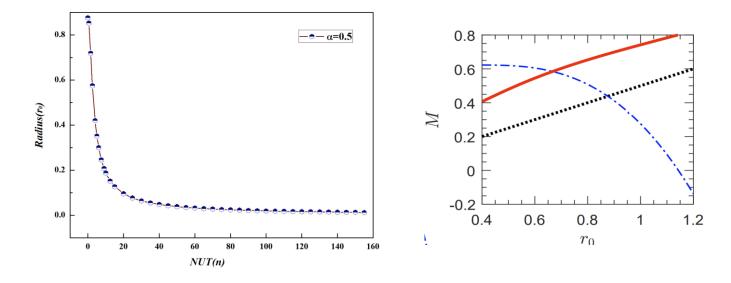


• It can be seen that when $n=0, r_0 = 0.876$, which covers the Schwarzschild case

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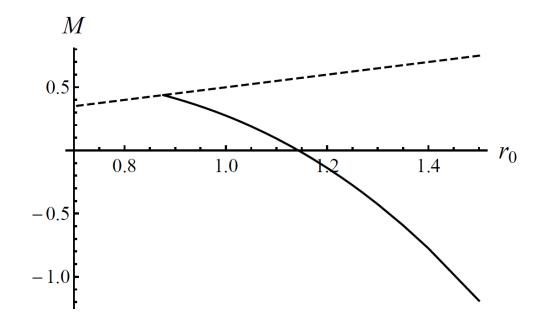


• It can be seen that when $n=0, r_0 = 0.876$, which covers the Schwarzschild case

•The Lichnerowicz is very important for us to find new black hole numerically.

•First, it can tell us if a black hole exists.

•Second, it can tell us where the new black hole will emerge.





Theory

$$L_4 = \sqrt{-g} \left(R + \alpha \ C^2 \right)$$

• Taub-NUT ansatz

$$ds_4^2 = -h(r)(dt + 2n\cos\theta d\phi)^2 + \frac{dr^2}{f(r)} + (r^2 + n^2)(d\theta^2 + \sin\theta^2 d\phi^2)$$

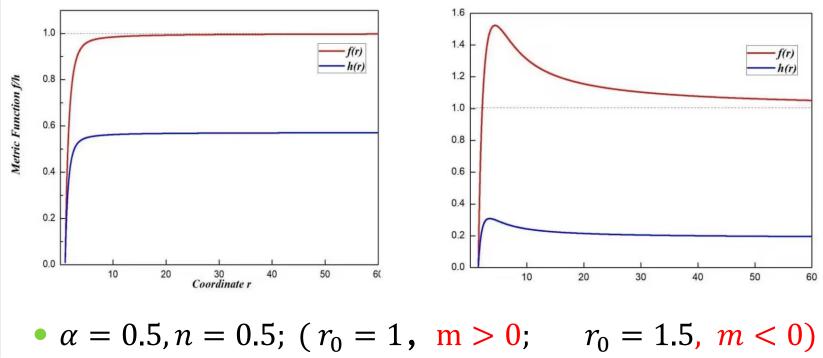
• Near horizon r_0 , the expansions are

$$f(r) = f_1(r - r_0) + f_2(r - r_0)^2 + f_3(r - r_0)^3 + \dots$$
$$h(r) = h_1(r - r_0) + h_2(r - r_0)^2 + h_3(r - r_0)^3 + \dots$$
with

$$f_{2} = \frac{-3n^{2}[1 + 4\alpha f_{1}(f_{1} - h_{1}) - f_{1}r_{0}]}{8\alpha f_{1}r_{0}(r_{0}^{2} + n^{2})} + \frac{8\alpha f_{1}(1 - 2f_{1}r_{0}) + 3r_{0}(f_{1}r_{0} - 1)}{8\alpha f_{1}(r_{0}^{2} + n^{2})}$$
$$h_{2} = \frac{h_{1}n^{2}[1 + 4\alpha f_{1}(f_{1} - h_{1}) - f_{1}r_{0}]}{8\alpha f_{1}^{2}r_{0}(r_{0}^{2} + n^{2})} + \frac{8\alpha f_{1}(-1 + 2f_{1}r_{0}) + r_{0}(f_{1}r_{0} - 1)}{8\alpha f_{1}^{2}(r_{0}^{2} + n^{2})}$$

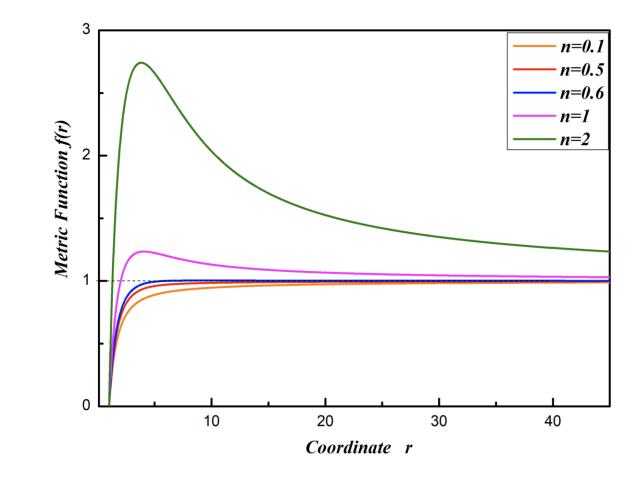


• Numerically, we shall integrate from horizon to infinity, using shooting method to get rid of the diverging mode $e^{+\sqrt{-\lambda r}}$ and keep the converging mode $e^{-\sqrt{-\lambda r}}$.





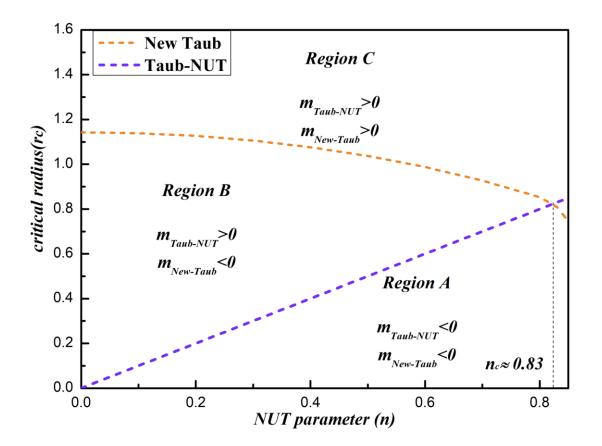
Metric profile with different NUT parameter n



• $\alpha = 0.5$, $r_0 = 1$



Critical line of m=0





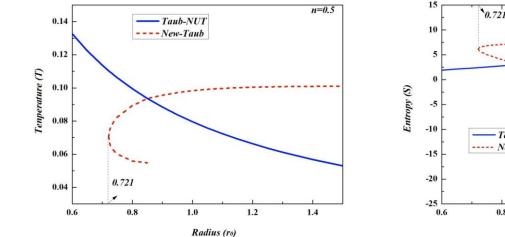
Temperature and entropy

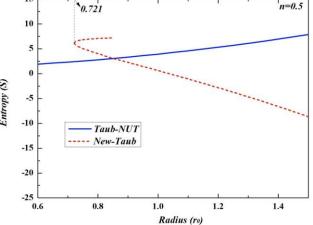
$$T = \frac{\sqrt{f'(r_0)h'(r_0)}}{4\pi} = \frac{\sqrt{1+r_0\delta}}{4\pi r_0}$$

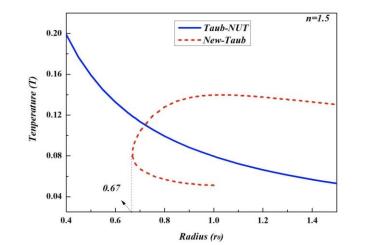
$$S = -\frac{1}{8} \int \sqrt{h} d\Omega \epsilon_{ab} \epsilon_{cd} \frac{\partial L}{\partial R_{abcd}} = \pi (r_0^2 + n^2) - 4\pi \alpha r_0 \delta$$

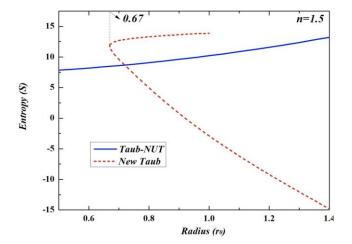












II scalarization: Black holes with scalar hair

• No hair theorem: it is well-know that Einstein gravity minimally coupled to massless scalar can not supports a black hole with non trivial scalar hair.

$$L_4 = \sqrt{g} \left(R - \frac{1}{2} (\partial \phi)^2 \right)$$

 The integration
 ∫ dx⁴√gφ∇_μ∇^μφ = −∫ dx⁴√g(∇^μφ)²
 EOM for scalar is

$$\nabla_{\!\mu}\nabla^{\mu}\phi=0$$

SO

$$\phi = 0$$

Scalarization in EsGB

 Recently, a way of scalarization of black holes by using Gauss-Bonnet terms was carried out.

> Antoniou, Bakopoulos, Kanti, Doneva and Yazadjiev, Silva, Sakstein, Gualtieri, Sotiriou, Berti,

Phys.Rev.Lett. 120, 131102,2018; Phys.Rev.Lett. 120, 131103,2018; Phys.Rev.Lett. 120, 131104,2018.

• It was shown that Einstein-scalar theory can excite new scalar hairy black hole by adding non-minimally coupled term

$$L_{4} = \sqrt{g} \left(R - \frac{1}{2} (\partial \phi)^{2} \right) + \frac{\lambda^{2}}{12} (1 - e^{-6\phi^{2}}) GE$$

$$\int_{2}^{6} \int_{2}^{6} \frac{1}{\sqrt{2}} (\partial \phi)^{2} + \frac{\lambda^{2}}{12} (1 - e^{-6\phi^{2}}) GE$$

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$$\int_{2}^{6} \frac{1}{\sqrt{2}} (\partial \phi)^{2} + \frac{\lambda^{2}}{\sqrt{2}} (\partial \phi)^{2} + \frac{\lambda^{$$

Scalarization of Taub-NUT

Theory

$$L_4 = \sqrt{g} \left(R - \frac{1}{2} (\partial \phi)^2 \right) + \frac{\lambda^2}{12} (1 - e^{-6\phi^2}) GB$$

Test scalar under Taub-NUT background

$$\nabla_{\alpha}\nabla^{\alpha}\delta\varphi + \frac{\lambda^2}{4}R_{GB}^2\delta\varphi = 0$$

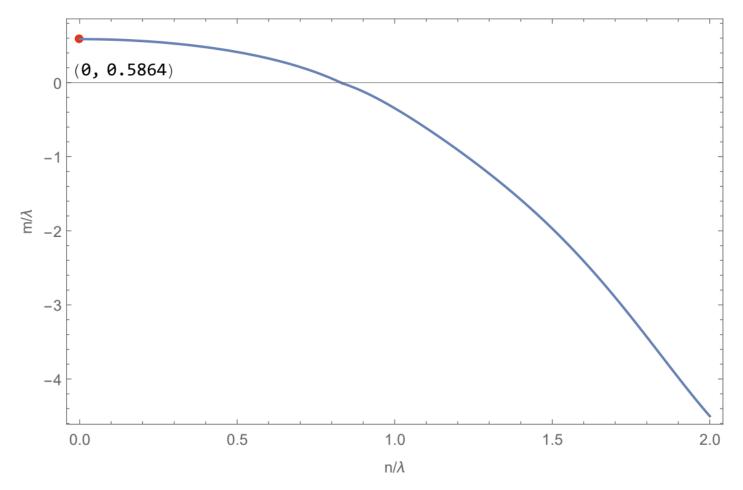
• Performing a spherical harmonics of the scalar, the radial equation is

$$\frac{1}{(r^2+n^2)}\frac{d}{dr}\left[\left(r^2-2mr-n^2\right)\frac{dU_l}{dr}\right] - \left[\frac{l(l+1)}{n^2+r^2} - \frac{\lambda^2}{4}R_{GB}^2\right]U_l = 0.$$

It becomes an eigenvalue problem, fixing (n, λ), the equation can admit solutions for a discrete parameter m.

Scalarization of Taub-NUT

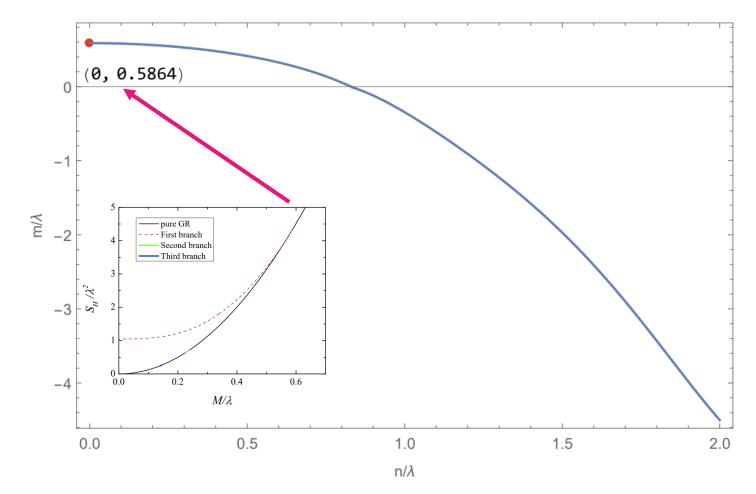
The equation can be solved numerically, which gives



• This analysis shows the Taub-NUT black hole is not stable in these points, where new black holes with scalar hair may emerge.

Scalarization of Taub-NUT

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• This analysis shows the Taub-NUT black hole is not stable in these points, where new black holes with scalar hair may emerge.

• For Taub-NUT ansatz, the equations of motion are very complicated.

$$4fh^{2}r\left[2n^{2}\Psi_{r}+r(n^{2}+r^{2})\Psi_{r}'\right] - 3n^{2}f^{2}(n^{2}+r^{2})\left(-1+2\Psi_{r}h'+4h\Psi_{r}'\right) - fh(n^{2}+r^{2})^{3}\varphi'^{2} - fh(n^{2}+r^{2})\left[2n^{2}+r^{2}-6r^{2}\Psi_{r}h'+4(n^{2}+r^{2})\Psi_{r}'\right] + f(n^{2}+r^{2})^{2}\left[1-(r+2\Psi_{r})h'\right] = 0,$$

$$\begin{aligned} 4f^{2}h\left[n^{2}+(n^{2}+r^{2})^{2}\varphi'^{2}\right]-h(n^{2}+r^{2})(n^{2}+r^{2}-4hr\Psi_{r})f'^{2}+2f^{2}r(n^{2}+r^{2})h'-\\ 4n^{2}f^{3}\left(2\Psi_{r}h'+4h\Psi_{r}'-1\right)+f(n^{2}+r^{2})^{2}f'h'+2fhr(n^{2}+r^{2})\left(1-6\Psi_{r}'h'\right)f'+\\ 2fh(n^{2}+r^{2})^{2}f''-8fh^{2}r(n^{2}+r^{2})\Psi_{r}f''-8fh^{2}\left[n^{2}\Psi_{r}+r(n^{2}+r^{2})\Psi_{r}'\right]f'=0, \end{aligned}$$

$$\begin{split} fh(n^2+r^2)^4\varphi'f' + f^2\left(n^2+r^2\right)^4\varphi'h' + 2f^2h\left(n^2+r^2\right)^2\left[2r(n^2+r^2)\varphi' + (n^2+r^2)^2\varphi''\right] + \\ \left\{fh(n^2+r^2)^2\left[3r^2\lambda^2h'f' - 2\lambda^2(n^2+r^2)f''\right] - 3n^2\lambda^2ff'\left(n^2+r^2\right)^2\left(f'h+fh'\right) + \\ 4n^2\lambda^2f^3\left[2h\left(n^2-3r^2\right) + r\left(n^2+r^2\right)h'\right] + \lambda^2h(n^2+r^2)(n^2+r^2-r^2h)f' + \\ 2\lambda^2fh^2r(n^2+r^2)\left[2n^2f' + r(n^2+r^2)f''\right] - \lambda^2f\left(n^2+r^2\right)^3f'h' + \\ 24n^2rf^2h(n^2+r^2)^2f' - 6n^2\lambda^2f^2h\left(n^2+r^2\right)^2f''\right\}\frac{dF(\varphi)}{d\varphi} = 0, \end{split}$$

$$(n^{2} + r^{2})h\left[r(n^{2} + r^{2}) + 2(n^{2} + r^{2} - 3r^{2}h)\Psi_{r}\right]f' - n^{2}f^{2}(n^{2} + r^{2} + 8rh\Psi_{r}) - (n^{2} + r^{2})^{2}f - (n^{2} + r^{2})fh\left[-r^{2} - 6n^{2}f'\Psi_{r} + (n^{2} + r^{2})^{2}\varphi'^{2}\right] = 0,$$

• Asymptotic form in the horizon

$$h(r) = \frac{1+\delta}{r_h} (r-r_h) + \mathcal{O}(r-r_h)^2,$$

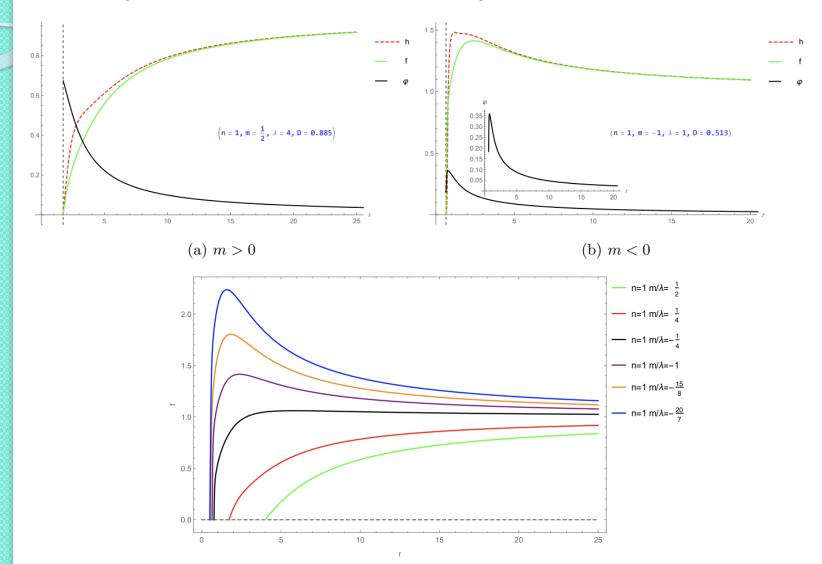
$$f(r) = \frac{r_h \left[e^{\left(-12\varphi_h^2\right)} \left(n^2 + r_h^2\right) \delta + 6\varphi_h^2 \left(1+\delta\right)^2 \lambda^4 \right]}{6n^2 \varphi_h^2 \left(1+\delta\right) \lambda^4} (r-r_h) + \mathcal{O}(r-r_h)^2,$$

$$\varphi(r) = \varphi_h - \frac{e^{6\varphi_h^2} r_h \delta}{2\varphi_h \left(1+\delta\right) \lambda^2} (r-r_h) + \mathcal{O}(r-r_h)^2.$$

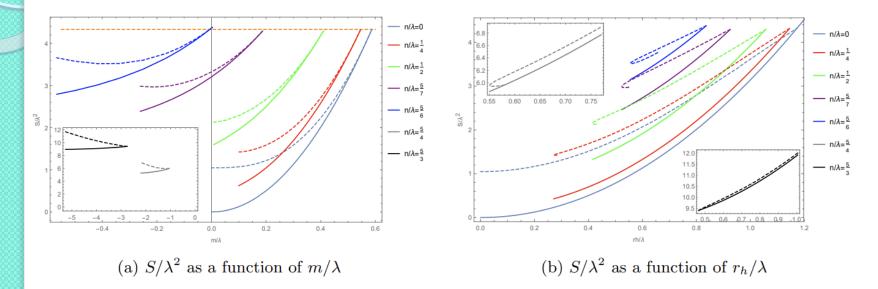
• Asymptotic form in the infinity

$$\begin{split} h\left(r\right) &= 1 - \frac{2M}{r} + \frac{D^2 - 2n^2}{r^2} + \frac{M\left(D^2 + 2n^2\right)}{r^3} + \frac{\left(4D^2M^2 + 6n^4 + 2n^2D^2\right)}{3r^4} \\ &+ \frac{M}{12r^5} \left\{ 63D^4 - 4D^2 \left(142M^2 + 65n^2 - 4\lambda^2\right) + 8 \left(80M^4 + 12n^2 - 51n^2\lambda^2\right) \right. \\ &+ 8M^2 \left(80n^2 - 21\lambda^2\right) \right\} + \dots, \\ f\left(r\right) &= 1 - \frac{2M}{r} - \frac{2n^2}{r^2} + \frac{M\left(D^2 + 6n^2\right)}{3r^3} + \frac{2D^2M^2 + 2n^2\left(D^2 + 3n^2\right)}{3r^4} \\ &+ \frac{M}{60r^5} \left\{ 55D^4 - 4D^2 \left(130M^2 + 50n^2 - 28\lambda^2\right) + 8 \left(80M^4 - 51n^2\lambda^2\right) \right. \\ &+ 8M^2 \left(80n^2 - 2\lambda^2\right) \right\} + \dots, \\ \varphi\left(r\right) &= \frac{D}{r} + \frac{DM}{r^2} - \frac{D\left[D^2 - 2\left(4M^2 + n^2\right)\right]}{6r^3} + \frac{DM\left(6M^2 - 2D^2 + 3n^2\right)}{3r^4} \\ &+ \frac{D}{120r^5} \left\{ 9D^4 + 24 \left[16M^4 + n^4 + 3n^2\lambda^2 + 3M^2\left(4n^2 - \lambda^2\right)\right] \\ &- 4D^2 \left(58M^2 + 13n^2\right) \right\} + \dots. \end{split}$$

The equations can be solved numerically



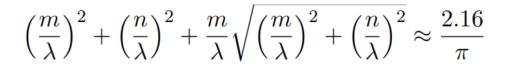
The entropy of the black hole

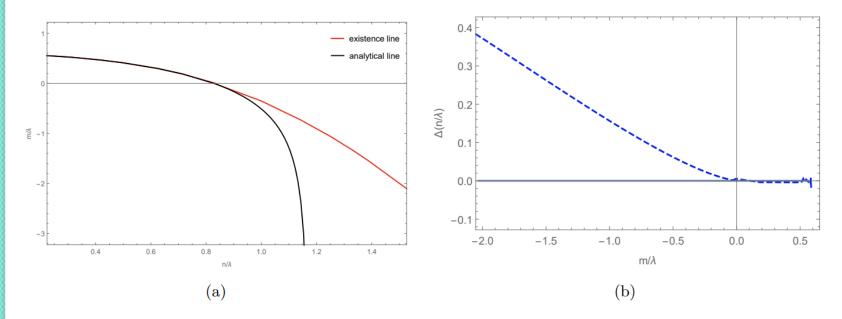


• The entropy of scalarized black hole is universal at the bifurcation point for different NUT parameter, and it is the highest value of the scalarized black hole.

 $S_{\text{scalar-free}}(m,n) < S_{\text{scalarized}}(m,n) \le S_{\text{scalarized}}(m_{\max}(n)) = S_{\text{scalar-free}}(m_{\max}(n))$

• The universal value at bifurcation points shows a relation between m and n





• The universal value only holds for positive mass parameter.



Conclusion and outlook

We constructed new Taub-NUT black holes with scalar hair or spin-2 hair in higher derivative gravity theories.

The thermodynamics, especially the definition of mass and NUT charge for new Taub-NUT black hole in higher derivative gravity is worth further investigation.

Thank you!